



FERMI NATIONAL ACCELERATOR LABORATORY

US DEPARTMENT OF ENERGY

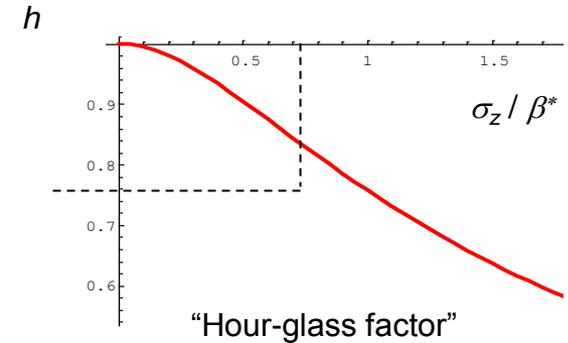
Muon Collider Ring Lattice Design

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Vadim Alexakhin (JINR, Dubna)

$$\langle \mathcal{L} \rangle = f_0 \frac{n_b N_\mu^2}{4\pi\epsilon_\perp \beta^*} h \times \frac{1}{2} \mathcal{F}_{rep} \sim \frac{P_\mu \xi}{C\beta^*} h\tau$$



P_μ – average muon beam power (limited by the P-driver power)

$\xi = \frac{r_\mu N_\mu}{4\pi\gamma\epsilon_\perp}$ – beam-beam parameter (limited by particle stability, $\xi < 0.1$?)

C – collider circumference (limited from below by available B-field)

τ – muon lifetime

β^* – beta-function at IP (limited from below by chromaticity of final focusing and aperture restrictions in IR magnets),

small β^* requires small $\sigma_z \Rightarrow$ large $\sigma_p / p = \epsilon_{||} / \sigma_z$

The recipe:

- Pack as many muons per bunch as the beam-beam effect allows (in practice this means 1 bunch/beam)
- Make beams round to maximize the beam-beam limit
- Develop new chromatic correction scheme to reduce β^*
- Do not leave free spaces to reduce C (also good for neutrino radiation)

What we would like to achieve compared to other machines:

	MC	Tevatron	LHC
Beam energy (TeV)	0.75	0.98	7
β^* (cm)	1	28	55
Momentum spread (%)	>0.1	<0.01	0.0113
Bunch length (cm)	1	50	15
Momentum compaction factor (10^{-3})	0.05	2.3	0.322
Geometric r.m.s. emittance (nm)	3.5	3	0.5
Particles / bunch (10^{11})	20	2.7	1.15
Beam-beam parameter, ξ	0.1	0.025	0.01

Muon collider is by far more challenging:

- much larger momentum acceptance with much smaller β^*
- ~ as large Dynamic Aperture (DA) with much stronger beam-beam effect
- **New ideas for IR magnets chromaticity correction needed!**

- 1996 by Carol J., A. Garren
 - 1996 by K.Oide
 - “Dipole first” (2007)
 - Eliana’s “synthetic” (2009)
 - Asymmetric dispersion
 - “Flat top”
- } ~ satisfy the requirements

1996 designs (especially by K.Oide) had extremely high sensitivity to field errors

Montague chromatic functions :

$$A_x = \frac{\partial}{\partial \delta_p} \alpha_x - \frac{\alpha_x}{\beta_x} \frac{\partial}{\partial \delta_p} \beta_x, \quad B_x = \frac{1}{\beta_x} \frac{\partial}{\partial \delta_p} \beta_x,$$

$$W_x = \sqrt{A_x^2 + B_x^2},$$

$\alpha_{x,y} = -\beta_{x,y}'/2$, $\beta_{x,y}$ are Twiss lattice functions,
 δ_p is relative momentum deviation.

Equations for chromatic functions

$$A_x' = 2\phi_x' B_x - \beta_x (K_1 - D_x K_2),$$

$$B_x' = -2\phi_x' A_x$$

K_1 , K_2 are normalized quadrupole and sextupole gradients,
 D_x is dispersion function: $D_x = dx_{c.o.}/d\delta_p$

The receipt:**Kill A's before they transform into B's !**

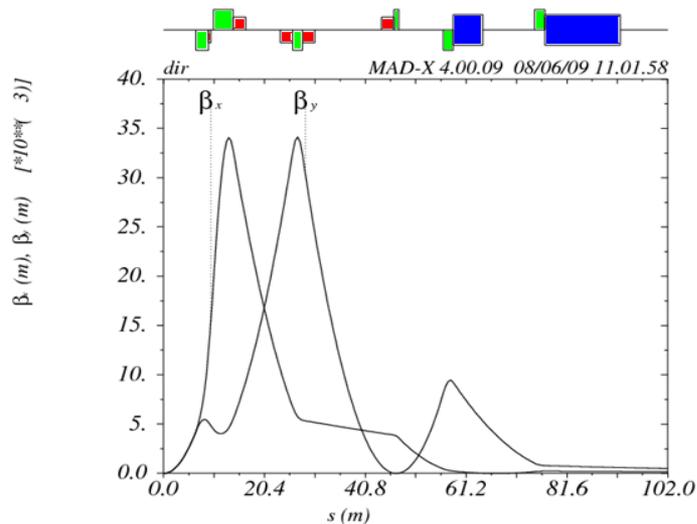
- difficult to achieve in both planes
- horizontal correction requires 2 sextupoles 180° apart to cancel spherical aberrations

$B_{x,y}$ are most important since they determine modulation of phase advance $\phi_{x,y}$

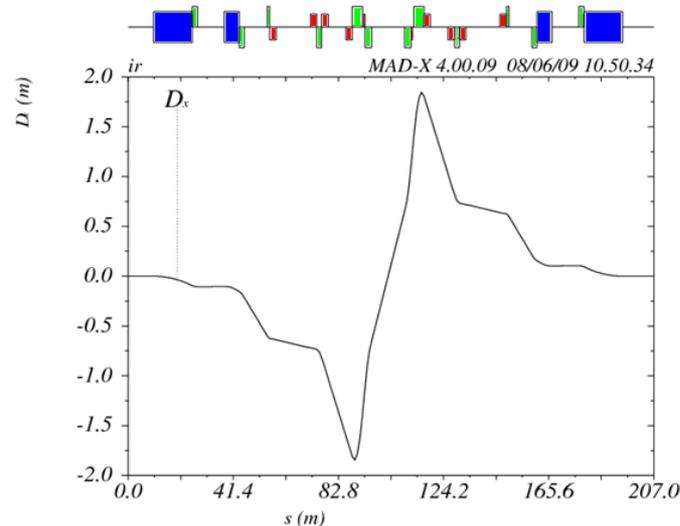
$$\phi_x(s, \delta_p) = \frac{1}{1 + \delta_p} \int_0^s \frac{ds'}{\beta_x(s', \delta_p)}$$

$A_{x,y}$ are created first, and then converted into $B_{x,y}$ as phase advances $\phi_{x,y}$ grow

