Electron Microscopy with Vortex Beams Carrying Orbital Angular Momentum

Jefferson Lab seminar February 12, 2015

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Thanks to





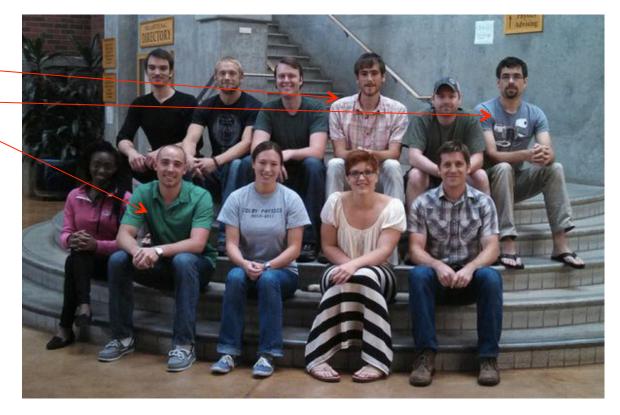
LDRD Carbon Cycle 2.0 "Electron Microscopy with Vortex Beams"



The work was supported by the grant DE-SC0010466 funded by the U.S. Department of Energy, Office of Science.

This data thanks to





<u>Dr. Peter Ercius</u> <u>Dr. Colin Ophus</u> Dr. Martin Linck (CEOS) Dr. Andreas Schmid





The Molecular Foundry A National Nanoscience Research Facility



Colin

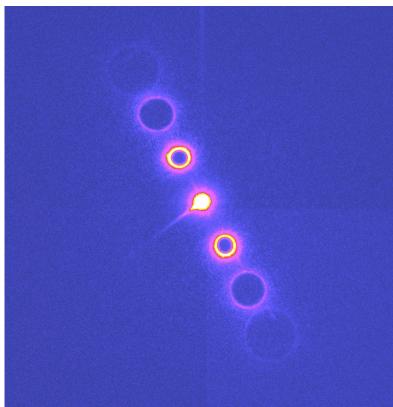


Peter

Overview

What is a vortex beam?:

- helical phase
- quantized orbital angular momentum
- How is a vortex beam produced?:
- spiral phase plate
- holographically (diffractive optics)
- detected?
- Is a vortex beam useful?:
- optical trapping/manipulation
- phase imaging/coherence filter
- magnetic sensitivity



What is a vortex beam?



Example:

Suppose we want a regular beam with flat wavefronts, a plane wave:

wavefunction amplitude represented by color and transparency

concave converging wavefronts

beam focus (waist)

convex diverging wavefronts

Another Example:

Suppose we want a "vortex beam" with azimuthal phase topological screw dislocation in wavefronts

wavefunction amplitude represented by color and transparency

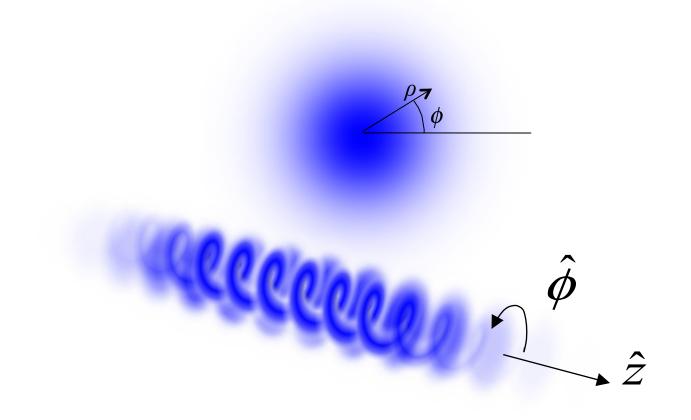
concave converging wavefronts

beam focus (waist)

convex diverging wavefronts

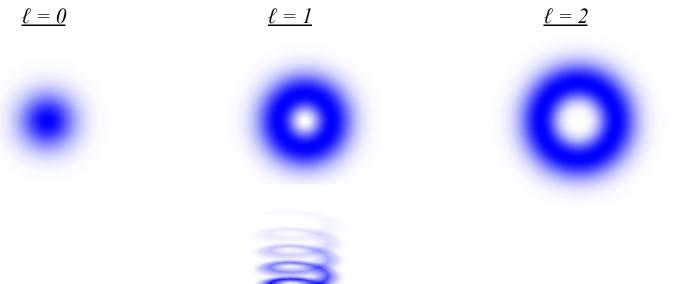
consider a plane wave: imprint with azimuthal phase

 $u(\mu(,\phi,z) \propto e^{-\left(\binom{\partial^2 p^2}{\partial \mu}\right)} e^{ikikz - im_{\ell}\phi}$



simplest optical vortex - Laguerre-Gaussian beams:

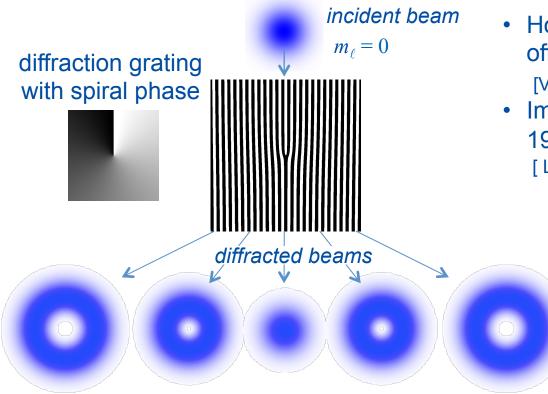
 $LG_{0m_{\ell}}(\rho,\phi,z) \propto \rho^{m_{\ell}} e^{-\left(\rho^{2} / w^{2}\right)} e^{ikz} e^{-im_{\ell}\phi}$



"screw"-shaped wavefronts

- "thread pitch" λ
- "number of starts" m_{ℓ}
- "lead" $m_\ell \lambda$

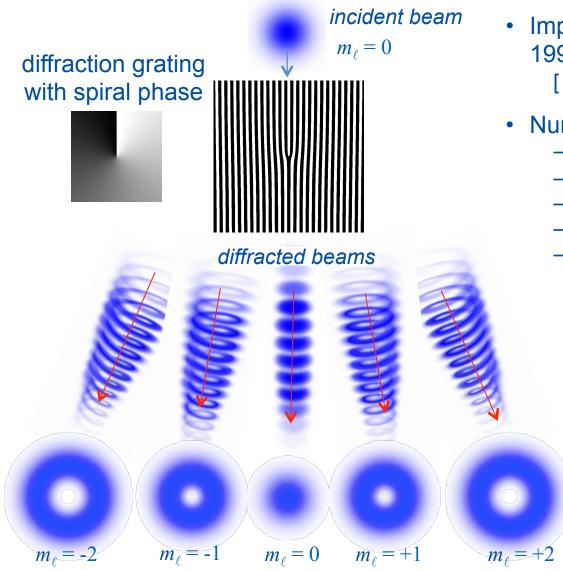
Light Optics: beams with helical phase made by physical holograms



- Holograms (spatial light modulator) offer precise way to produce
 [V. Bazhenov et al., JETP Letters. 52, 429 (1990)]
- Important topic in light optics since 1992*

