

Development of laser stripping at the SNS



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On behalf of the SNS laser stripping team

October 28, 2010

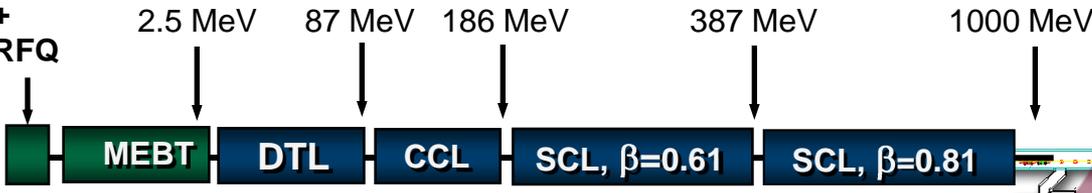
Outline

- **History of laser stripping**
- **Proof of principle experiment**
- **Physical model**
- **Proof of practicality experiment**
- **Experimental studies**
- **Future plans and ideas for SNS and other projects**

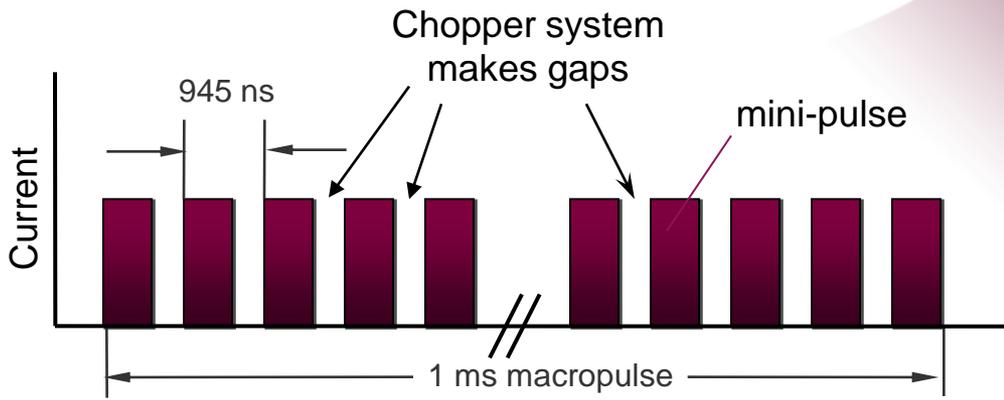
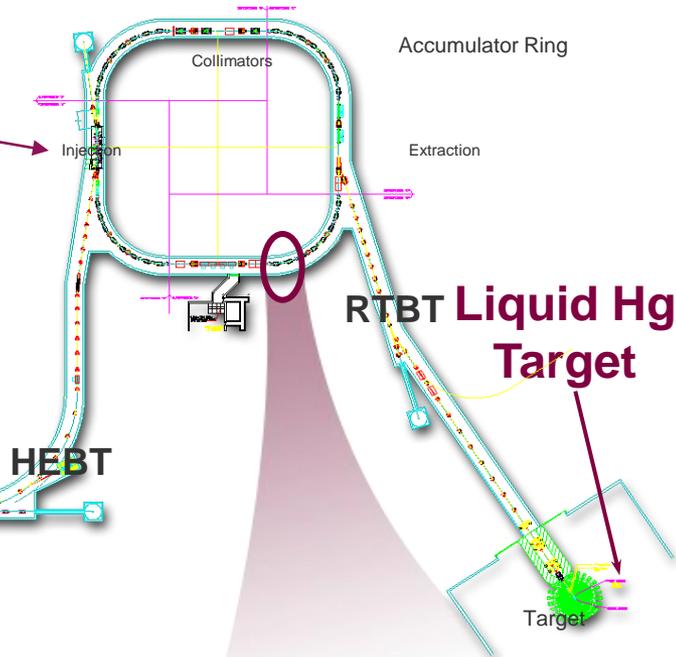
SNS Accelerator Complex

1 GeV
LINAC

Ion Source
+ RFQ

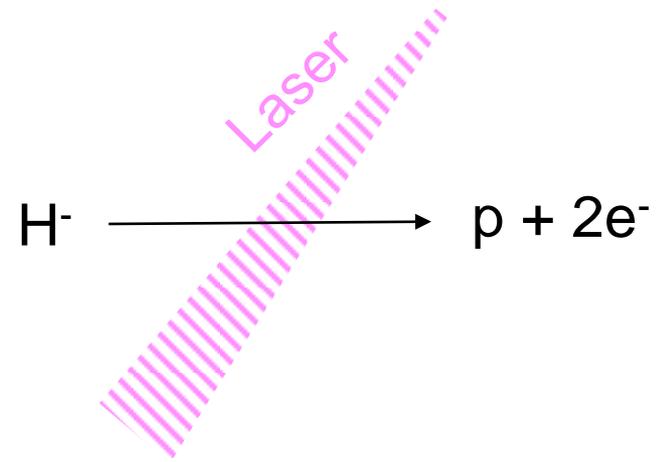
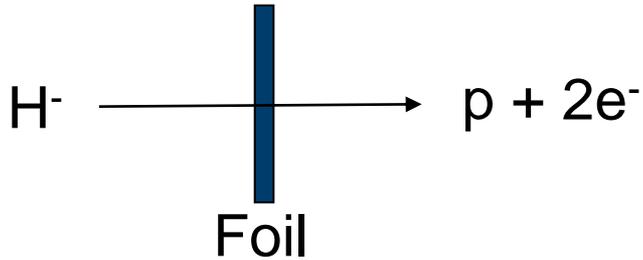