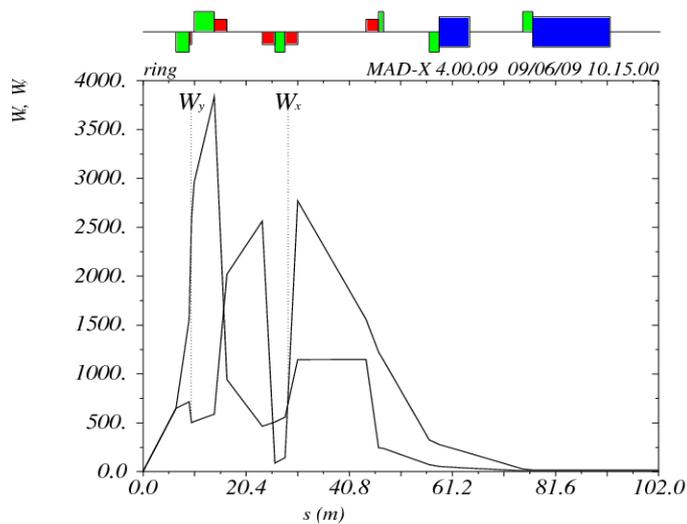
Half-IR optics



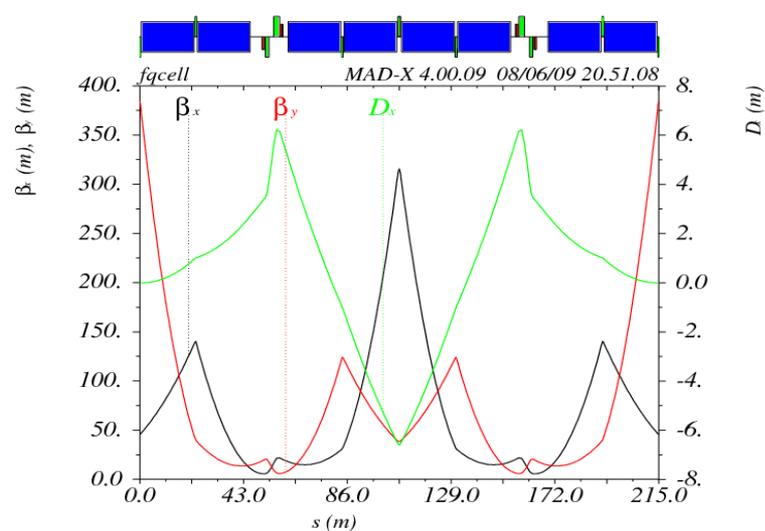
IR dispersion function



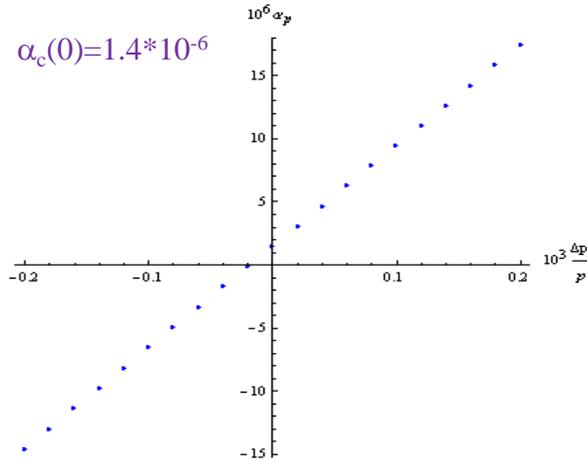
Chromaticity correction



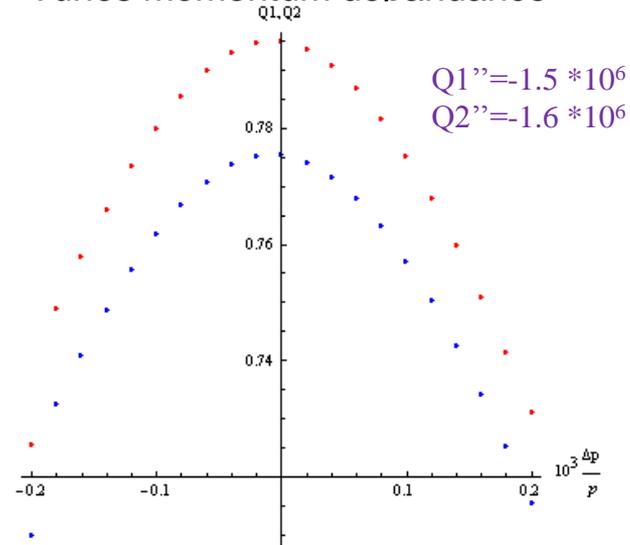
Arc cell



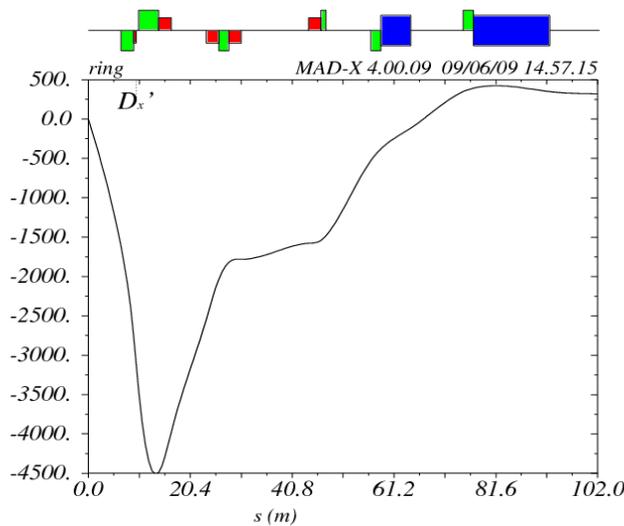
Momentum compaction factor



Tunes momentum dependance



Second order dispersion

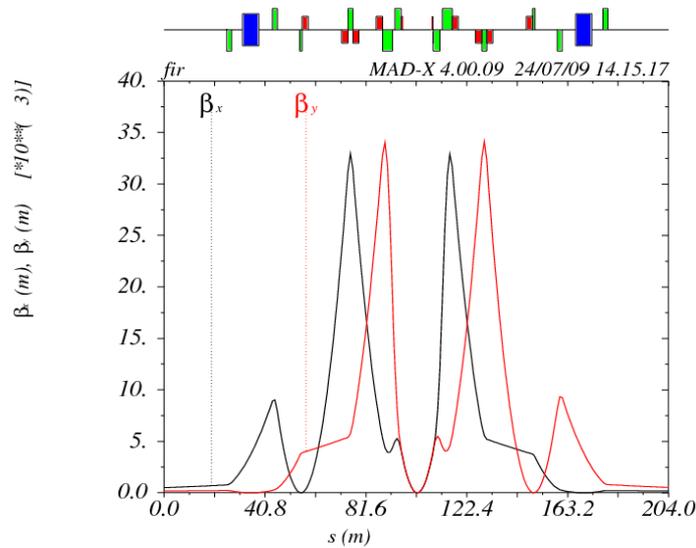


$$\xi_2 = \frac{1}{8\pi} \int_0^C [S_0(s)D_0(s) - K_0] \Delta\beta_1^C(s) ds + \frac{1}{4\pi} \int_0^C \beta_0(s) S_0(s) \Delta D_1^C(s) ds - \xi_1$$

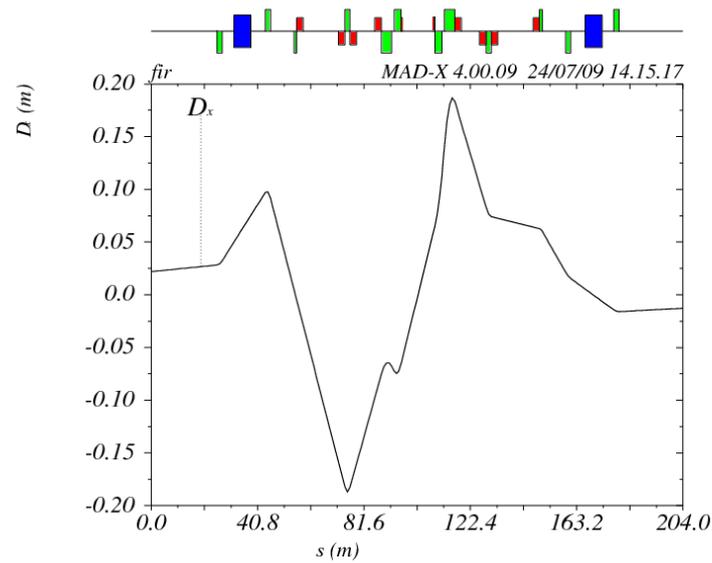
$$\Delta D_1^C(s) = -\sqrt{\beta_0(s)} \int_s^{s+C} \frac{\sqrt{\beta_0(s')}}{\sin(\mu_0/2)} [S_0(s')D_0(s') - K_0(s')] \times D_0(s') \cos\left[\frac{\mu_0}{2} - |\mu(s') - \mu(s)|\right] ds'$$

$$\frac{\Delta\beta_1^C(s)}{\beta_0(s)} = \frac{-1}{2 \sin \mu_0} \int_s^{s+C} [S_0(s')D_0(s') - K_0(s')] \beta_0(s') \times \cos[\mu_0 - 2|\mu(s') - \mu(s)|] ds'$$

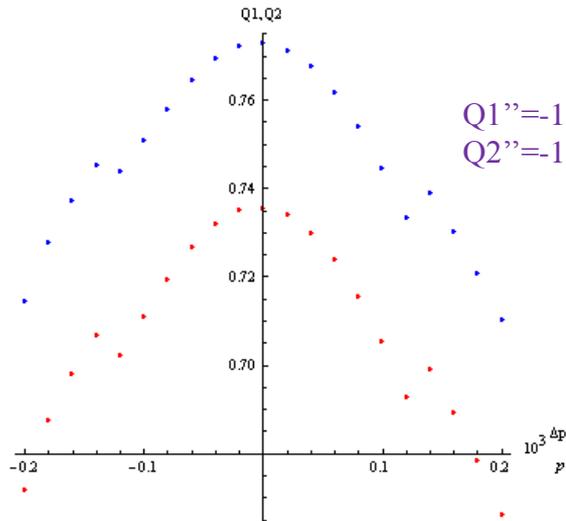
IR beta-functions



Dispersion function



Tunes vs Momentum Deviation



$Q1'' = -1.4 * 10^6$
 $Q2'' = -1.5 * 10^6$

Second order chromaticity calculated using numerical integration

$Q1'' = -1.33 * 10^6$
 $Q2'' = -1.39 * 10^6$

$$\xi_2 = \frac{1}{8\pi} \int_0^C [S_0(s)D_0(s) - K_0] \Delta\beta_1^C(s) ds + \frac{1}{4\pi} \int_0^C \beta_0(s) S_0(s) \Delta D_1^C(s) ds - \xi_1$$

- Chromaticity of the larger β -function should be corrected first (before φ is allowed to change)
 - and in one kick to reduce sensitivity to errors!