[L. Allen et al., PRA 45, 8185 (1992)]

Light Optics: beams with helical phase made by physical holograms



 Important topic in light optics since 1992*

[L. Allen et al., PRA 45, 8185 (1992)]

- Numerous applications:
 - microscopy
 - astronomy
 - micromanipulation
 - communications
 - quantum state manipulation

L. Allen *et al.*, PRA 45, 8185 (1992): "Vortex beams" carry orbital angular momentum (OAM)

$$\psi(\rho, \phi, z) = f(\rho, z) e^{-im_{\ell}\phi}$$

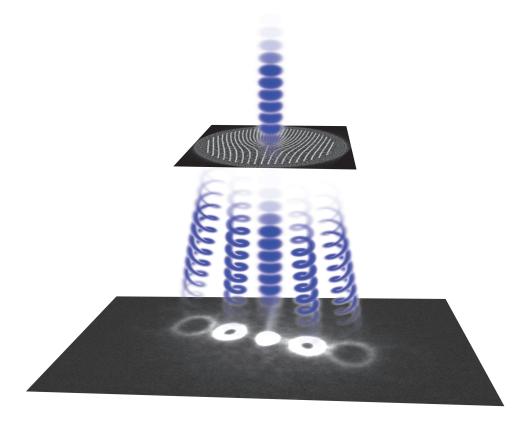
OAM operator: $\hat{L}_{z} = i\hbar \frac{\partial}{\partial \phi}$

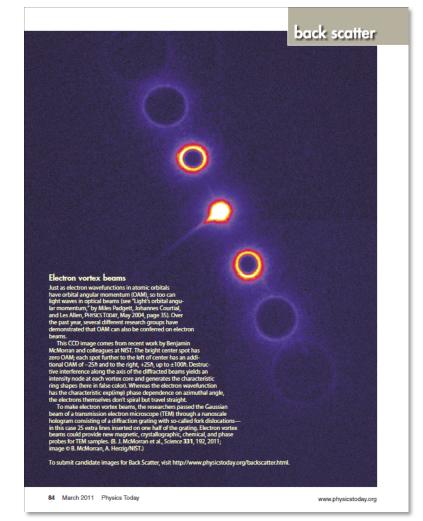
$$\hat{L}_{z} \begin{bmatrix} \psi_{m_{\ell}} \end{bmatrix} = \hbar m_{\ell} \begin{bmatrix} \psi_{m_{\ell}} \end{bmatrix}$$

$m_{\ell} = -1$ $\underline{L_z} = -\hbar$	$m_\ell = 1$ $L_z = +\hbar$	$m_{\ell} = 2$ $L_z = +2\hbar$	$m_{\ell} = 5$ $\underline{L_z = +5\hbar}$
	CORRECTION		
 "screw"-shaped beams "thread pitch" λ 			
• "number of starts" m_{ℓ} • "lead" $m_{\ell} \lambda$			

Can also make electron vortex beams

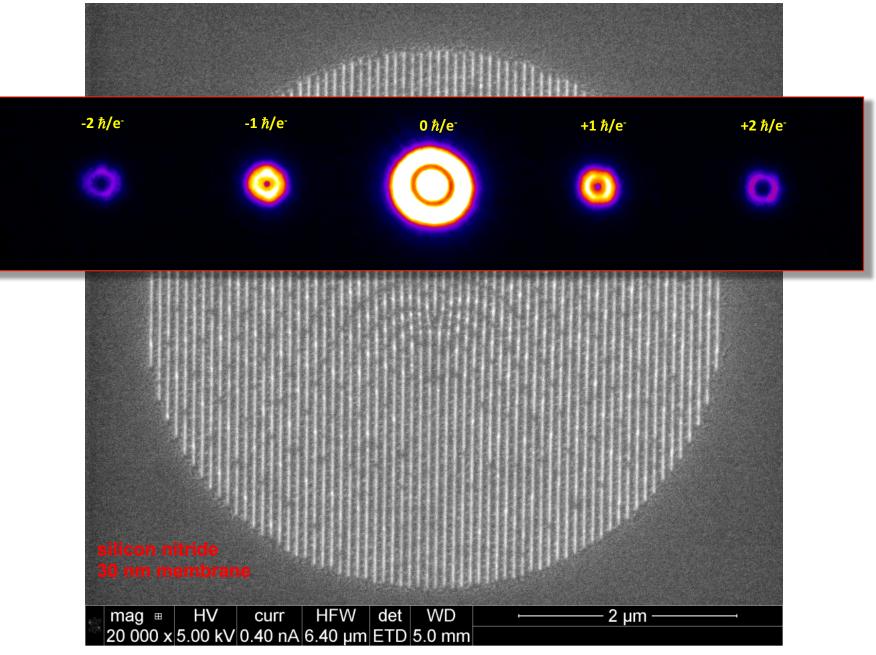
- 300 keV electrons (λ = 2 pm)
- nanofabricated diffractive holograms
- 39 µrad diffraction angle
- TEM



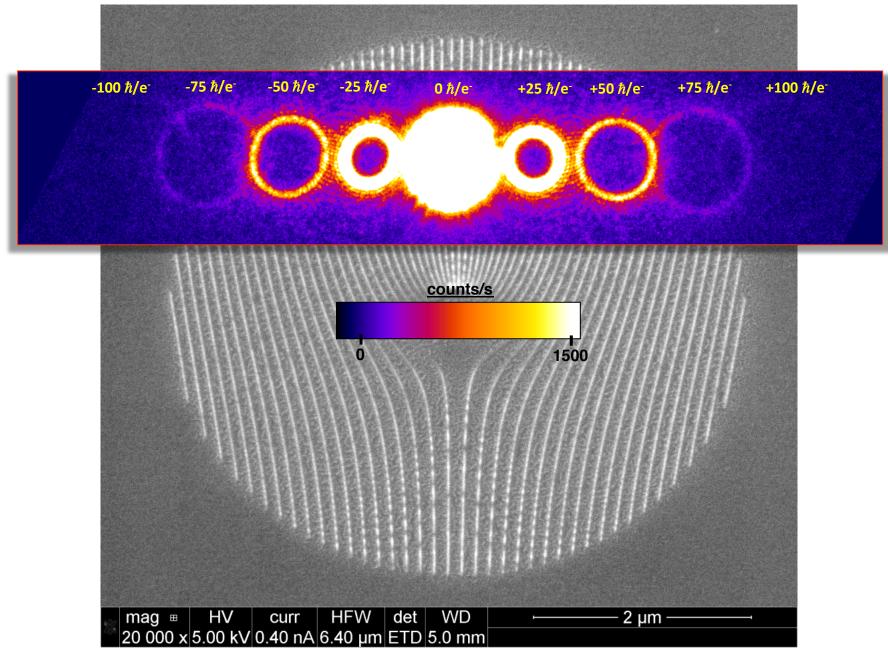


B.J. McMorran *et al.*, *Science* **331**, 192 (2011)