Injection point

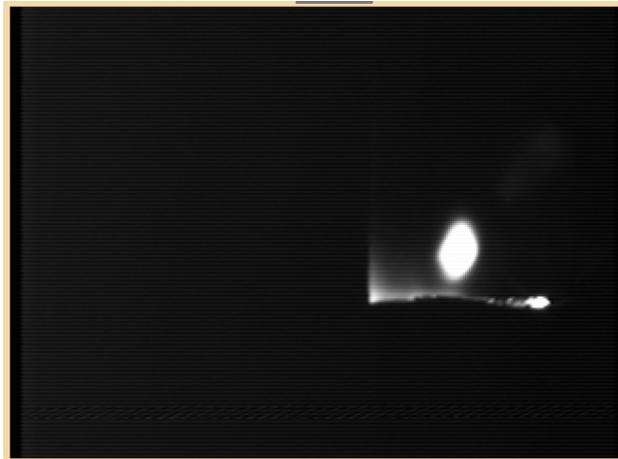


402.5 MHz
micro-pulses

Goals



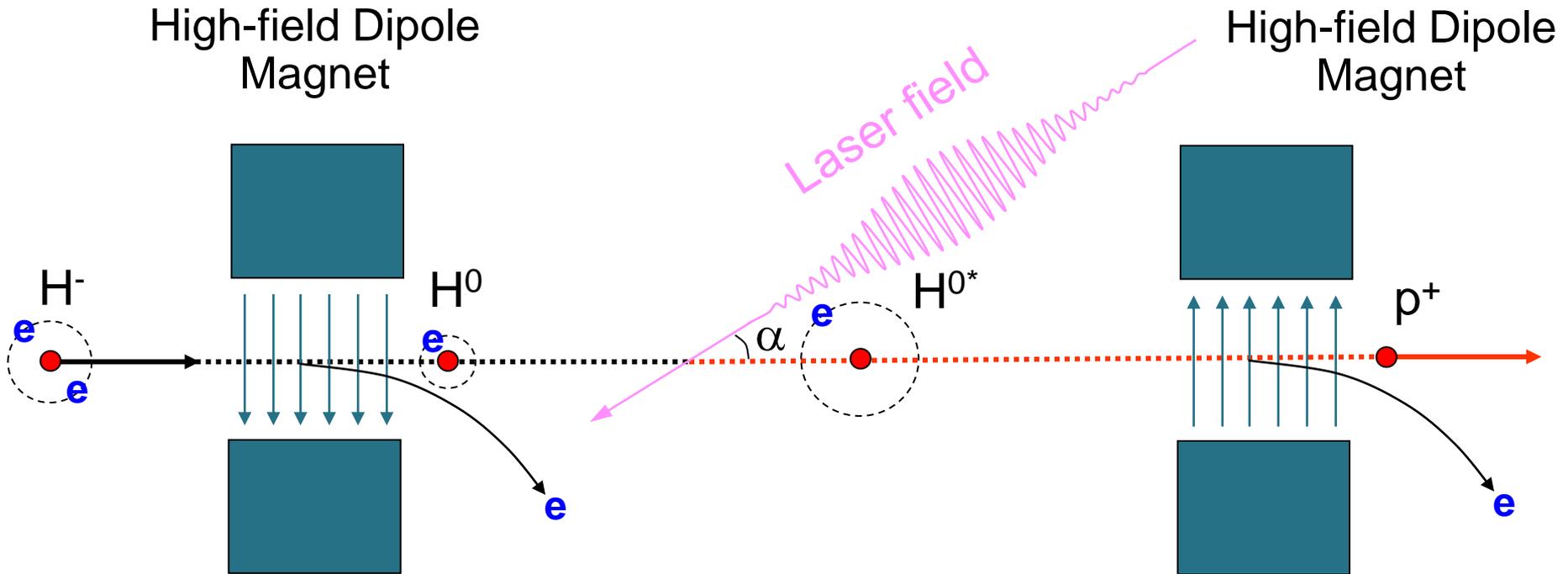
- Replacement of stripping foil by laser assistant stripping



← Foil (Beam on)

Three-step laser stripping scheme

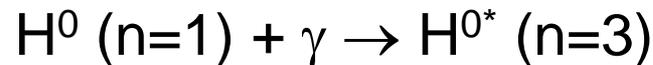
V. Danilov (2003)



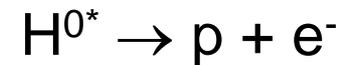
Step 1: Lorentz Stripping



Step 2: Resonant laser Excitation

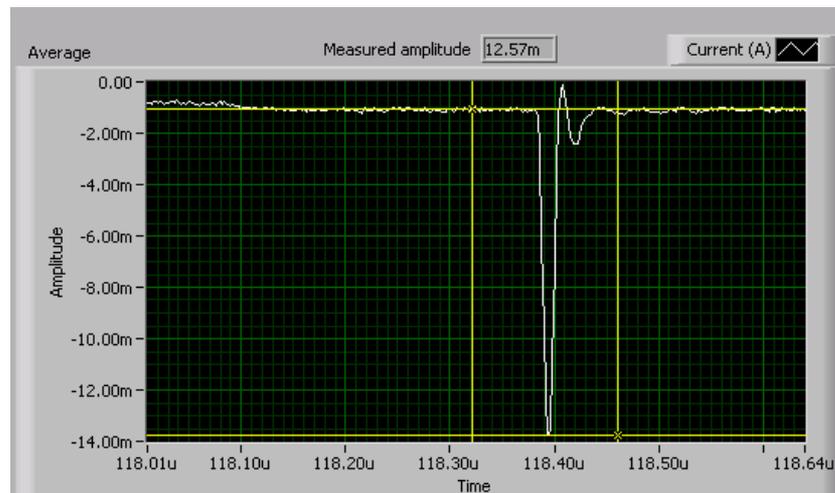
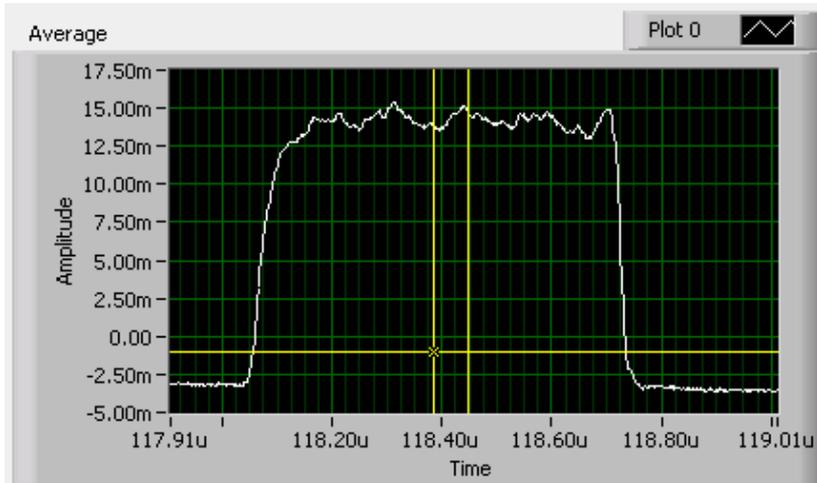


Step 3: Lorentz Stripping



Proof of principle experiment (2006)

mini-pulse



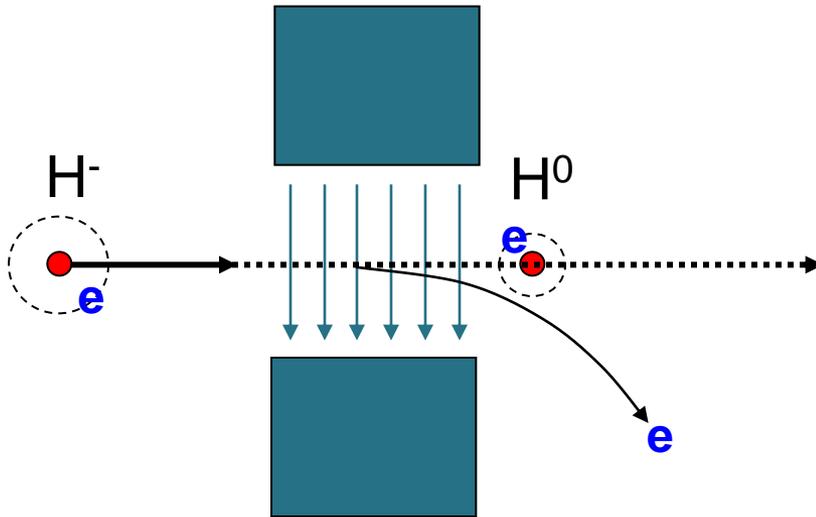
The maximal achieved efficiency: 0.85 ± 0.1 (1st run) and 0.9 ± 0.05 (2nd run)
Stripping time: about 10 ns

Theory of the Laser Stripping

Step 1: Lorentz Stripping



High-field Dipole Magnet



7.1.7 Lorentz Stripping of H^- Ions [1]

M.A. Furman, LBNL

When a H^- ion moves in a magnetic field B it experiences a Lorentz force that bends its trajectory and also tends to break it up since the protons and electrons are bent in opposite directions, and the binding energy of the extra electron is only 0.755 eV. The breakup is a probabilistic process and quantum-mechanical in nature. In the ion rest frame, the stripping force is effected by the electric field E that is the Lorentz-transform of the magnetic field B in the lab, $E = \kappa' \beta \gamma B$, where $\kappa' \simeq 0.3 \text{ GV/T}\cdot\text{m}$. For the H^- ion, $E[\text{MV/cm}] = 3.197 p[\text{GeV}/c] B[\text{T}]$.