- To avoid spherical aberrations it must be $\beta_y \Rightarrow$ then small β_x will kill all detuning coefficients and RDTs (this will not happen if $\beta_y \leftrightarrow \beta_x$)

$$\frac{dQ_y}{dE_y} \sim b_2^2 \beta_x \beta_y^2, \quad \frac{dQ_x}{dE_x} \sim b_2^2 \beta_x^3$$

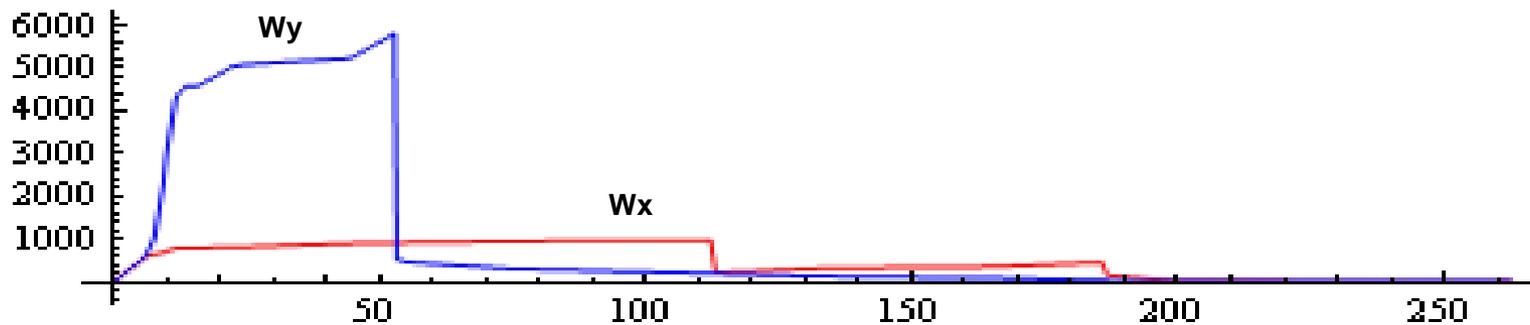
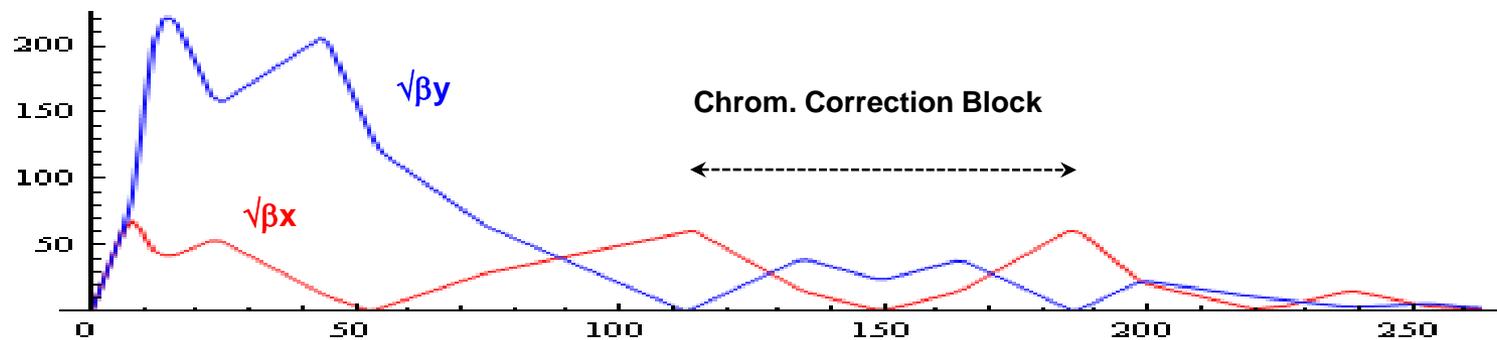
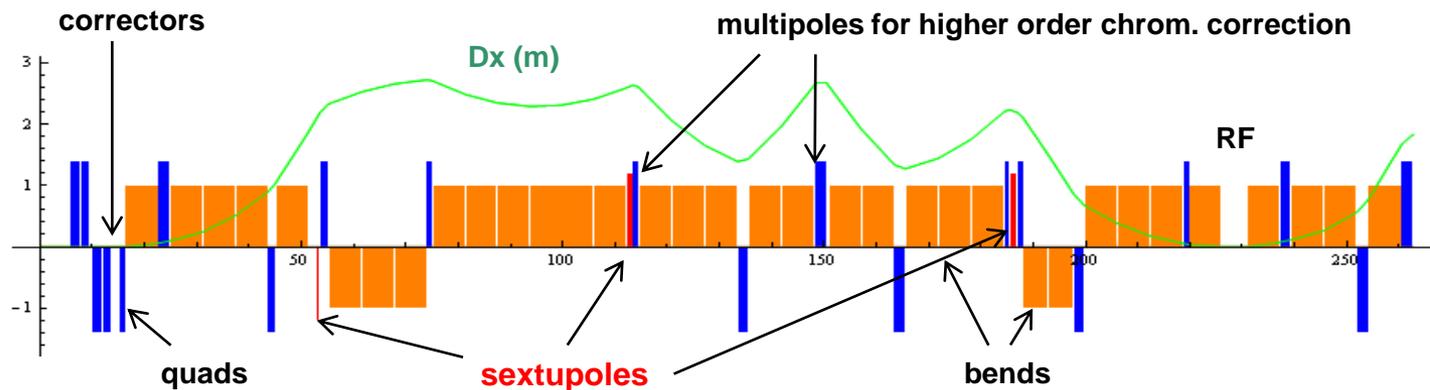
- Chromaticity of β_x should be corrected with a pair of sextupoles separated by -I section to control DDx (smallness of β_y is welcome but not sufficient)
- Placing sextupoles in the focal points of the other β -function separated from IP by $\varphi = \pi \times \text{integer}$ reduces sensitivity to the beam-beam interaction.

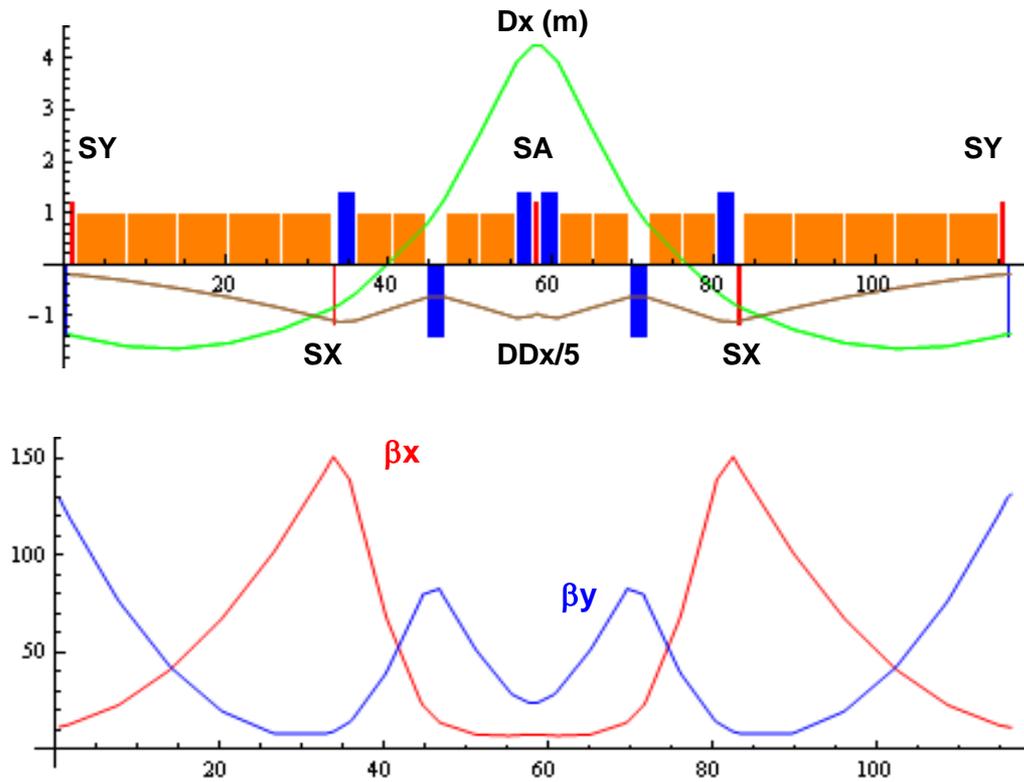
These considerations uniquely determine the IR layout.

Eliana came very close to it, just minor corrections were needed.

Requirements adopted for the latest version:

- full aperture $A = 10\sigma_{\text{max}} + 2\text{cm}$ (A.Zlobin adds 1cm on top of that)
- maximum quad gradient 12% below quench limit at 4.5°K as calculated by A.Zlobin
- bending field 8T in large-aperture open-midplane magnets, 10T in the arcs
- IR quad length < 2m (split in parts if necessary!) – no shielding from inside
- Sufficient space for magnet interconnects (typically 30-40cm)



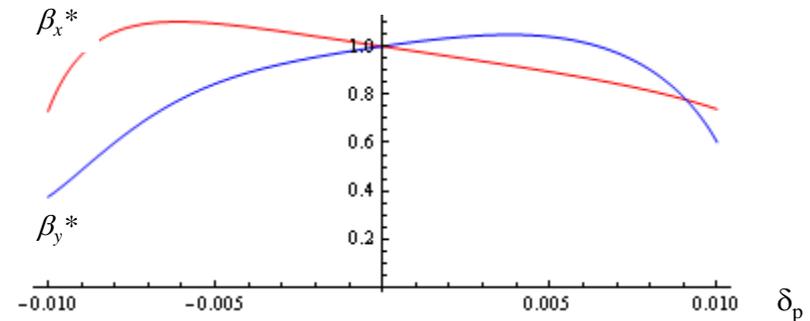
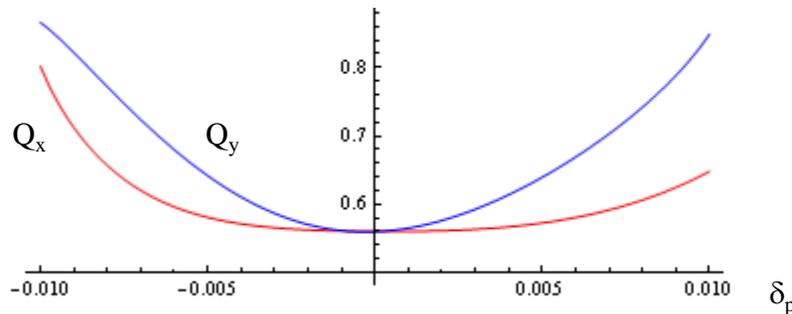


$$\alpha_c = \frac{1}{C} \int_0^C \frac{D_x}{\rho} ds,$$

$$\frac{d\alpha_c}{d\delta_p} = \frac{1}{C} \int_0^C \left[\frac{DD_x}{\rho} + \frac{1}{2} (D'_x)^2 \right] ds$$