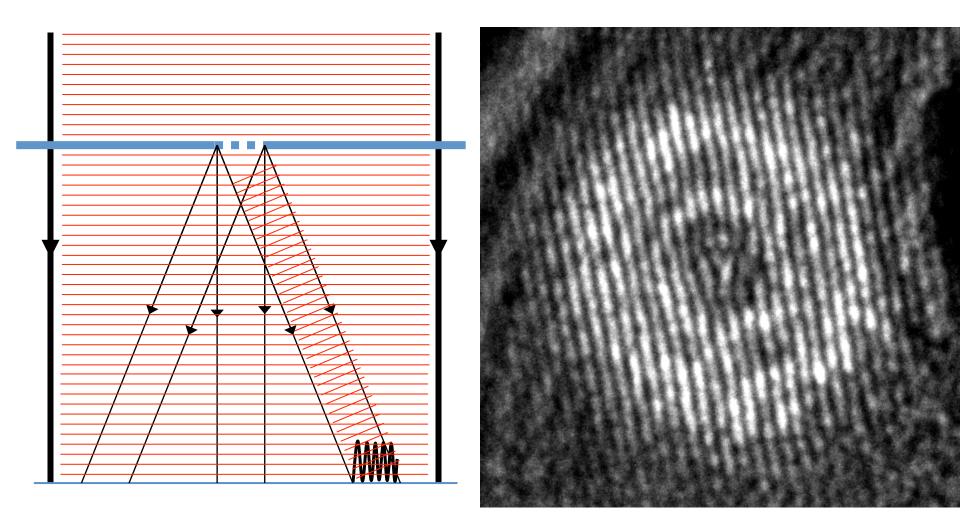
Nanofabricated diffraction hologram for electron vortex beams:



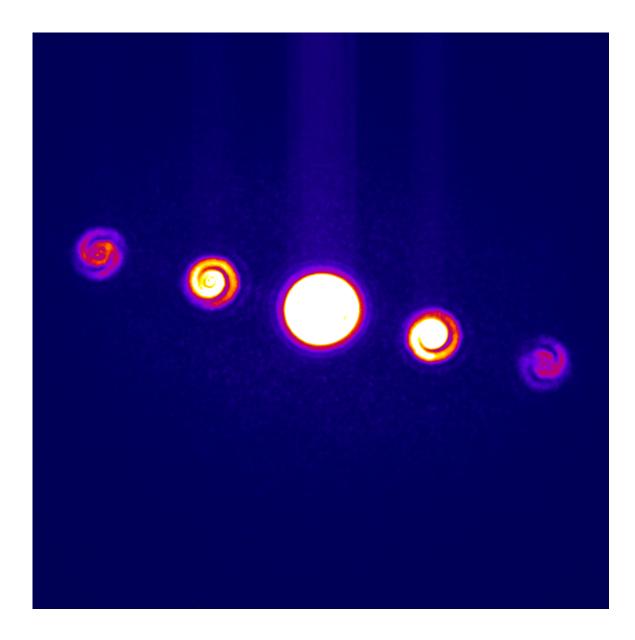
Nanofabricated diffraction holograms for electrons



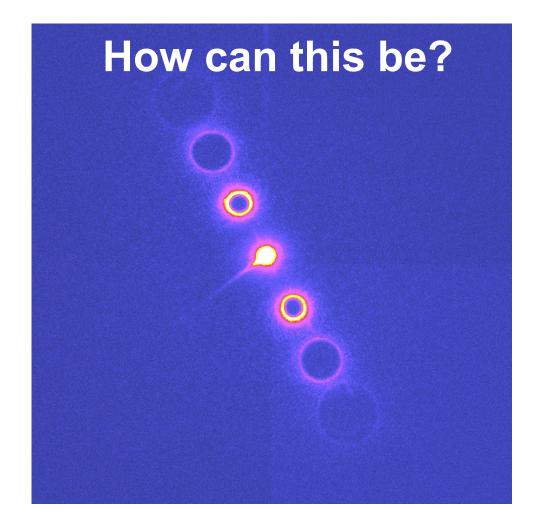
Confirming the helical phase of an electron vortex beam using interferometry



Measuring helical phase using electron interferometry



Stable orbital motion of an unbound *massive* particle Stable orbital motion of an unbound charged particle Orbital eigenstates in free space?



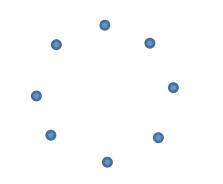
The electron is <u>not</u> orbiting in circles (not a helical trajectory)

The wavefunction represents the shape of the electron

The electron's propagation looks like this:



in the electron's reference frame orbital motion does not require closed trajectories



The electron vortex beam: a superposition of straight line trajectories

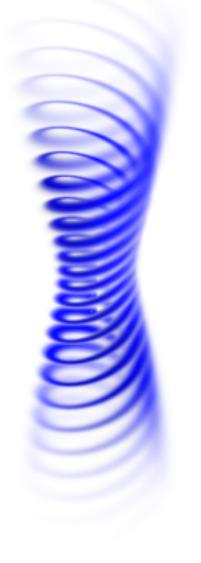


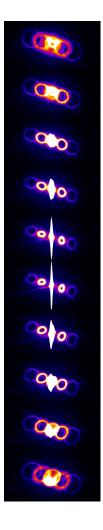
How we model free electrons with orbital angular momentum

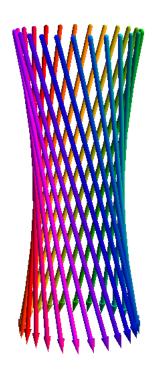
helical wavefunction model

current loop model

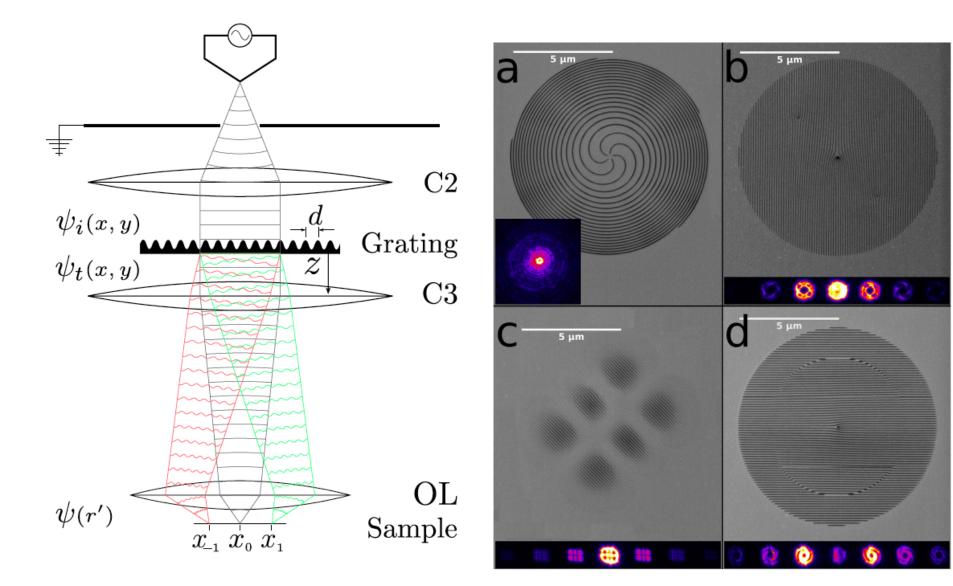
semiclassical trajectory model







This is one example of a more general ability to engineer electron wavefronts (or other short wavelength beams)