The lifetime of the ion in an electric field can be calculated by applying the WKB approximation to the tunneling probability [2]-[4]. It has also been measured in several experiments [5]-[8] whose results, for the ion's lifetime τ in its own rest frame is well parametrized as

$$\tau = \frac{A}{E} \exp\left(\frac{C}{E}\right) \quad (1)$$

In the region of values of E where they overlap, the measurements in [3, 7, 8] are fairly consistent with each other, but are not consistent with [6]. Ref.[7], which covers the range $E = 1.87 - 2.14 \text{ MV/cm}$, has $A = 7.96 \times 10^{-14} \text{ s MV/cm}$ and $C = 42.56 \text{ MV/cm}$, while Ref.[8], which covers $E = 1.87 - 7.02 \text{ MV/cm}$, has $A = (2.47 \pm 0.09) \times 10^{-14} \text{ s MV/cm}$ and $C = 44.94 \pm 0.10 \text{ MV/cm}$. The mean decay length in the lab is given by

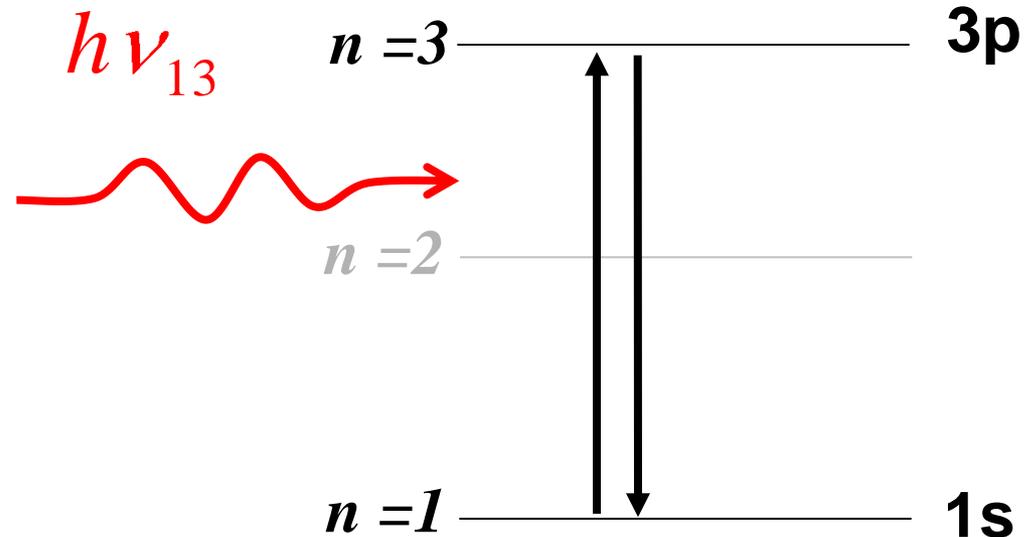
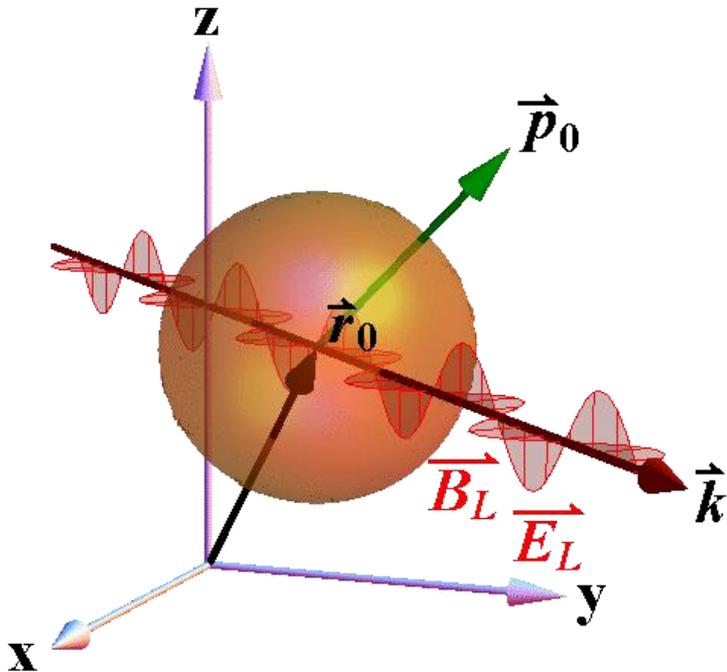
$$\lambda = c\beta\gamma\tau \quad (2)$$

Two level atom (Schrodinger equation)

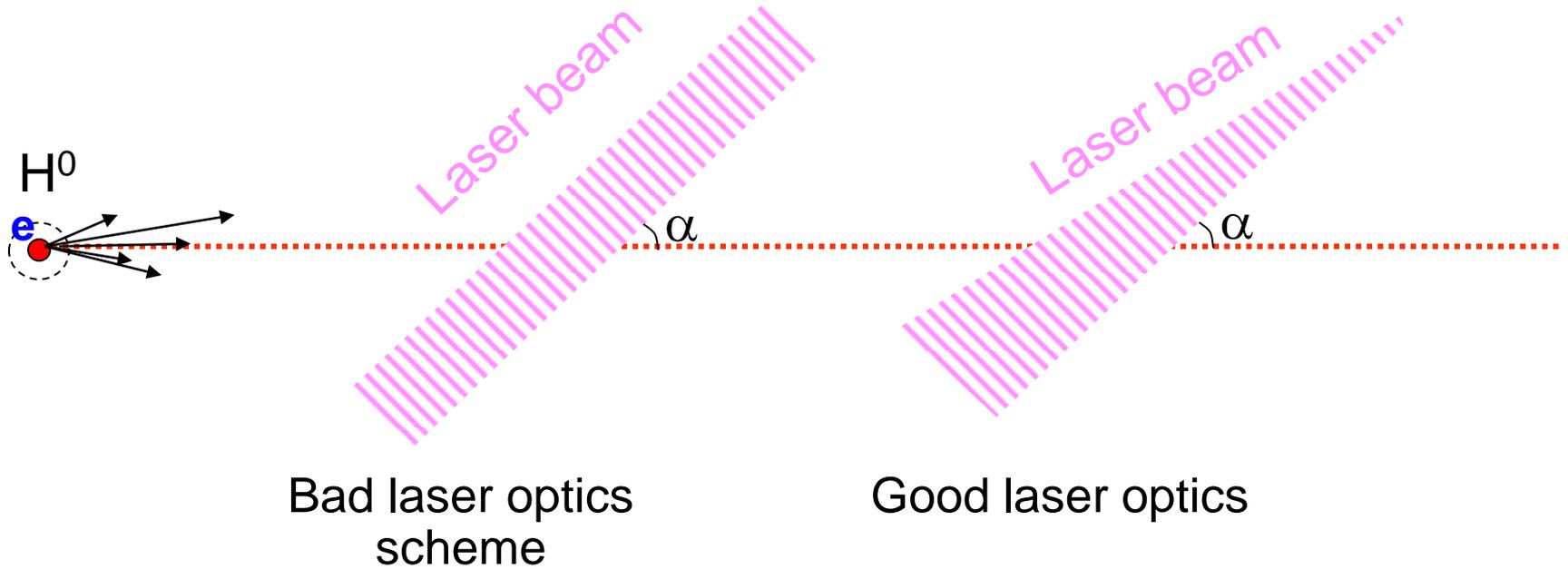
The model is valid only for $E=B=0$ and linear polarized laser field

$$i\hbar \frac{\partial \Psi}{\partial t} = \hat{H}(\vec{E}_L) \Psi$$

$$\begin{cases} \dot{c}_1 = c_3 s_{13}(E_L) \\ \dot{c}_3 = c_1 s_{31}(E_L) \end{cases}$$

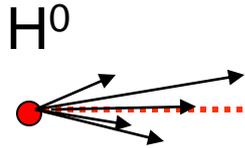


First way to compensate energy spread

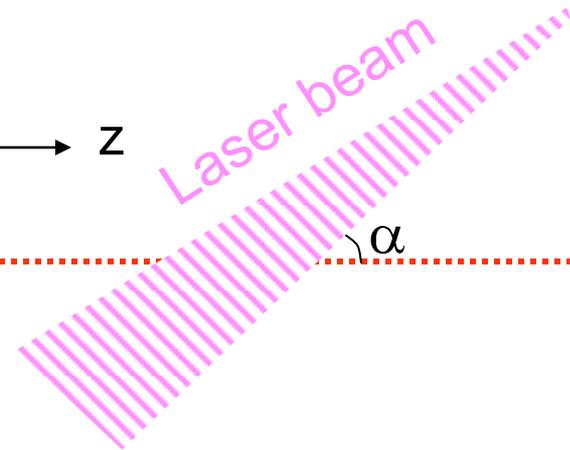
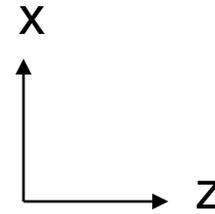
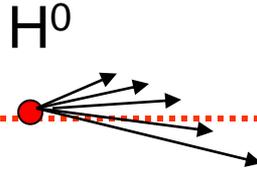