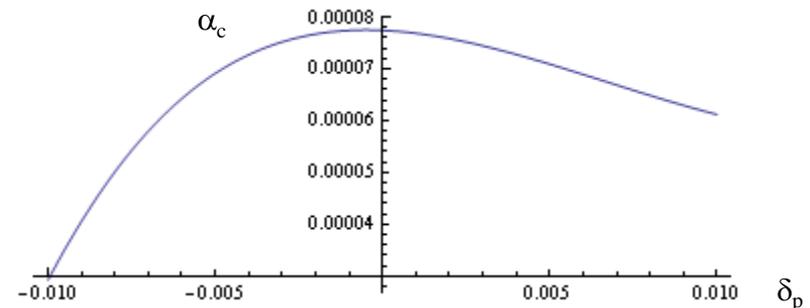
- Central quad and sextupole SA control the momentum compaction factor and its derivative (via Dx and DDx) w/o significant effect on chromaticity
- Large β -functions ratios at SX and SY sextupole locations simplify chromaticity correction
- Phase advance 300° / cell \Rightarrow spherical aberrations cancelled in groups of 6 cells
- Large dipole packing factor \Rightarrow small circumference ($C=2.6$ km with 9.2T dipole field)

Fractional parts of the tunes

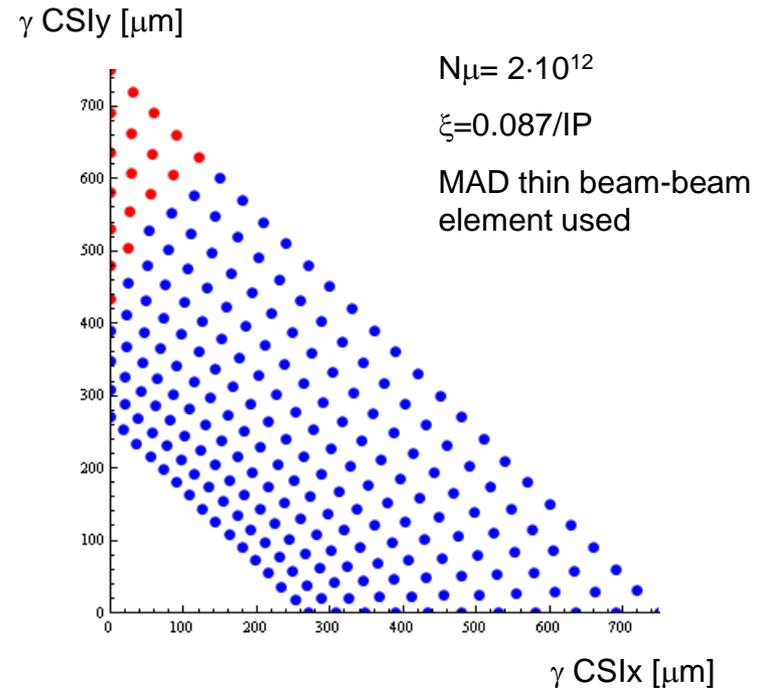
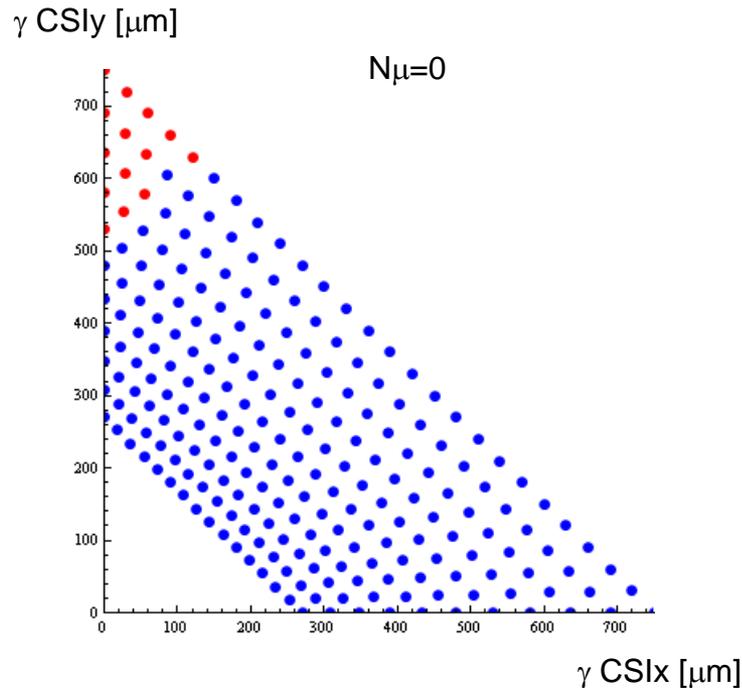


With 2 IPs the central tunes are 18.56, 16.56

- neutral for beam-beam effect
- good for the orbit stability and DA



- Static momentum acceptance is $\pm 1\%$, but the baseline scheme calls for only $\pm 0.3\%$
- The momentum compaction factor can be lowered to $\sim 5 \cdot 10^{-5}$, or made even smaller negative



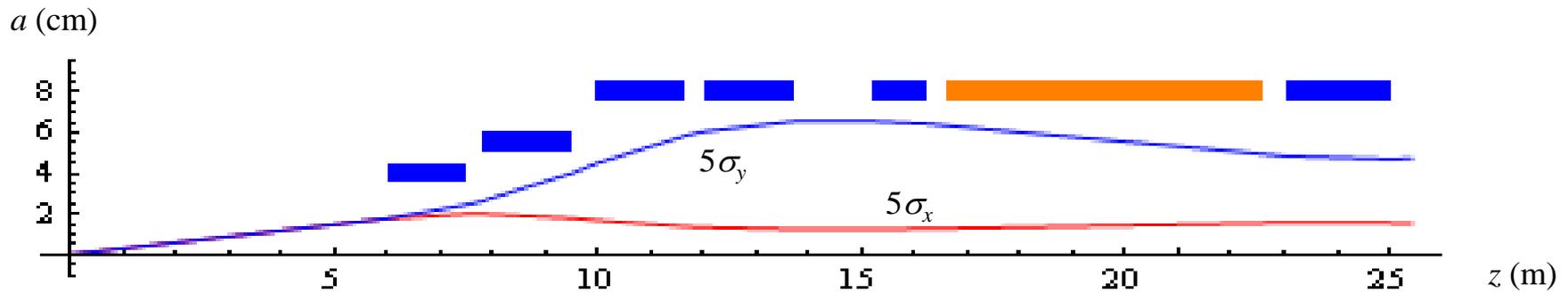
1024 turns DA (Lie3 tracking method): blue – Courant-Snyder Invariants of stable muons, red – lost muons

$$DA = (\gamma \text{ CSI} / \varepsilon_{\perp N})^{1/2} = 4.5\sigma \text{ for } \varepsilon_{\perp N} = 25 \mu\text{m}$$

- Dynamic Aperture is marginally sufficient for $\varepsilon_{\perp N} = 50 \mu\text{m}$
- DA can be further increased with vertical nonlinear correctors
- With chosen tunes, 18.56, 16.57, beam-beam increases β^* from 1cm to 1.27cm – with thin beam-beam element, For a long bunch $\beta^* \rightarrow 0,8\text{cm}$ w/o increase in the FF quads

Requirements adopted for this design:

- full aperture $2A = 10\sigma_{\text{max}} + 2\text{cm}$ (Sasha Zlobin wants + 1cm more)
- maximum tip field in quads = 10T ($G=200\text{T/m}$ for $2A=10\text{cm}$)
- bending field 8T in large-aperture open-midplane magnets, 10T in the arcs
- IR quad length < 2m (split in parts if necessary!)



