

Differential Algebraic Methods for Space Charge Modeling and Applications to the University of Maryland Electron Ring

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Motivation

- Extract relevant information from the transfer map of a given system using normal form methods.
- Include effects not normally included in a transfer map.
- Provide new insights into the effects of space charge on particle beams.

Outline

- Overview of the Research
- The University of Maryland Electron Ring (UMER)
- Mathematical Tools
- UMER's modeling, and early experiments
- Adding space charge to the simulation
- Parallelization of the space charge code
- Alternatives for space charge calculation
- Summary

UMER



- 10 Kev electron ring
- 15 Beam position monitors.
- 16 Intercepting Phosphor screens
- 1 fast phosphor screen
- 3 Current monitors

Geometry (Injection)

• The off-axis nature of the injection line made modeling the Y-section also difficult.



Geometry (dipoles)



 COSY dipoles assume that the steering is perfect, so a beam that enters straight will exit straight. If we want to use them for steering we need to adjust for how a dipole will change the path of the beam.



Differential Algebras and Numerical Derivatives

First we introduce the $_1D_1$ Differential Algebra, which is a first order one variable vector.

$$(q_0, q_1) + (r_0, r_1) = (q_0 + r_0, q_1 + r_1)$$

$$n(q_0, q_1) = (nq_0, nq_1)$$

$$(q_0, q_1)(r_0, r_1) = (q_0r_0, q_0r_1 + r_0q_1)$$

It follows that,

$$(f(x), f'(x)) = f(x+d)$$

 $d = (0, 1)$

As an example,

$$f(x) = 2x^{2} - x + 3$$

$$f'(x) = 4x - 1$$

$$f(2) = 9$$

$$f'(2) = 7$$

$$f(2+d) = f(2,1)$$

$$= 2(2,1)(2,1) - (2,1) + (3,0)$$

$$= 2(4,4) - (2,1) + (3,0)$$

$$= (8,8) - (2,1) + (3,0)$$

$$= (9,7)$$

Normal Form Methods



 $\mathcal{N} = \mathcal{A} \circ \mathcal{M} \circ \mathcal{A}^{-1}$



External Fields

• These are modeled using Strang Splitting.







External Fields (cont'd)



$\mathcal{M}(L/2)$ +O(L³)



Earth's Magnetic Field



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Steering (no fields)





Steering (Earth's Field)



Steering (Earth's field with Dipole Corrections)



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Steering Comparison

• COSY fitted steering v. field integration offsets.



Steering Comparison

• COSY fitted steering v. field integration offsets.



Comparison of Predicted to Measured Trajectories



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Extended Comparison



Tunes/Matching Dynamic Aperture



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Dispersion Evolution



Dispersion Evolution (cont'd)



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Tunes at Higher Order

Pencil Beam		$7 \mathrm{mA}$ Beam	
Order	$\mu(\delta)$	Order	$\mu(\delta)$
0	.747	0	.63722
1	-7.1058	1	-4.4947
2	111.479	2	-25.83
3	-4787.3359	3	-969.6789
$23 \mathrm{mA}$ Beam		80 mA Beam	
Order	$\mu(\delta)$	Order	$\mu(\delta)$
0	.6796	0	.6900
1	-6.172	1	-5.0174
2	-11.2667	2	-31.2928
3	2034.3401	3	-1221.4241

Resistance to Errors (Magnet Placement)



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Resistance to Errors (Magnet Current)



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Space Charge

- Space charge has eluded normal form analysis for some time
- A Differential Algebraic approach could allow us to determine how it affects other quantities
- Allows us to understand UMER better
- Allows for investigations into the intensity frontier

The Distribution Function

$$\rho(x, y) = \sum_{i} \delta(x - x_{i}) \delta(y - y_{i}) \rightarrow \rho(x, y) = \sum_{j} \sum_{k} C_{jk} x^{j} y^{k}$$

• If two distributions have the same moments then they are mathematically indistinguishable.



Distribution: (Cont'd)

 $M_{ij} = \iint x^i y^j \sum \sum C_{lm} x^l y^m dx dy$ m

$$M = TC$$

$$C = T^{-1}M$$

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Potential Calculation

 Since we now have a Taylor series for the distribution can't we just integrate and find the potential.