Second way to compensate energy spread

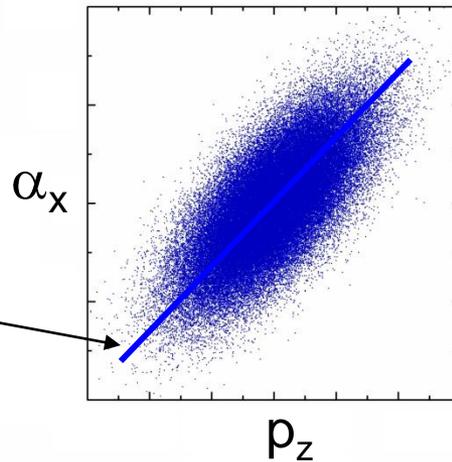
No angle-energy correlation



angle-energy correlation is not zero



Perfect correlation



Necessary correlation is given by dispersion function

$$D_x' = 2.6 \text{ for } T = 1 \text{ GeV beam}$$

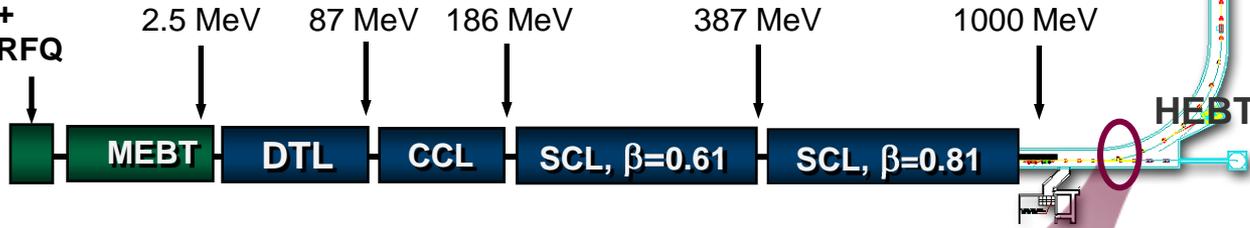
Proof of plasticity experiment

SNS Accelerator Complex

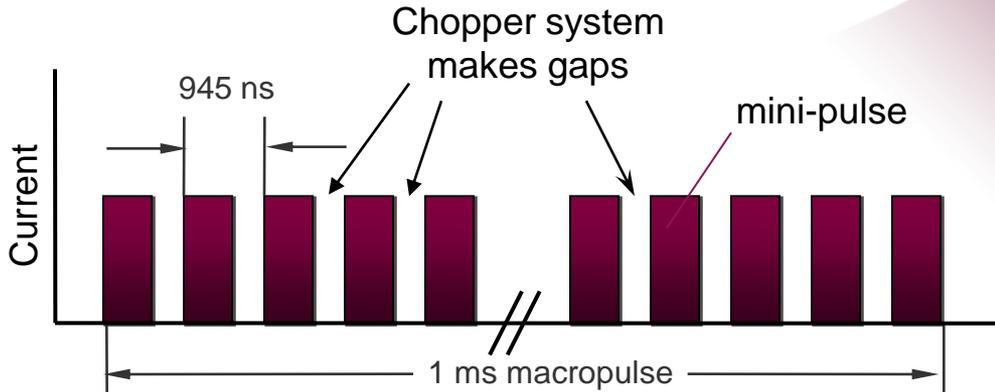
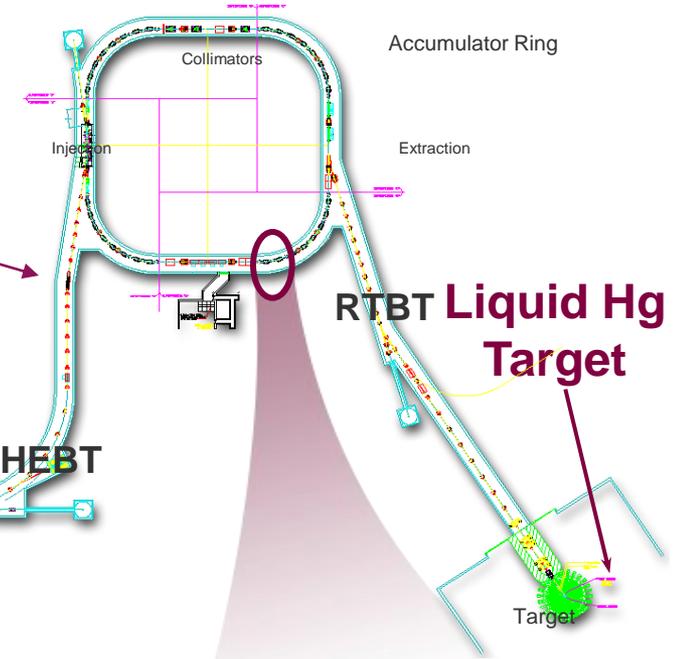
1 GeV
LINAC