Gradient (T/m)	250	187	-131	-131	-89	82
Quench @ 4.5°K	282	209	146	146	(with inner radius 5mm larger)	
Quench @ 1.9°K	308	228	160	160		
Margin @ 4.5°K	1.13	1.12	1.12			
Margin @ 1.9°K	1.23	1.22	1.22			

- Is the margin sufficient? If not lower beam energy or increase β^* to allow for smaller aperture
- We don't need 5sigma+ half-aperture, 3sigma+ is enough: can accommodate $\epsilon_{\perp N}=50 \mu\text{m}$!
- **No dipole field from 6 to 16.5m, is it worthwhile to create ~2T by displacing the quads?**

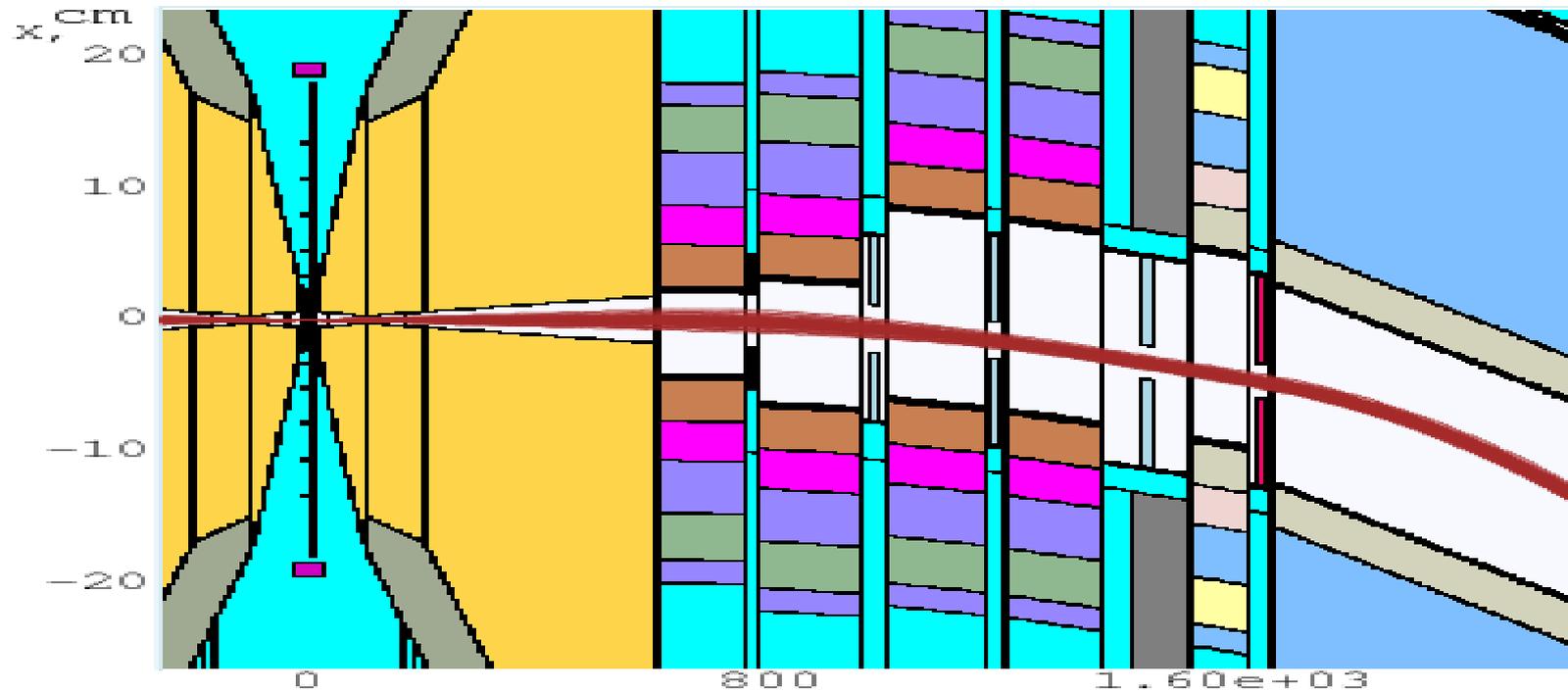
Fears:

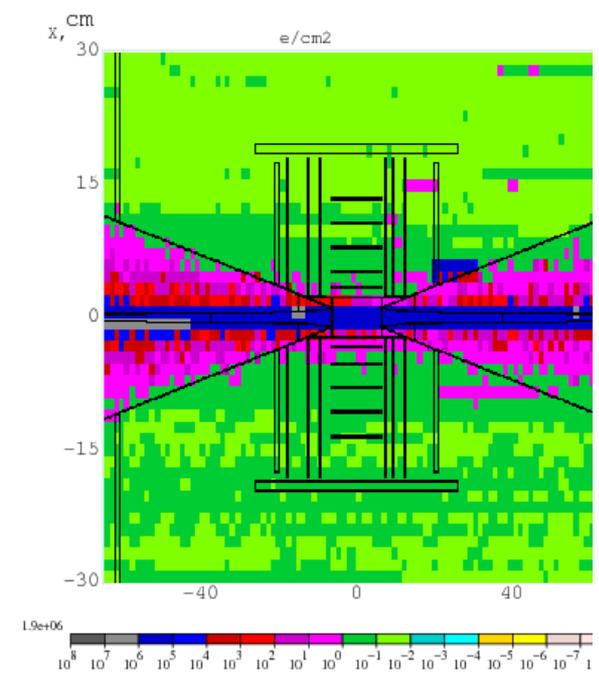
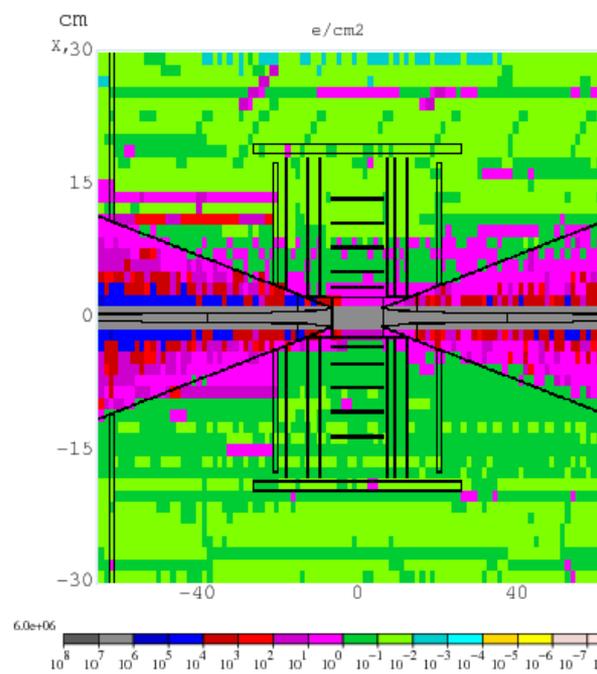
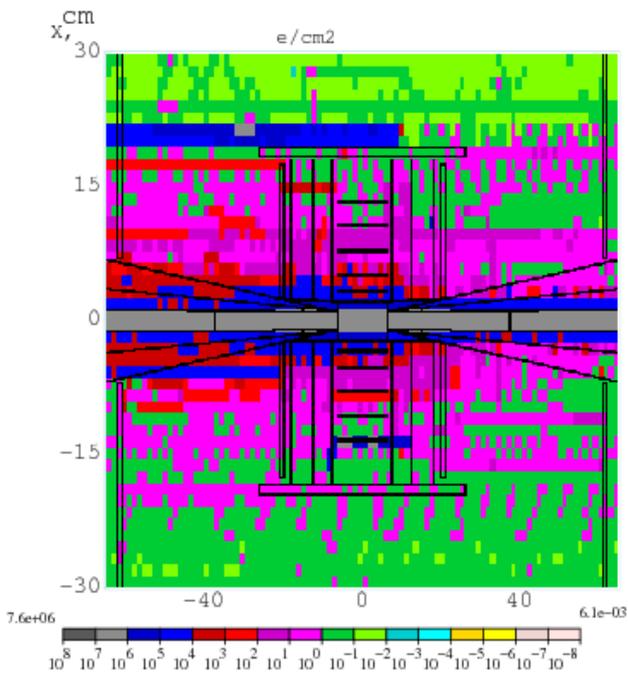
- Dipole field component in the FF quads will deflect the decay electrons more than muons so that they may hit the detector instead of passing through.
- There will be more X-radiation from the decay electrons.