 No, Since the expansion is occurring inside the distribution Singularities become an issue.

$$G(r,r') = \ln(|r-r'|); \lim_{r \to r'} (G(r,r')) = -\infty$$







$$\int_{c}^{d} \int_{a}^{b} \ln(\sqrt{(x-x_{0})^{2}+(y-y_{0})^{2}})dxdy$$

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$$u_{1} = \frac{x - x_{0}}{b - x_{0}}; u_{2} = \frac{y - y_{0}}{d - y_{0}}; \lambda_{1} = (b - x_{0}); \lambda_{2} = (d - y_{0});$$
$$\lambda_{1}\lambda_{2}\int_{0}^{1}\int_{0}^{1}\ln(\sqrt{(\lambda_{1}u_{1})^{2} + (\lambda_{2}u_{2})^{2}})dxdy$$

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• This is done using the following set of coordinate transformations:

$$w_{1} = u_{1}; w_{2} = \frac{u_{2}}{u_{1}}$$
$$\lambda_{1}\lambda_{2} \int_{0}^{1} \int_{0}^{1} w_{1} \ln(\sqrt{\lambda_{1}^{2}w_{1}^{2} + \lambda_{2}^{2}w_{1}^{2}w_{2}^{2}}) dw_{1}dw_{2}$$



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• This is done using the similar set of coordinate transformations:

$$w_{2} = u_{2}; w_{1} = \frac{u_{1}}{u_{2}}$$
$$\lambda_{1}\lambda_{2}\int_{0}^{1}\int_{0}^{1}w_{2}\ln(\sqrt{\lambda_{2}^{2}w_{2}^{2} + \lambda_{1}^{2}w_{2}^{2}w_{1}^{2}})dw_{1}dw_{2}$$

$$\lambda_{1}\lambda_{2}\int_{0}^{1}\int_{0}^{1}w_{1}\ln(w_{1}) + \ln(\sqrt{\lambda_{1}^{2} + \lambda_{2}^{2}w_{2}^{2}})dw_{1}dw_{2}$$
$$+ \lambda_{1}\lambda_{2}\int_{0}^{1}\int_{0}^{1}w_{2}\ln(w_{2}) + \ln(\sqrt{\lambda_{1}^{2}w_{1}^{2} + \lambda_{2}^{2}})dw_{1}dw_{2}$$

$$\lim_{x\to 0} (x\ln(x)) = 0$$

Accuracy



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Example

$$\frac{r_m}{r_0} = 1 + 5.87 \times 10^{-5} \frac{I}{(\gamma^2 - 1)^{\frac{3}{2}}} (\frac{z}{r_0})^2$$

Method	Growth	
Edge Point x	35.27 %	
Edge Point y	35.30 %	
Map Element x	31.21 %	
Map Element y	31.34 %	





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Tune Measurement



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Space Charge Serial



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Space Charge Parallel



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Parallel Scaling (cont'd)



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Limitations

























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Fast Multipole Method

- Uses multipole expansions for distant particles
- Uses direct coulomb interactions for close particles
- Currently used in solid state physics, fluids, and chemistry



FMM: Examples



$$\phi(z) = \sum_{i} \left(Q \log(z_i) + \sum_{k} \frac{a_k}{z_i^k} \right) + \sum_{j} q_j \log(z_j)$$



Timing Comparison



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Distribution Tracking Experiments





Experimental Comparison



Summary

- Differential Algebraic methods can be used to model the single and multiple particle effects of a system.
- The University of Maryland Electron Ring is a useful tool for verifying machine codes.
- Adding the effects of space charge to the transfer map of a system is both possible and feasible
- New insights can be found using this method in conjunction with normal form analysis
- Can even analyze complex distributions using the fast multipole method

Questions?



Map Overview

 A map is a method of advancing particles which takes the form

 $z_f = M z_i$

$x_{f} = (x_{f} | x_{i})x_{i} + (x_{f} | p_{xi})p_{xi} + (x_{f} | x_{i}^{2})x_{i}^{2} + (x_{f} | x_{i}p_{xi})x_{i}p_{xi} + (x_{f} | p_{xi}^{2})p_{xi}^{2}$ $p_{xf} = (p_{xf} | x_{i})x_{i} + (p_{xf} | p_{xi})p_{xi} + (p_{xf} | x_{i}^{2})x_{i}^{2} + (p_{xf} | x_{i}p_{xi})x_{i}p_{xi} + (p_{xf} | p_{xi}^{2})p_{xi}^{2}$

Legendre Polynomials

• Legendre Polynomials have the following orthogonality property.

$$\int_{-1}^{1} P_m(x) P_n(x) dx = \frac{2}{2n+1} \delta_{mn}$$

 If we assume that the distribution can be modeled as a sum of legendre polynomials, we can easily find the coefficients.

$$\rho(x) = \sum_{n} C_{n} P_{n}(x)$$

$$\int_{-1}^{1} P_m(x)\rho(x)dx = C_m \frac{2}{2m+1}$$

Method Comparison





Three Dimensions