Ion Source

+ RFQ



Location of the experiment

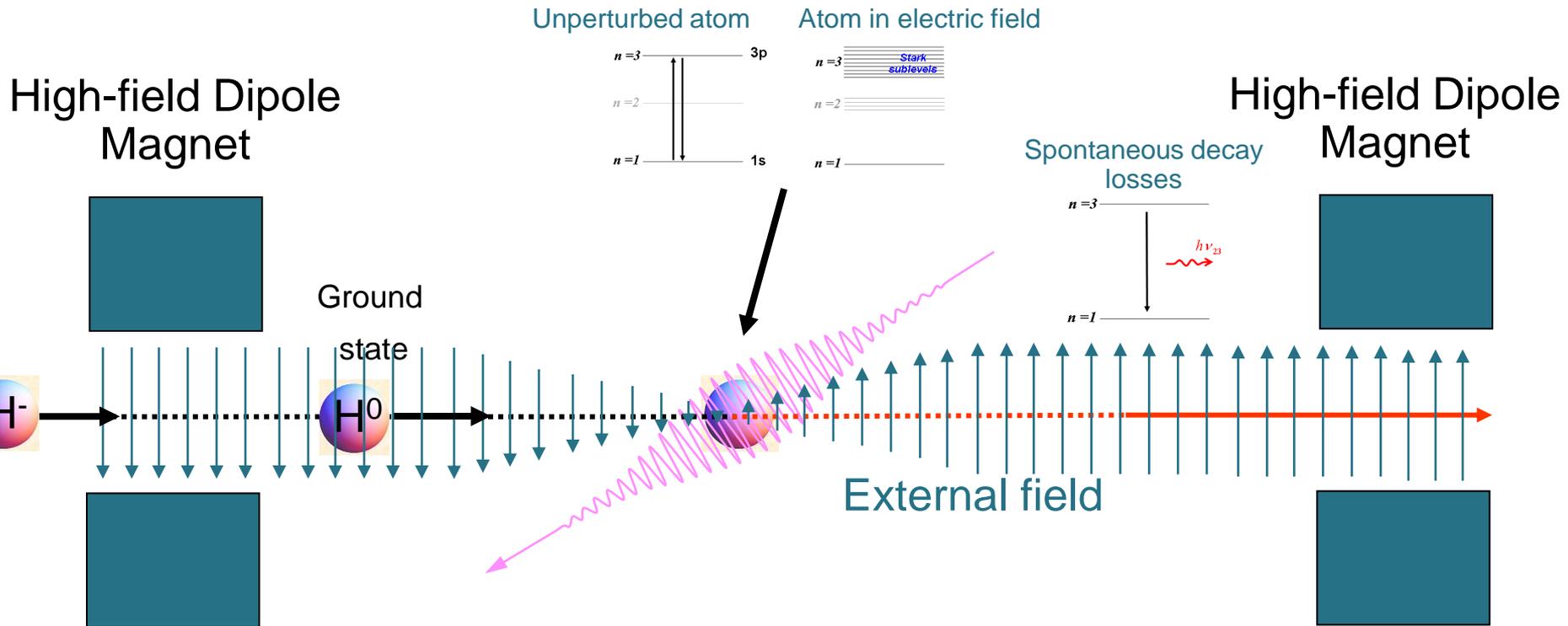


402.5 MHz micro-pulses

Proof of principles
Stripping time appr. 10 ns

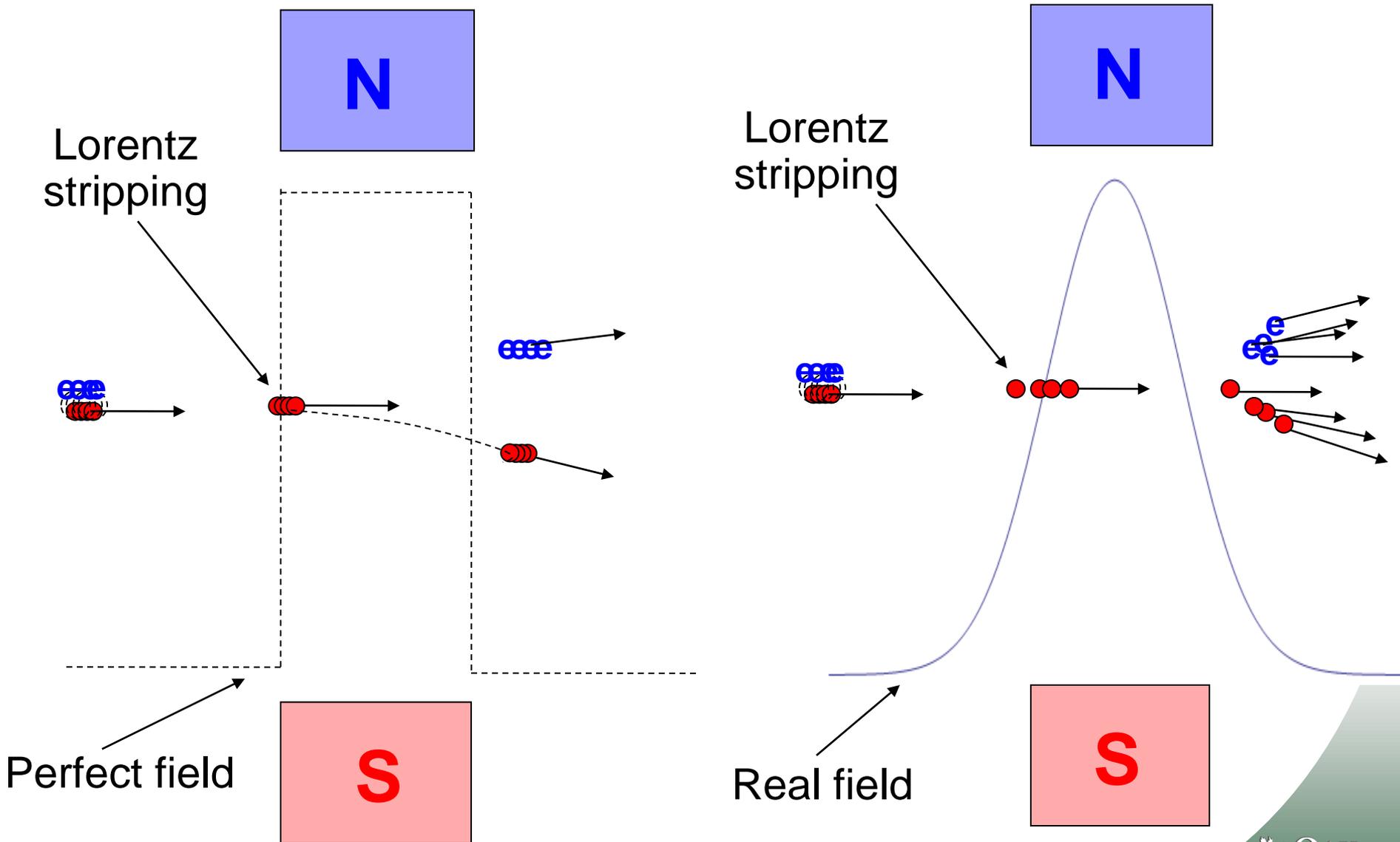
Proof of practicality
Stripping time appr. 1 us

Advanced physical model



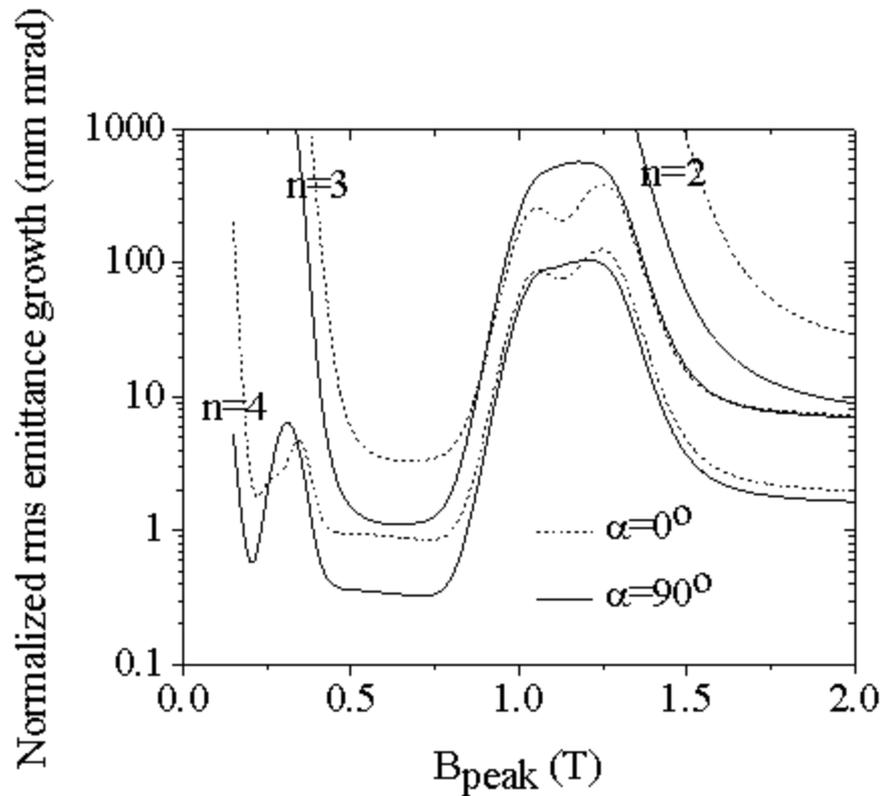
- Stark effect at the interaction point
- Spontaneous decay losses
- Lorentz stripping of H^0 excited beam

Lorentz stripping of H^0 excited beam in the fringe magnetic field



Emittance growth at the fringe field

Fringe field:
$$F(s) = \frac{B}{1 + \exp[a_1 + a_2(\frac{s}{D}) + a_3(\frac{s}{D})^2 + a_4(\frac{s}{D})^3 + a_5(\frac{s}{D})^4 + a_6(\frac{s}{D})^5]}$$