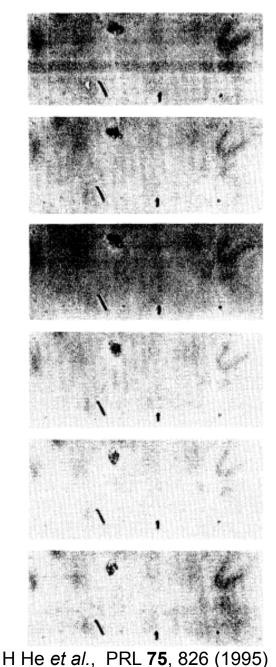


What do we do with an electron vortex beam?



exchange angular momentum with specimens

Optical ("light") vortex beams carry real angular momentum



- An optical vortex beam can torque objects
- This torque is quantized; i.e., $\ell\hbar$ per photon
- Optical beams (and particle beams) can carry two types of angular momentum:
 - -- spin angular momentum (SAM), AKA polarization
 - -- orbital angular momentum (OAM)

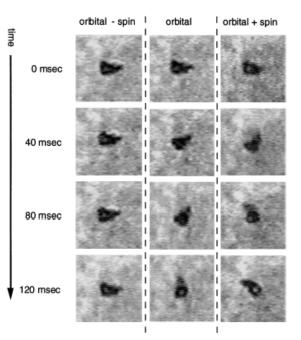
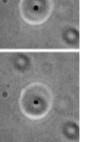


Fig. 2. Successive frames of the video image showing the stop-start behavior of a $2-\mu$ m-diameter Teflon particle held with the optical spanner.

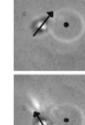
N. B. Simpson et al., Opt. Lett., 22, 52 (1997)

Optical ("light") vortex beams carry real angular momentum

Time

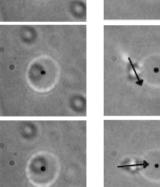


Orbit



Spin





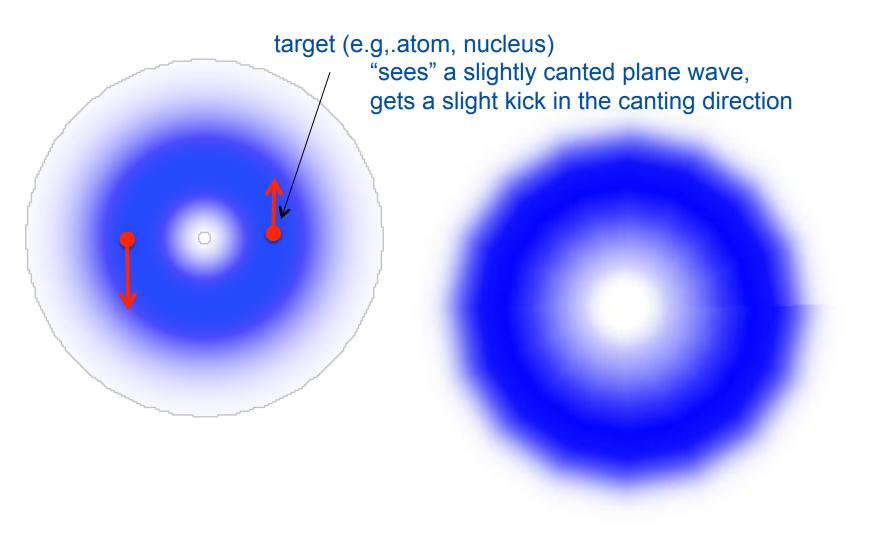
- Spin angular momentum (SAM) is "intrinsic"
 - circular polarization
 - transferred wherever photon is absorbed
- Orbital angular momentum (OAM) is "extrinsic"
 - property of wavefunction shape
 - transfer depends on relative beam/specimen size

Miles Padgett, "Light's twist", Proc. Royal Soc. Lon. A 470, 20140633 (2014):

"For beams that are large compared with the particle, the behaviour of SAM and OAM is different. Whereas the transfer of SAM causes particles to spin around their own axis, the transfer of OAM causes them to orbit around the beam axis"

A. T. O'Neil et al., PRL 88, 053601 (2002)

"Extrinsic" OAM – carried by the entire wave



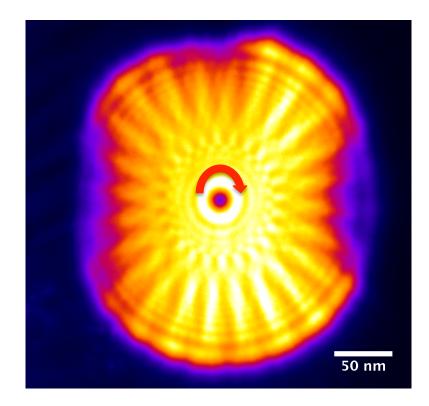
Can focused electron vortex beams induce angular momentum-dependent transitions?



How could we test this?

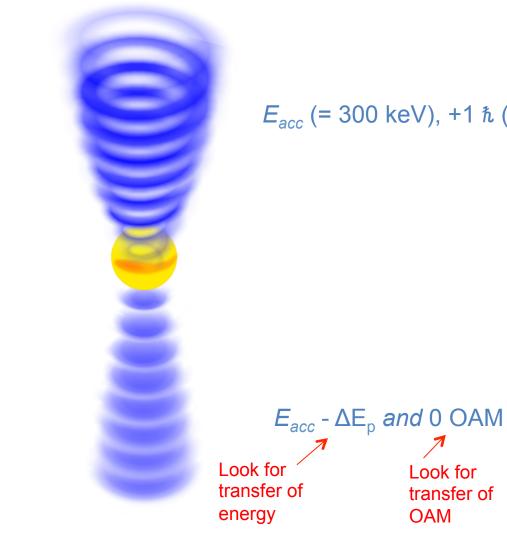
Central node (dark spot) indicates OAM

- Carries quantized orbital angular momentum ↔ azimuthal phase dependence
- Intensity drops to a minimum as $r \to 0$
- present throughout the entire path of the beam (remains through focus)
- The only way to make central dark spot disappear is by losing OAM



Basic experiment: Look at the donut hole!