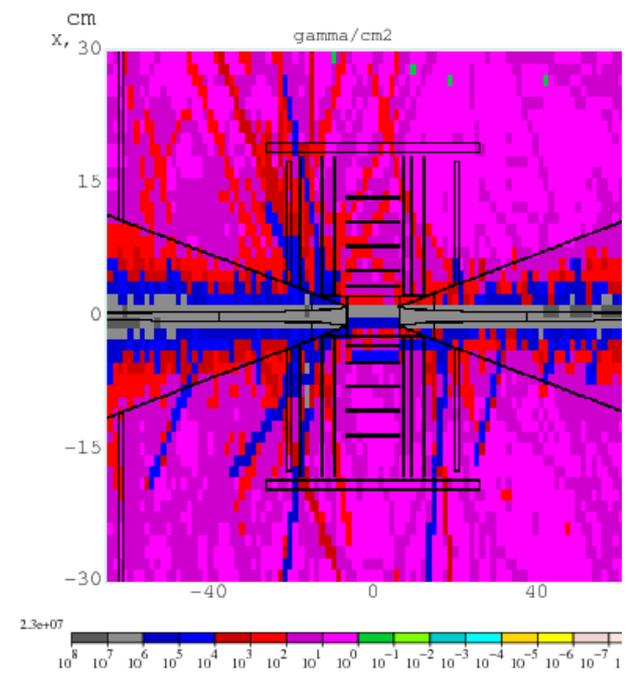
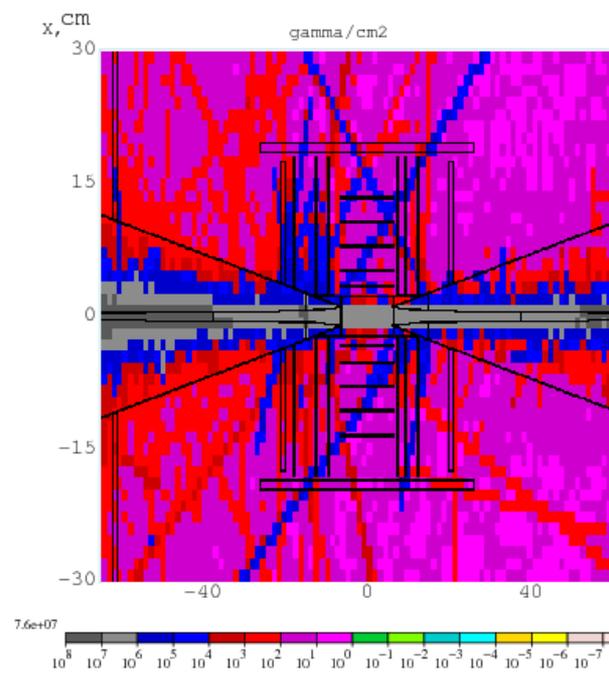
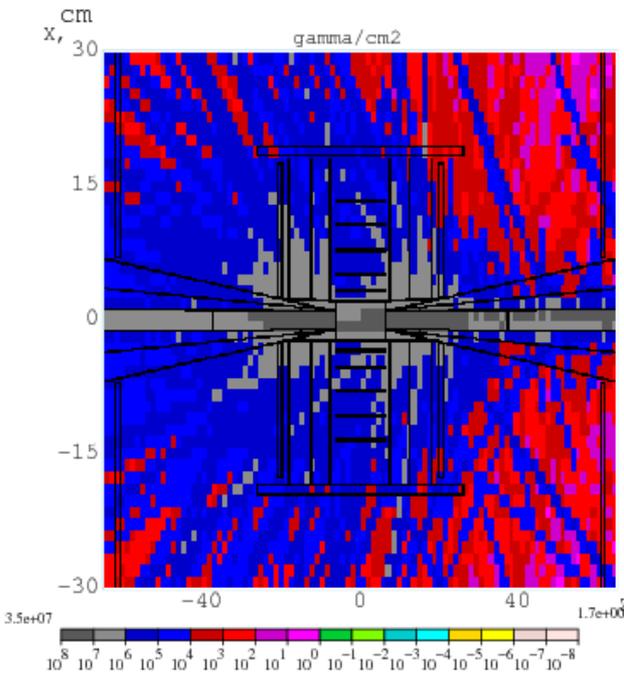
Three cases presented here:

- Initial cone configuration (6° , 5σ inner radius up to 2m from IP), no masks between FF quads - reported at the November workshop at FNAL,
- Cone angle increased to 10° , 5σ inner radius up to 1m from IP, 5σ masks inserted between FF quads
- The same as above + FF quads displacement by 1/10 of the aperture

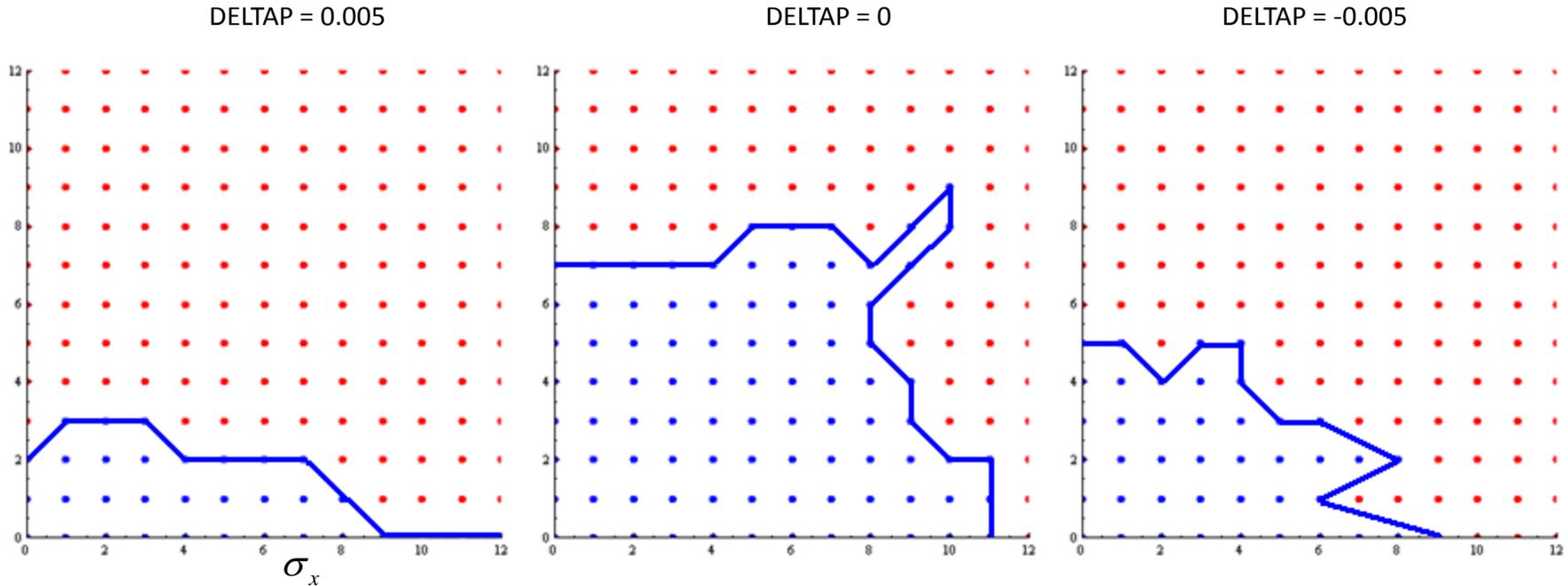
To speed up calculations processes involving neutrons were excluded





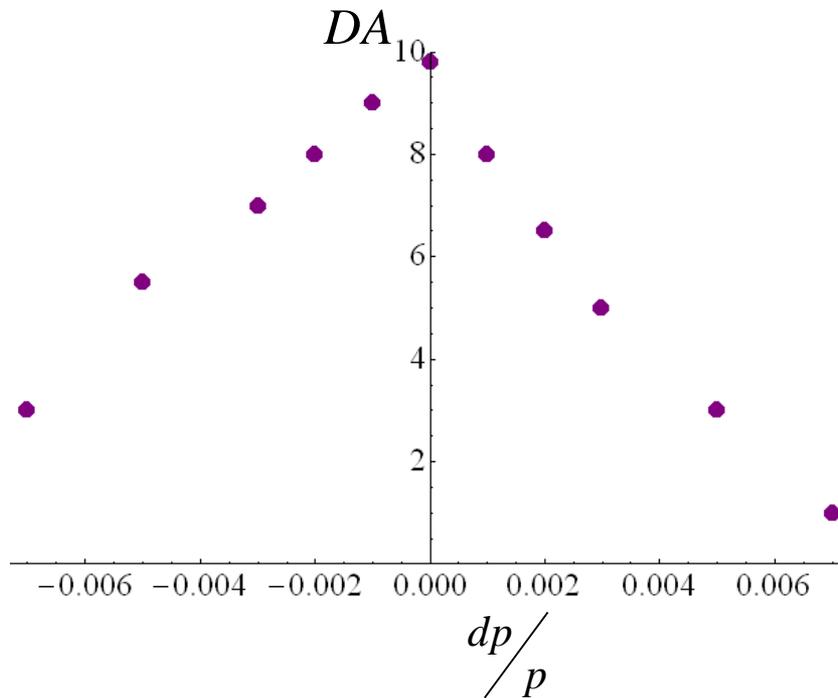


Dynamic Aperture vs Constant Momentum Deviation

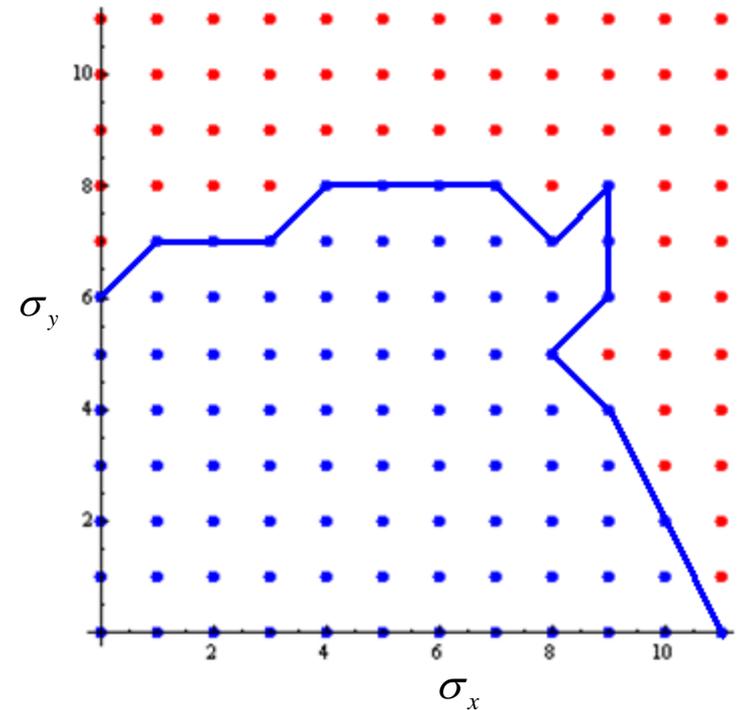


(Calculated using MAD-8 with lie4 method, BeamBeam included, 1024 turns)

$$E_{\perp N} = 10\pi \cdot \text{mm} \cdot \text{mrad} \quad \xi_{BB} = 0.095 \quad \sigma_z = 0.01m \quad \sigma_{x,z} = 5.7 \mu\text{m}$$