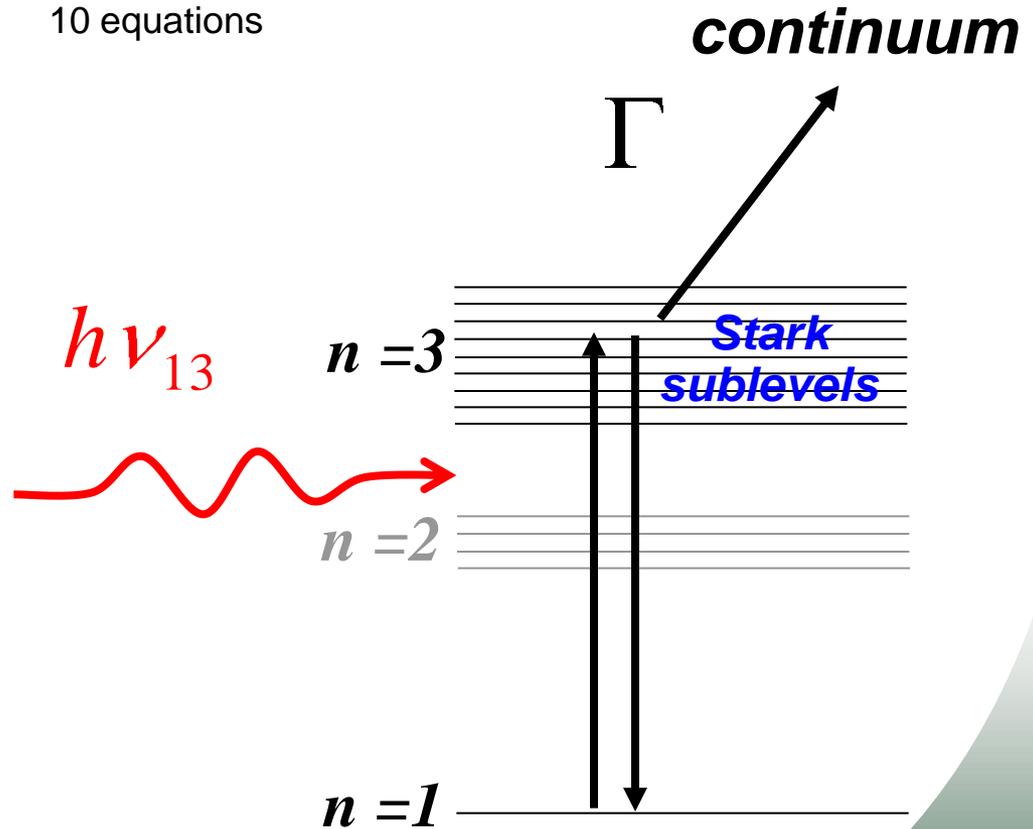
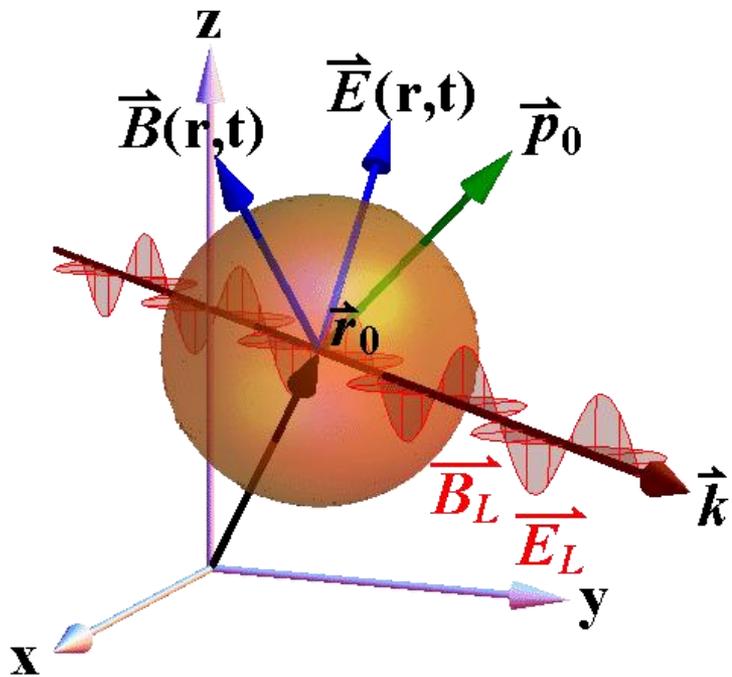
T=4 GeV

Schrödinger equation approach

Spontaneous decay can not be taken into account

$$i\hbar \frac{\partial \Psi}{\partial t} = \hat{H}(\vec{E}, \vec{E}_L) \Psi \quad \rightarrow \quad \dot{c}_m(t) = \sum_{n=1}^N c_n(t) S_{mn}(t) \quad N = 1^2 + 3^2 = 10$$

10 equations

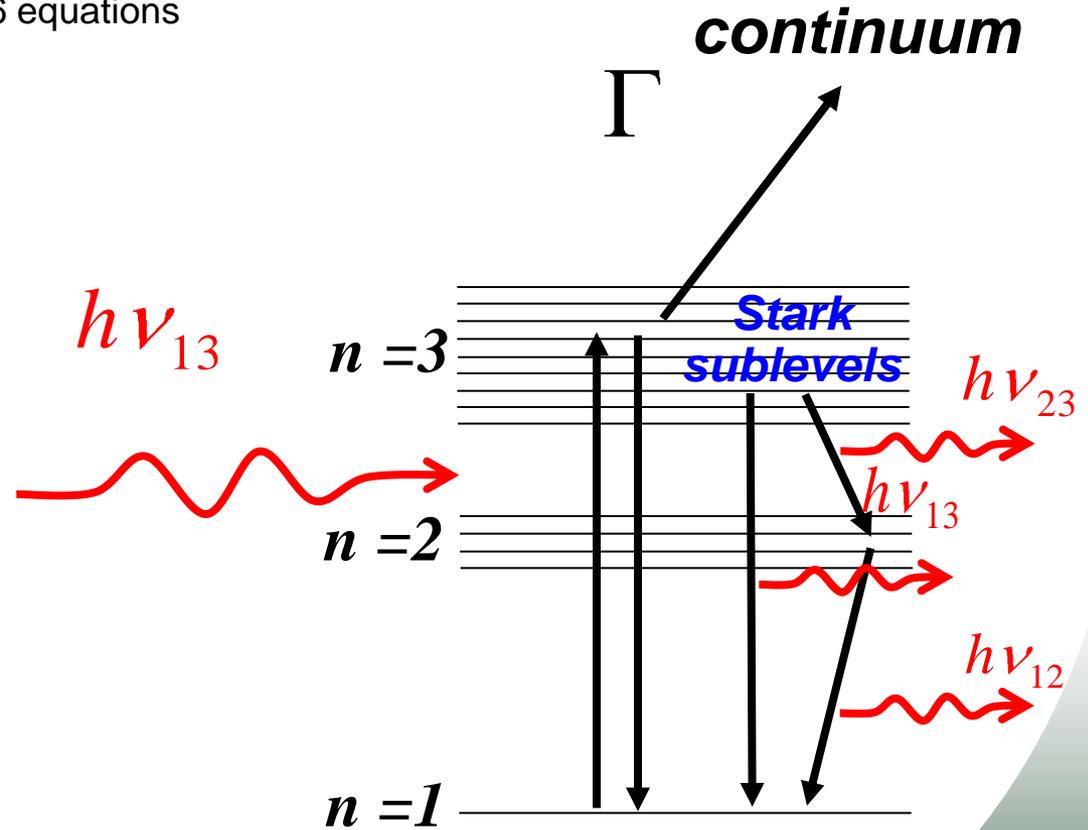
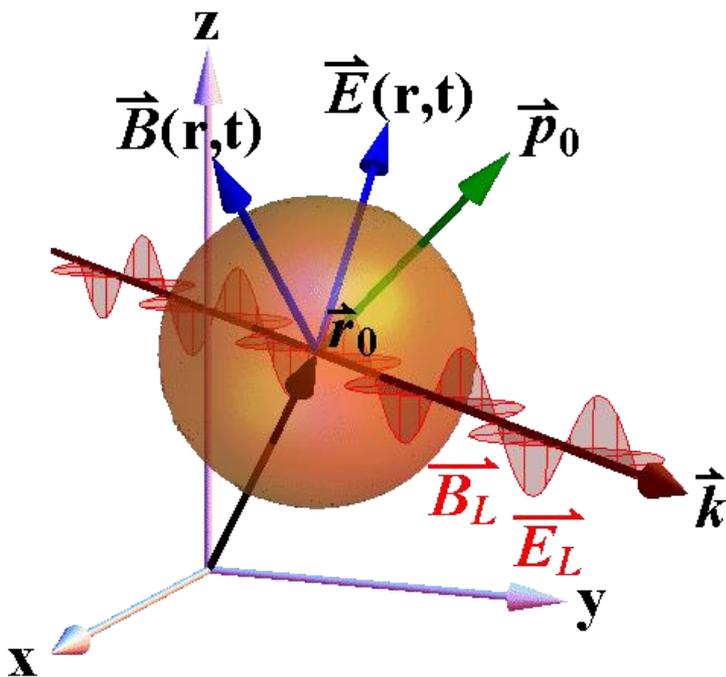