Electron vortex beam incident on plasmonic structure transfers ΔE_p and ±1 \hbar (OAM)



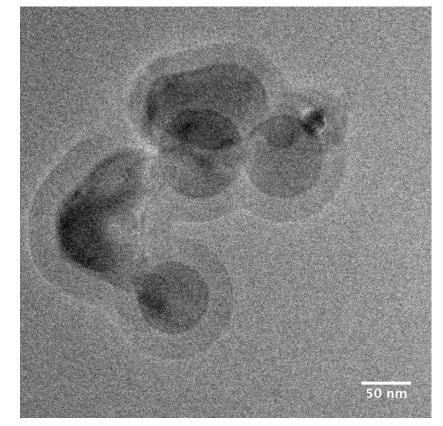
*E*_{acc} (= 300 keV), +1 ħ (OAM)

Electron vortices on chiral nanocluster

<u>Al/Al₂O₃ nanoparticles</u> 3.2 – 3.8 eV surface plasmon resonance [Scholl *et al.*, Nature **483**, 421 (2012)]

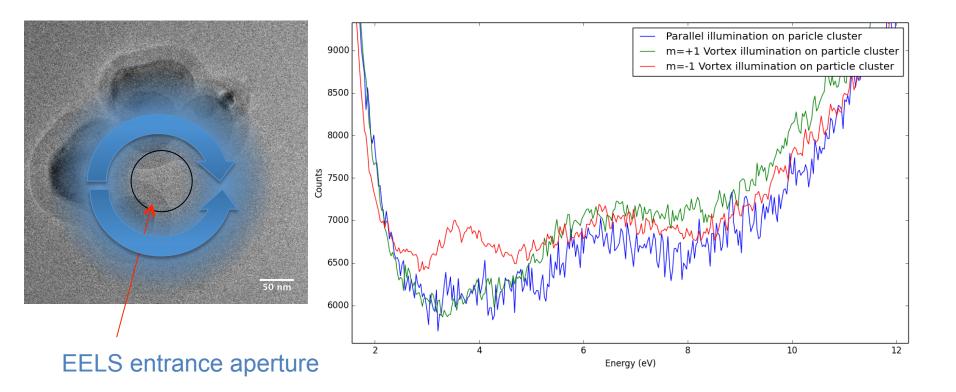
 \rightarrow high-enough E_{ρ} for us to measure in EELS!

chiral nanocluster:



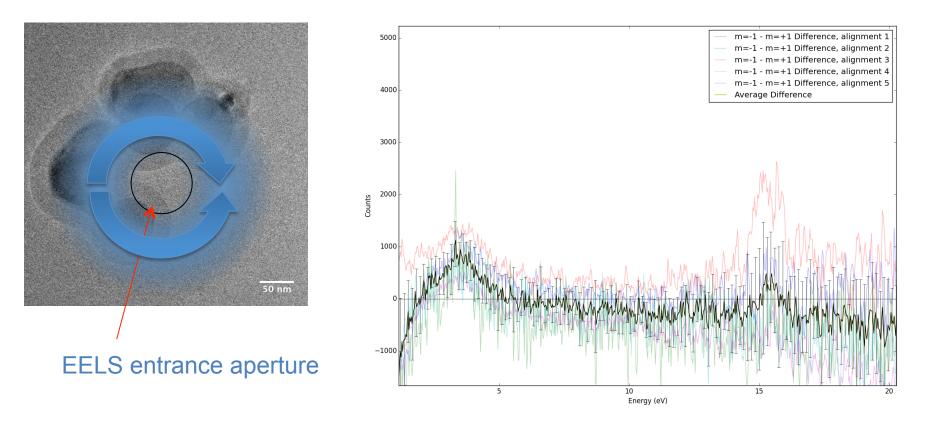
Electron vortices on chiral nanocluster

- Illuminate with three electron beams: -1, 0, and +1 OAM NOTE: <u>-1 and +1 beams have identical intensity distributions</u>
- · Record EELS spectra only for electrons scattered into the dark spot



Electron vortices on chiral nanocluster

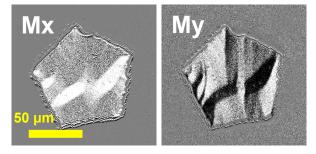
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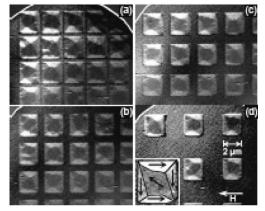
T. R. Harvey et al., M&M (2014)

magnetization is due to angular momentum of charges see it using beams that carry angular momentum

polarized light

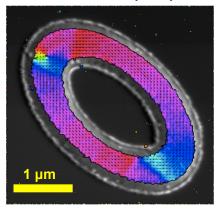


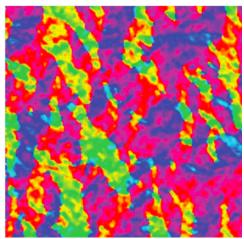
polarized X-rays



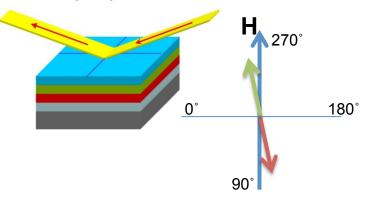
P. Fischer et al., Materials Today 9, 26 (2006)

spin polarized electrons





spin polarized neutrons



Magnetic dichroism at higher resolutions? OAM-dependent *atomic* scattering

Calculate scattering amplitude:

$$P_{i \rightarrow f} = \frac{2\pi}{\hbar} \left| \left\langle n'\ell'm' \right| H_{int} \left| n\ell m \right\rangle \right|^{2} \qquad (\text{Fermi's Golden rule})$$

interaction Hamiltonian: $H_{int} = \left\langle \psi_{2} \left| V(r) \right| \psi_{1} \right\rangle$
 $\psi_{2}(\mathbf{r}) = A_{2}e^{i\mathbf{k}_{2}\cdot\mathbf{r}}$
scattered electron $\psi_{1}(\mathbf{r}) = A_{1}(\rho)e^{i\mathbf{k}_{1}\cdot\mathbf{r}}e^{-im_{\ell}\phi}$
incoming electron vortex
Coulomb scattering: $V(r) = \frac{e^{2}}{4\pi\varepsilon_{0}} \left(\frac{-Z}{|\mathbf{r}|} + \sum_{j}^{Z}\frac{1}{|\mathbf{r} - \mathbf{r}_{j}|}\right)$

electron vortex STEM probe ≈ circularly polarized X-rays → circular dichroism using electron orbital angular momentum

new selection rules!

Transmitted electron energy loss (EELS):

$$P_{i \to f} \propto \delta \left(E_f - E_i - \Delta E \right) \begin{pmatrix} \ell' & 1 & \ell \\ m' & \Delta m_\ell & m \end{pmatrix}^2$$

X-ray absorption (XMCD):

$$P_{i \to f} \propto \delta \left(E_f - E_i - E_\gamma \right) \begin{pmatrix} \mathbf{l'} & \mathbf{l} & \mathbf{l} \\ m' & \boldsymbol{\sigma}_{\pm} & m \end{pmatrix}^2$$

XMCD-like contrast?

photons transfer energy and angular momentum to atom → X-ray magnetic circular dichroism (XMCD)

Helicity-Dependent X-ray Absorption Magnetic Imaging (resolution ~10 nm) absorption coefficient 30ÅFe/W(110) 3d M^+ band - M⁻ imaging with Difference circular polarized X-rays Negative Positive helicity helicity 700 760 720 740 $2p_{3/2}$ Photon energy (eV) $2p_{1/2}$ P. Fischer et al., Materials Today 9, 26 (2006)

electrons with orbital angular momentum (1 ħ) may also provide this sensitivity?