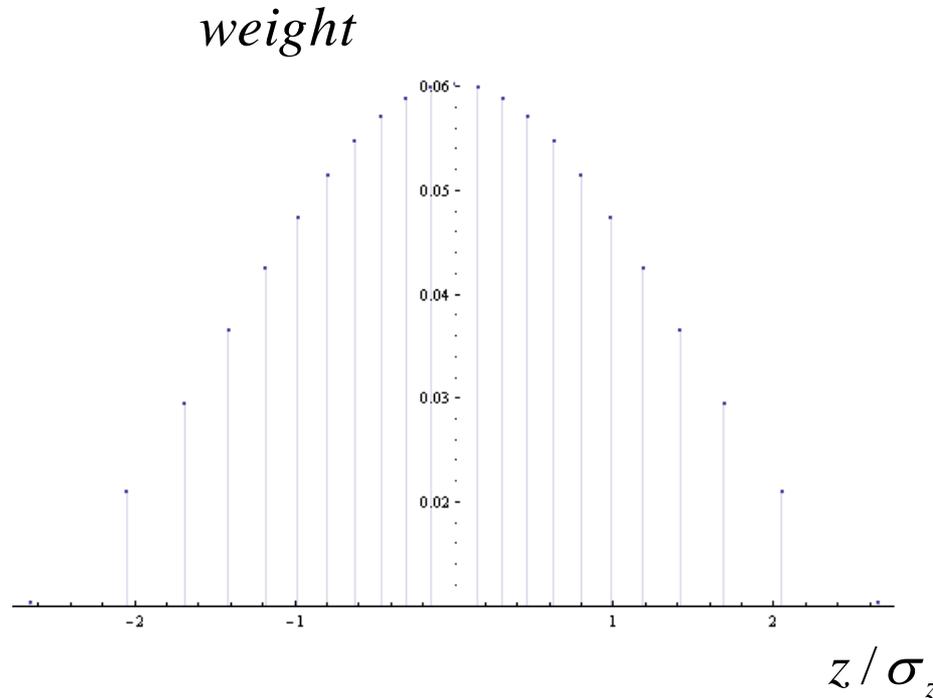
DA diagonal, MAD-8 calculation (4D tracking) for different constant dp/p , BeamBeam included, 1024 turns



MAD-X calculation 6D tracking with synchrotron oscillations, no BeamBeam, 1024 turns

$$f_{RF} = 800 \text{ MHz} \quad V_{RF} = 4.16 \text{ MV}$$

$$Q_s \cong 10^{-3}$$



Number of slices = 23

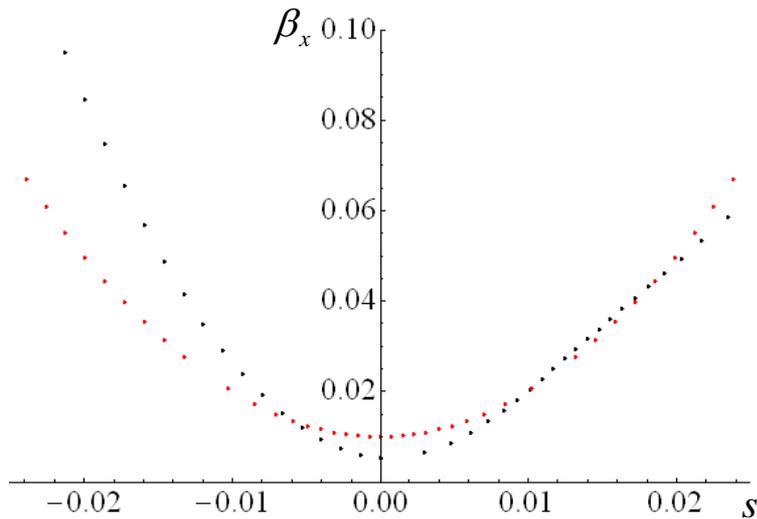
$$\sigma_z = 0.01m$$

$$N_b = 2 \cdot 10^{12}$$

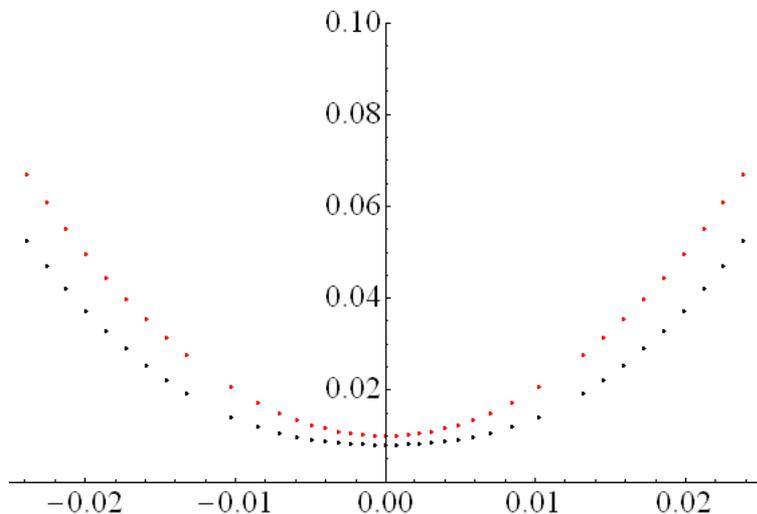
Thin lens model
strength of slice:

$$\frac{1}{f_i} = \frac{-2N \cdot weight_i \cdot r_0}{\gamma \sigma_{x_i} (\sigma_{x_i} + \sigma_{y_i})}$$

In quasi Strong-Strong approach we iteratively calculate new self consistent beta-functions.

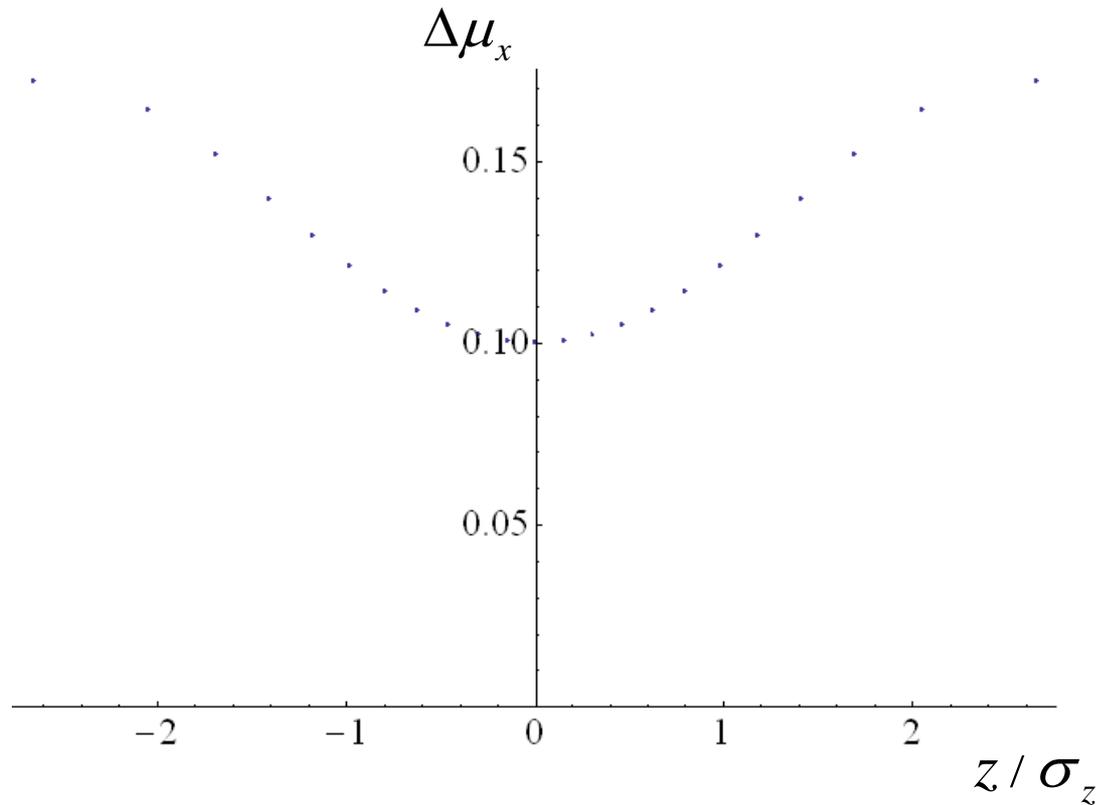


← First slice of the bunch
(black dots)
(red dots represent initial beta-function)



← Middle slice of the bunch
(black dots)