Density matrix approach

Everything can be taken into account

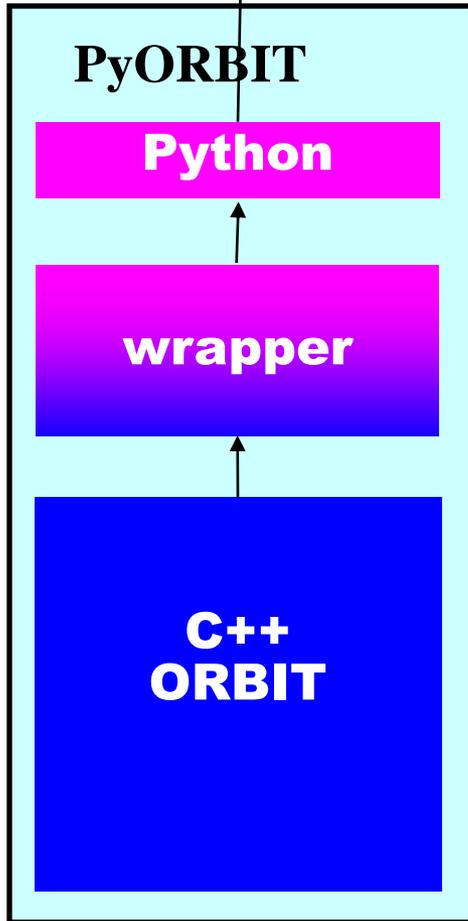
Master equation $\rightarrow \dot{\rho}_{nm}(t) = D[\rho] \quad m, n = 1, N \quad N = 1^2 + 2^2 + 3^2$

196 equations

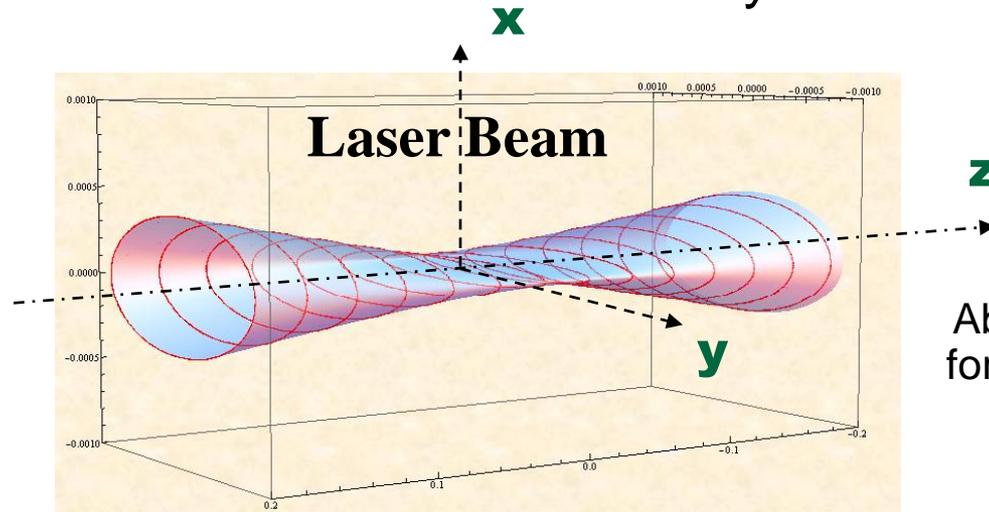


Computer application for laser stripping

user

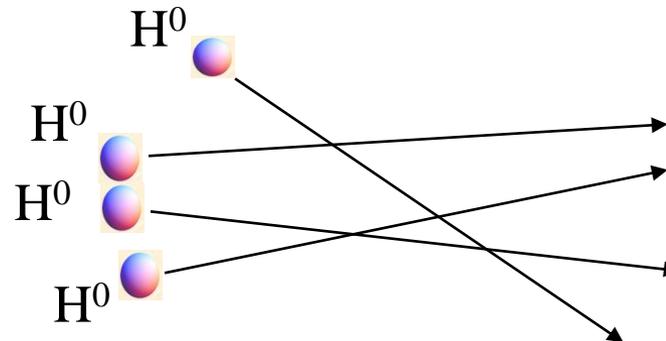


We can optimize parameters of laser beam and H^+ beam for better stripping efficiency



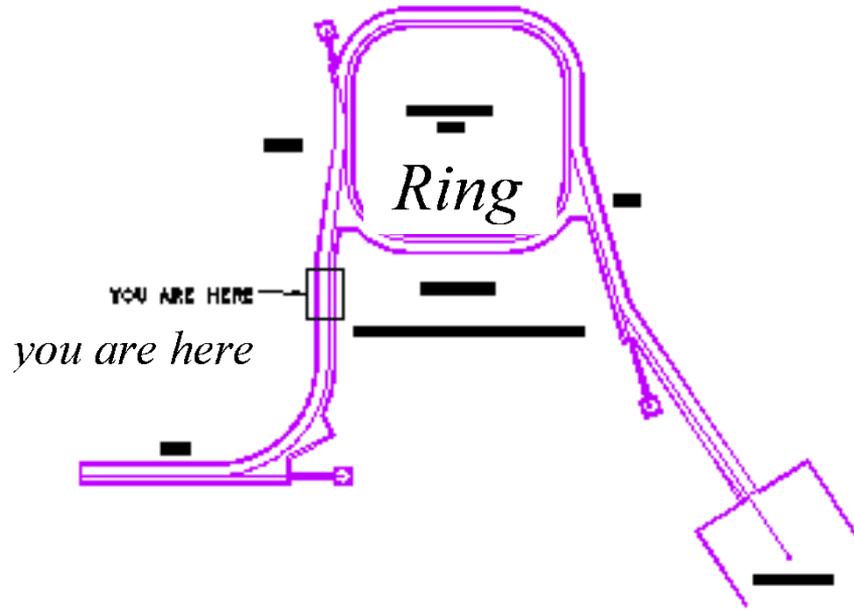
About 5 parameters for description of the laser beam