Electron vortex: inelastic collisions with atoms

Scattering amplitude:

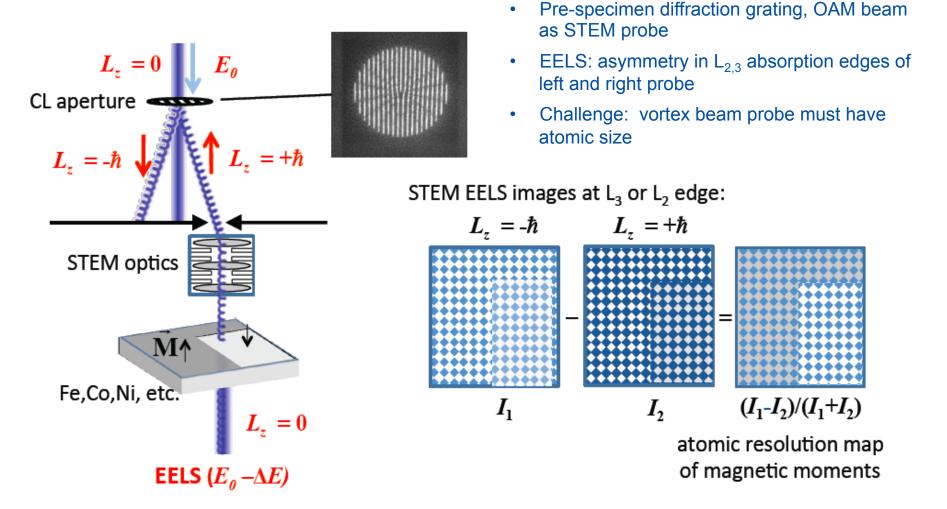
$$P_{i \rightarrow f} \propto \delta\left(E_{f} - E_{i} - \Delta E\right) \left| \sum_{j}^{Z} \sum_{\lambda=0}^{\infty} R_{n't',\lambda|m_{\ell}|,n\ell}(q) \Omega_{t'm',\lambda m_{\ell},\ell m} \right|^{2}$$
Electron
Energy
Loss
Spectroscopy
Need atomic-sized beam
(integral over radial coordinates)
Quantized energy transfer
a need atomic scale beam
b need atomic scale beam
c magnetic dichroism
c magnetic dichroism
c magnetic dichroism
c magnetic dichroism

TEAM 1

electron vortex beams

Electron Magnetic Orbital Dichroism (EMOD)

→ <u>STEM/EELS with vortex beam probe</u>



Calculations & Optical Experiments:

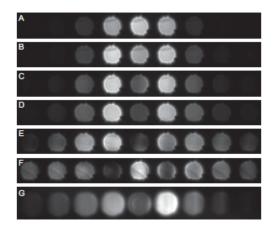
- To transfer OAM to single atoms, vortex beam needs to be atomic scale
- To use in a STEM instrument, need high current

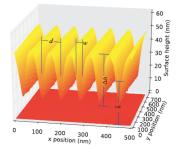
➔ optimize diffractive optics for electrons

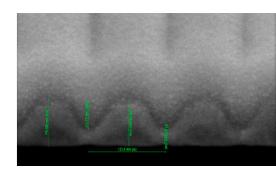
Optimizing diffraction efficiency

- **1.** make lots of dose arrays (various fabrication parameters):
- groove depth (5 nm 40 nm)
- groove width (12 nm 60 nm)
- groove spacing (24 nm- 150 nm)
- groove shape (sinusoidal to asymmetric triangular)
- grating material (SiNx (various sources), Si, Pt)
- fabrication method (FIB milling [various], FIB-deposited Pt, EBL [various
- 2. measure electron diffraction efficiency
- grating in specimen holder
- TEM LAD
- measure:
 - incident beam current
 - total transmitted current
 - current in individual diffracted beams
- 3. measure fabricated grating surface
- AFM, FIB & SEM
- 4. model
- 5. repeat





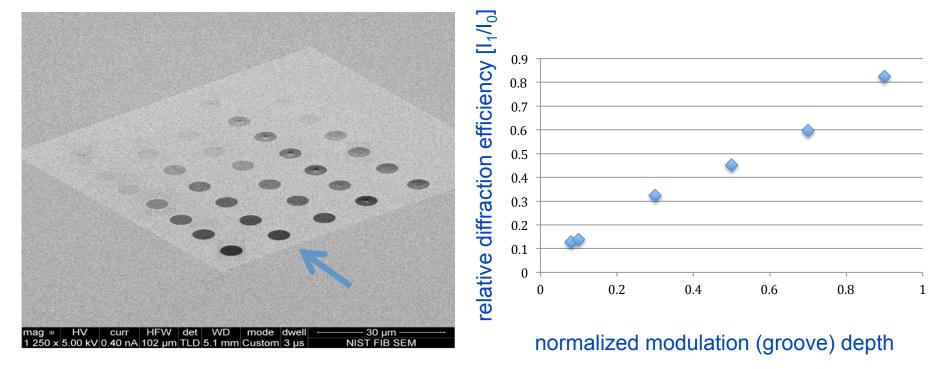




Maximizing diffraction efficiency of electron holograms

dose arrays (many)



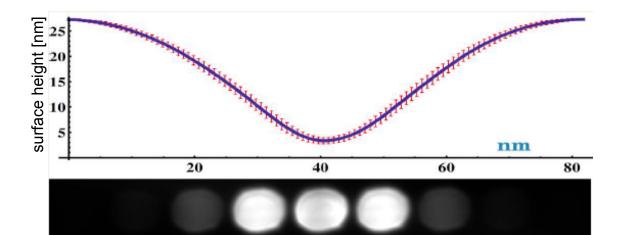


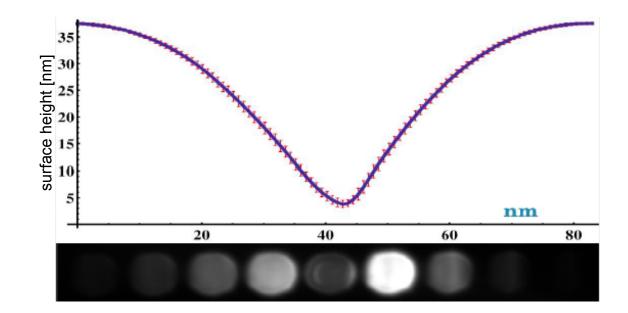
Grating groove shape influence diffraction efficiency

symmetric grating

AFM profile:

diffraction pattern:



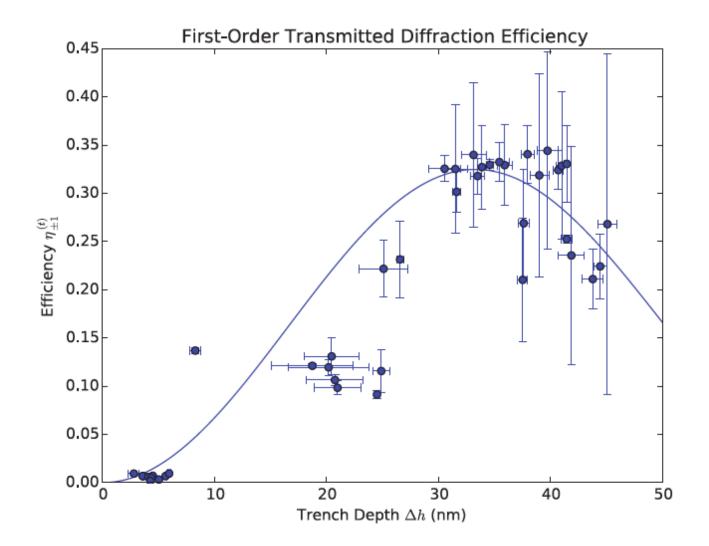


blazed grating

AFM profile:

diffraction pattern:

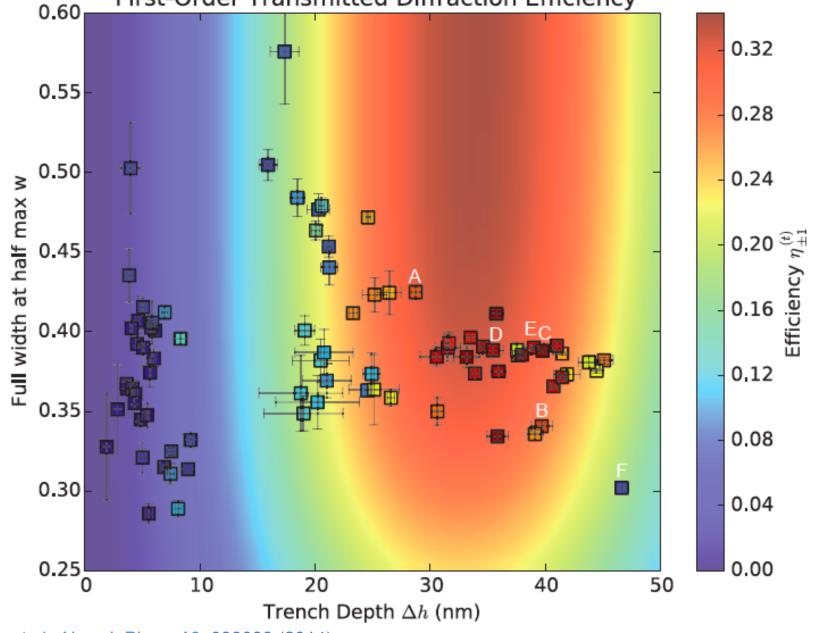
Diffractive electron optics diffraction efficiency



T. R. Harvey et al., New J. Phys. 16, 093039 (2014).

Optimizing diffraction efficiency – exploration of parameter space

First-Order Transmitted Diffraction Efficiency



T. R. Harvey et al., New J. Phys. 16, 093039 (2014).

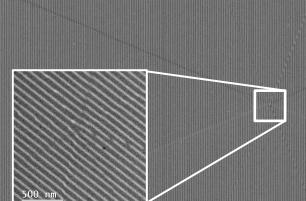
High spatial resolution High spatial coherence Large area Good isolation of diffracted orders Precisely controlled phase Smaller probes (high resolution)

2010 grating

- 5 µm diameter
- 100 nm line spacing
- 10 µm diameter
- 75 nm line spacing

det WD m ETD .0 mm

2011 grating



Current gratings (background image, drawn to

scale)

- 80 µm diameter
- 50 nm line spacing

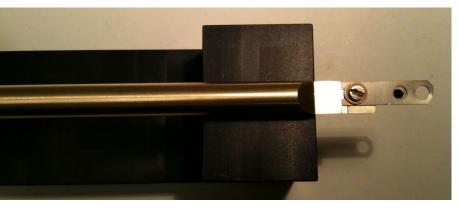
Enables:

- atomic resolution
- wide field-of-view

T. R. Harvey et al., New J. Phys. 16, 093039 (2014).

Aberration-corrected electron vortex STEM probes

Modified condenser/objective apertures:



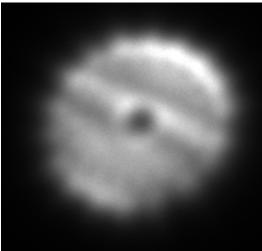
Installed in condenser lens system

- TEAM I (NCEM)
- Titan-X (NCEM)
- FEI Titan (UO-CAMCOR)
- FEI Titan (NIST-MML)
- JEOL 200CF (UIC)



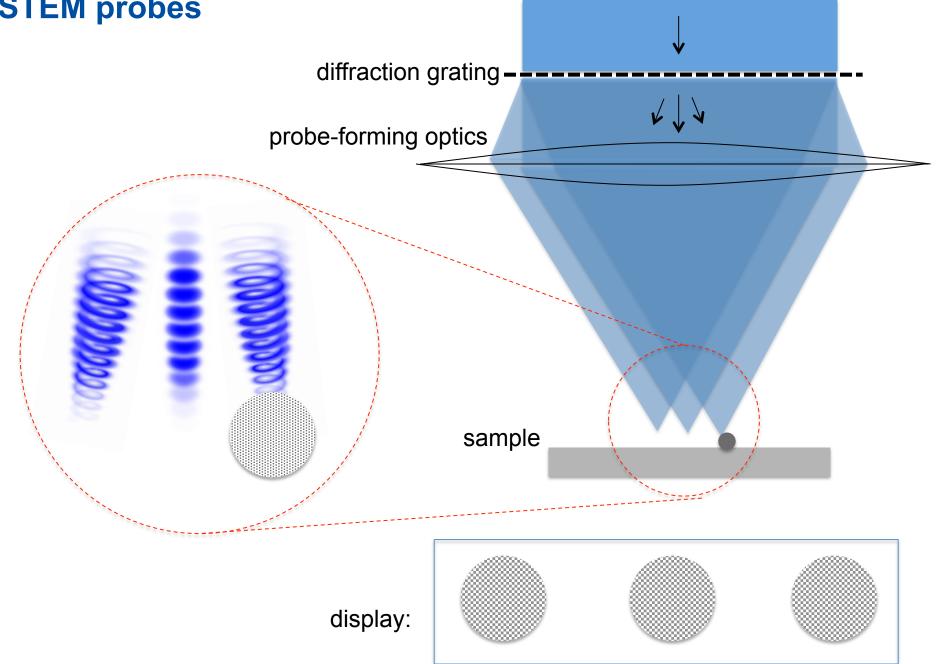
(Peter Ercius)

Aberration-corrected STEM probe (defocused):