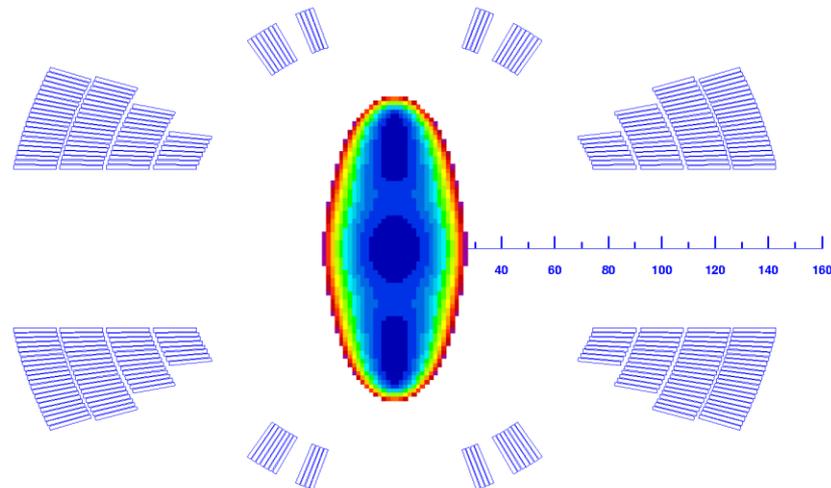
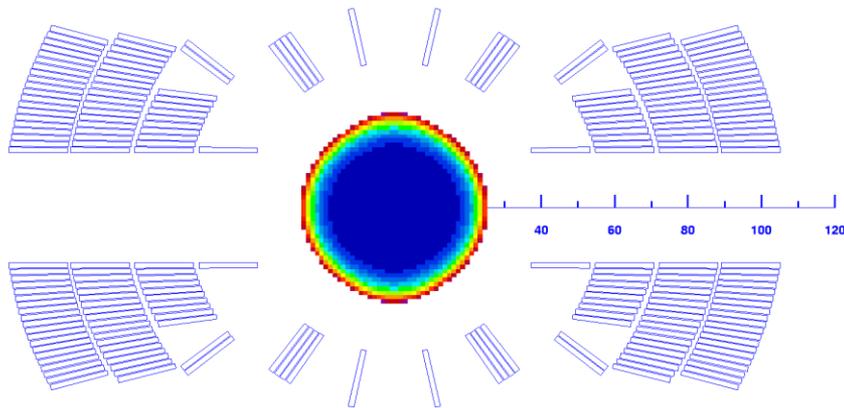
For each slice we introduce a thin lenses at points where it meets the slices of opposite bunch, find new beta-functions, assign them for both bunches and repeat till converged.



$$\Delta\mu_{x,central} = 0.1$$

(as far as we have round beams and almost the same fractional parts of tunes, vertical tune shifts are the same)

[Dynamics of Beam-Beam interaction in Mathematica file](#)



Magnetic field multipole expansion:

$$B_{\theta}(r, \theta) = B_{ref} \times 10^{-4} \sum_{n=1}^{\infty} \left(\frac{r}{r_{ref}} \right)^{n-1} [b_n \cos(n\theta) - a_n \sin(n\theta)]$$

$$B_r(r, \theta) = B_{ref} \times 10^{-4} \sum_{n=1}^{\infty} \left(\frac{r}{r_{ref}} \right)^{n-1} [b_n \sin(n\theta) + a_n \cos(n\theta)]$$

IR dipole:

Rref=40mm

b1=10000

b3=-5.875

b5=-18.320

b7=-17.105

b9=-4.609

b11=0.390

b13=0.103

Ring dipole:

Rref=20mm

b1=10000

b3=0.003

b5=-0.012

b7=0.154

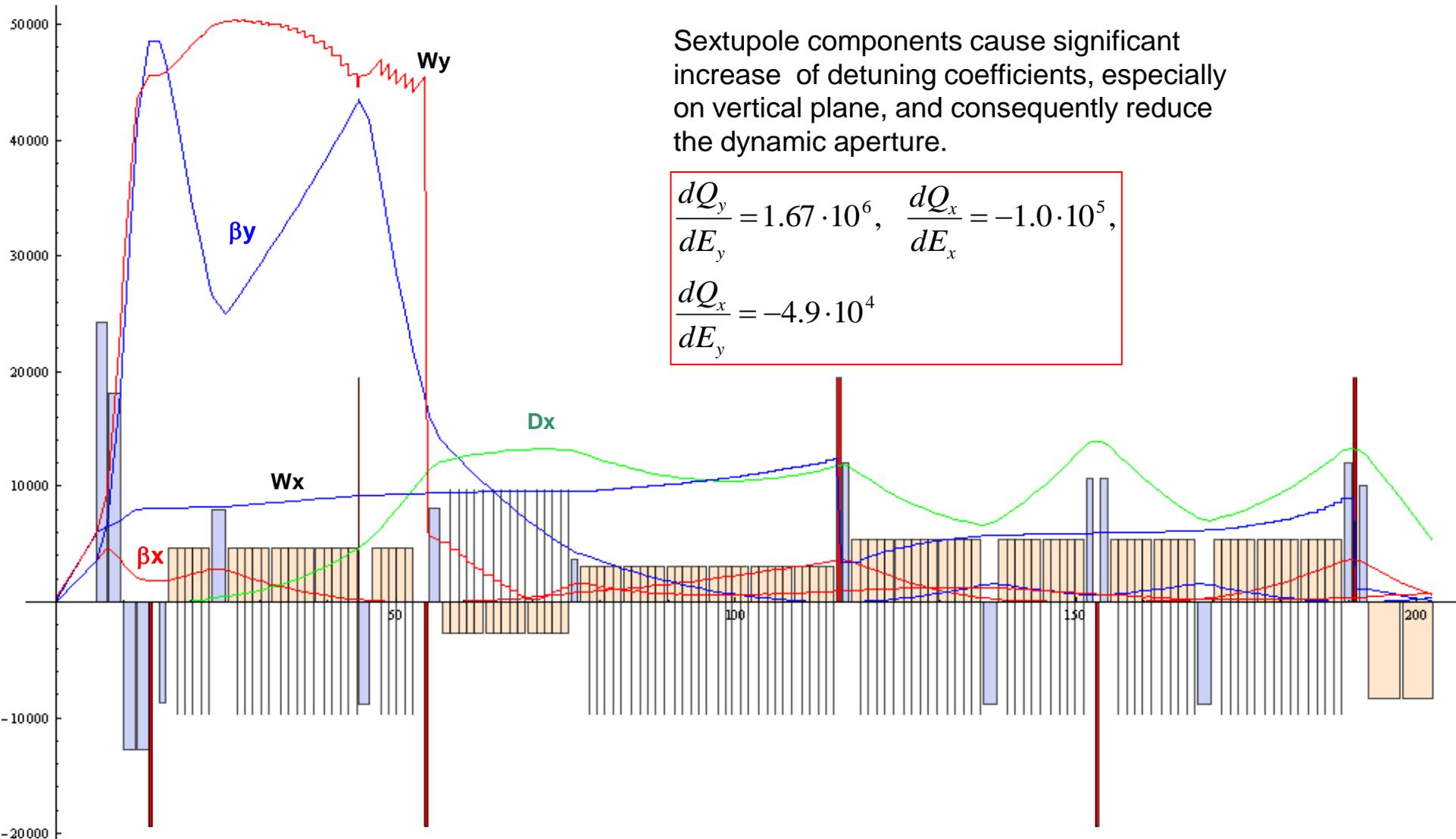
b9=-1.185

b11=-0.118

b13=0.053

(V.V.Kashikhin)

Magnets sliced in 5 pieces and thin multipoles introduced between them.



Sextupole components cause significant increase of detuning coefficients, especially on vertical plane, and consequently reduce the dynamic aperture.

$$\frac{dQ_y}{dE_y} = 1.67 \cdot 10^6, \quad \frac{dQ_x}{dE_x} = -1.0 \cdot 10^5,$$

$$\frac{dQ_x}{dE_y} = -4.9 \cdot 10^4$$

Lattice update including:

- FF quads displacement to produce dipole field $\sim 2\text{T}$
- Momentum compaction factor correction (making it slightly negative)
- Arc magnet length reduction to get 10T bending field \rightarrow even smaller circumference!

Next steps:

- Study the effect of magnet imperfections with Vadim Kashikhin's magnet design
- Possibility to change β^* in a wide range w/o changing the layout
- Collimation system design (are special sections necessary?)
- Study effects of fringe fields, add nonlinear correctors if necessary
- Design orbit correction and tuning circuits
- Study the effect of random misalignments and magnet imperfections

\sqrt{s} (TeV)	1.5	3
Av. Luminosity / IP ($10^{34}/\text{cm}^2/\text{s}$)	1.2*	5
Max. bending field (T)	9.2**	14
Av. bending field in arcs (T)	7.7	12
Circumference (km)	2.6	4
No. of IPs	2	2
Repetition Rate (Hz)	15	12
Beam-beam parameter / IP	0.087	0.087
β^* (cm)	1	0.5
Bunch length (cm)	1	0.5
No. bunches / beam	1	1
No. muons/bunch (10^{12})	2	2
Norm. Trans. Emit. (μm)	25	25
Energy spread (%)	0.1	0.1
Norm. long. Emit. (m)	0.07	0.07
Total RF voltage (MV) at 800MHz	60	700
μ^+ in collision / 8GeV proton	0.008	0.007
8 GeV proton beam power (MW)	4.8	4.3

*) With **increase** by the beam-beam effect

***) Not 10T just by mistake

$$\langle \mathcal{L} \rangle = f_0 \frac{n_b N_\mu^2}{4\pi\epsilon_\perp \beta^*} h \times \frac{1}{2} \mathcal{F}_{rep} \sim \frac{P_\mu \xi}{C\beta^*} h\tau$$

P_μ – average muon beam power ($\sim \gamma$)

$$\xi = \frac{r_\mu N_\mu}{4\pi\gamma\epsilon_\perp} \quad \text{– beam-beam parameter}$$

C – collider circumference ($\sim \gamma$ if $B=\text{const}$)

τ – muon lifetime ($\sim \gamma$)

β^* – beta-function at IP