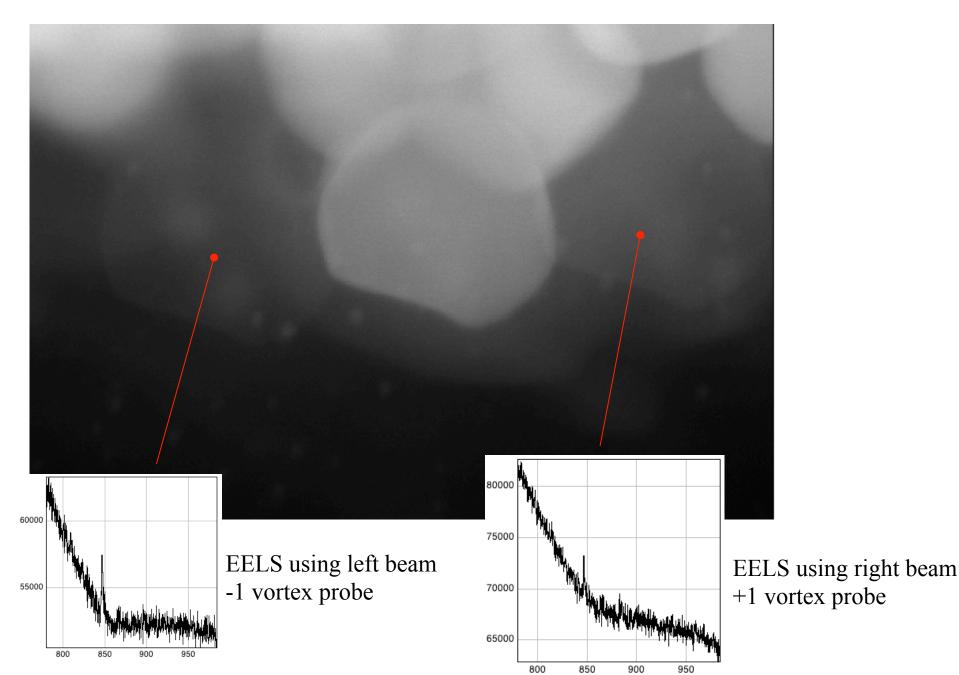


McMorran, B. United States Patent: 8680488 - System and

Imaging with multiple Angstrom-scale electron vortex STEM probes

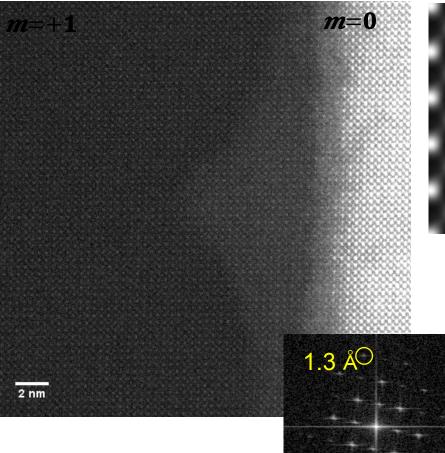


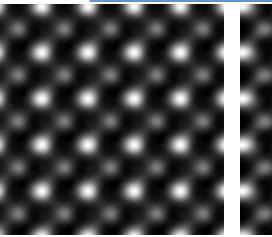
Multiple STEM probes \rightarrow multiple images of the same object



SrTiO₃ imaged by 0th & 1st order focused probes

Mean unit cells





Side beam

Center beam



Understanding electron vortex in external fields

Unlike optical OAM, electron OAM has a magnetic moment

 $\mu_z = g \,\mu_B \,\ell$

Couples to external fields:

G. M. Gallatin and B. McMorran, Phys. Rev. A 86, 012701 (2012):

full path integral solution

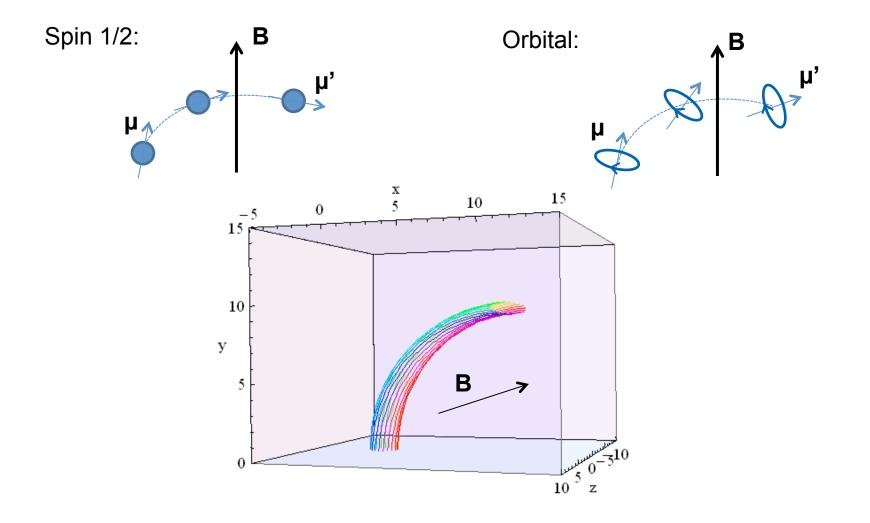
magnetic field \rightarrow rotating coordinate system

ladder operators \rightarrow

propagate OAM wavefunction through field equivalent to propagiting conventional wavefunction then adding OAM

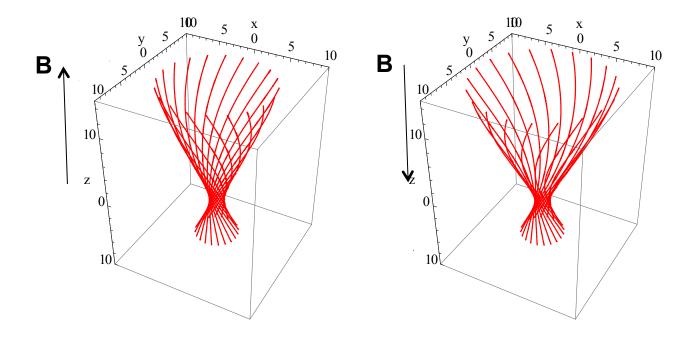
Use semiclassical ray picture to model eLG in external fields: Transverse magnetic field

magnetic moment does not precess same as beam path \rightarrow gyromagnetic ratio g is not 2 like electron spin.



Use semiclassical ray picture to model eLG in external fields: Longitudinal magnetic field

Zeeman interactions with longitudinal fields (working with G. Gallatin, NIST)



Key Points

Nanofabricated diffractive optics

→ fine control over quantum states of free electrons BUT this requires a transverse spatial coherence

- Free electrons can carry quantized orbital angular momentum
 - → theoretically, this can induce electronic transitions in atoms, and perhaps even nuclear transitions in atomic nuclei

BUT angular momentum is delivered on the lengthscale of the beam

- Other possible experiments at JLab:
 - Generation of electrons with both spin and orbital angular momentum (electron vector beams)
 - Utilize spin-orbit coupling to induce spin polarization
 - Generation of EUV and X-ray vortex beams using electron vortex beams?

Thanks to





LDRD Carbon Cycle 2.0 "Electron Microscopy with Vortex Beams"



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Electron Vortex Collaborations

Ben McMorran Tyler Harvey Jordan Pierce Jordan Chess

Peter Ercius (LBNL) Colin Ophus (LBNL) Andreas Schmid (LBNL)

Martin Linck (CEOS)

Juan-Carlos Idrobo (ORNL)

Ondrej Krivanek (Nion)

Phil Batson (Rutgers)

Nigel Browning (PNNL) Rodney Herring (UVic)



UNIVERSITY

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Robert Klie

Preliminary Lorentz TEM results in amorphous FeGd w/ James Lee, Sujoy Roy (ALS) Sergio Montoya, Eric Fullerton (UCSD)