Hydrogen Beam



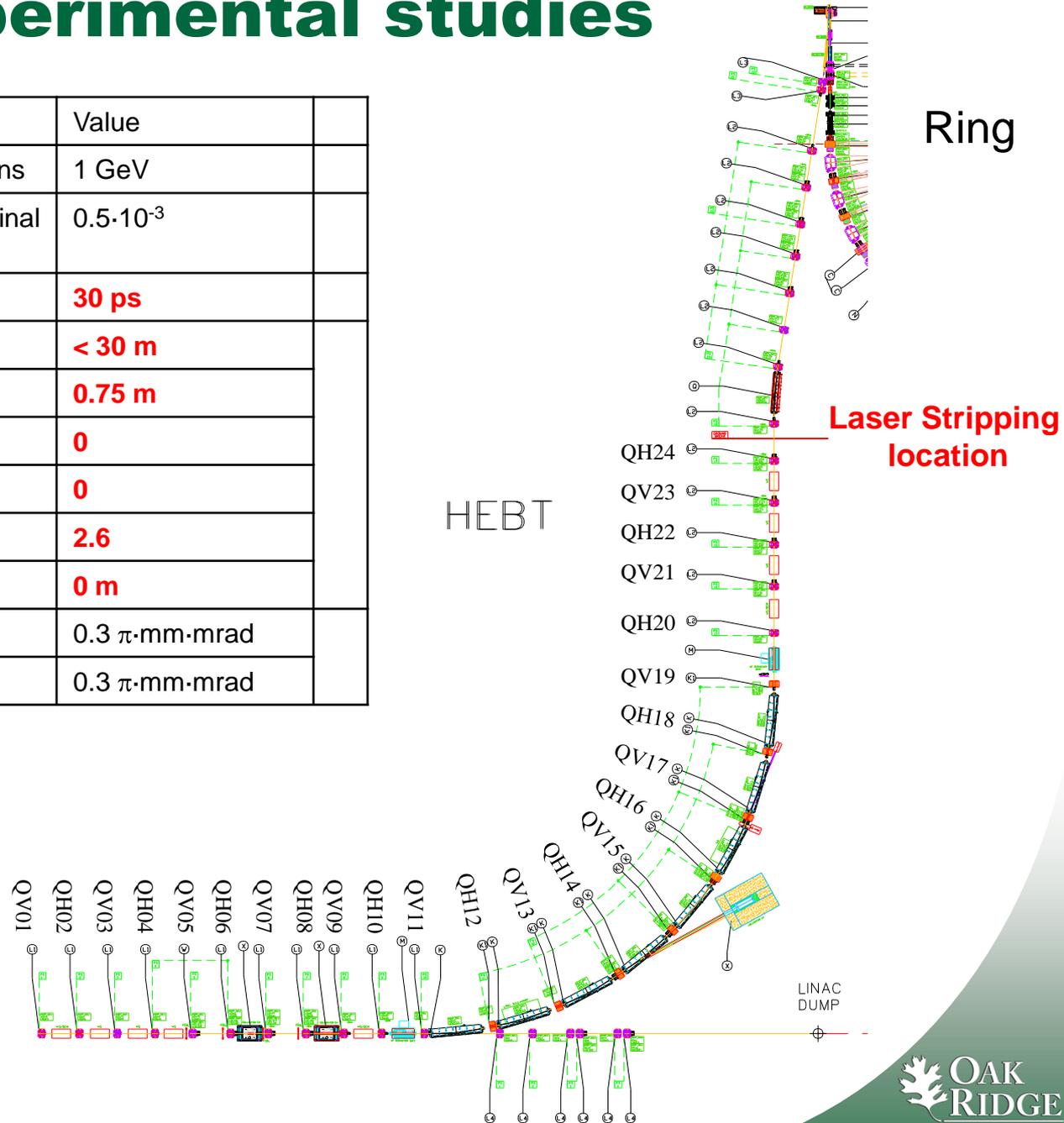
Emittance parameters for description of the H^0 beam

Location for the new experiment



Experimental studies

Parameter	Value	
The longitudinal energy of the ions	1 GeV	
Relative spread of the longitudinal energy σ_T/T	$0.5 \cdot 10^{-3}$	
τ_{FWHM} micro-bunch duration	30 ps	
β_x	< 30 m	
β_y	0.75 m	
α_x	0	
α_y	0	
D_x' dispersion derivative	2.6	
D_x dispersion	0 m	
ϵ_x	$0.3 \pi \cdot \text{mm} \cdot \text{mrad}$	
ϵ_y	$0.3 \pi \cdot \text{mm} \cdot \text{mrad}$	



Preliminary tune of dispersion function

Laser Stripping interaction point



Calculation and optimization
 reset step =

Power supply of quadrupoles

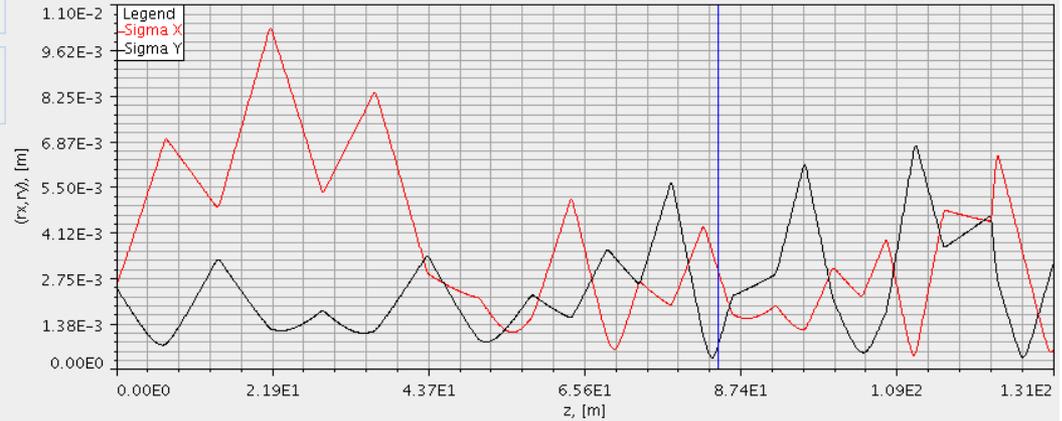
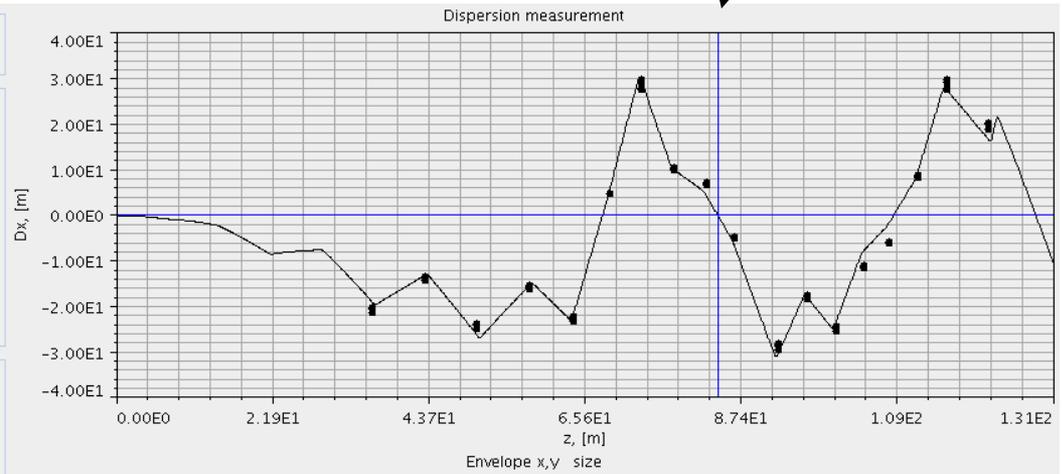
optimize	abs(I)	B+B_Book
<input checked="" type="checkbox"/> QH_12t18	246.764025	-1.540894
<input checked="" type="checkbox"/> QV_13t19	353.279226	2.202
<input checked="" type="checkbox"/> QH_20	247.225846	-3.24177
<input checked="" type="checkbox"/> QV_21	151.873247	2.0007
<input checked="" type="checkbox"/> QH_22	286.874364	-3.758714
<input checked="" type="checkbox"/> QV_23	256.115489	3.3659
<input checked="" type="checkbox"/> QH_24	236.47661	-3.1059

Dispersion measurement at the IP

measure	disp_x	disp_x_der
model =	0.034993	-2.6606
meas. =	2.242689	-2.961946
aver meas. =	2.197718	-2.940019
pvLogId =	14408234	

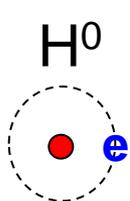
Twiss parameters at the IP

model	beta_y	radius_y	alpha_x
	1.892937	0.000727	4.472652

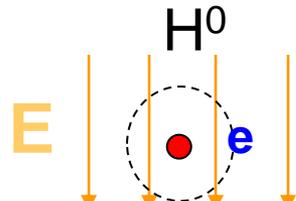


Laser stripping via a broad shape resonance

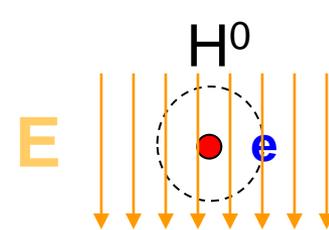
Hydrogen atom in a strong electric field



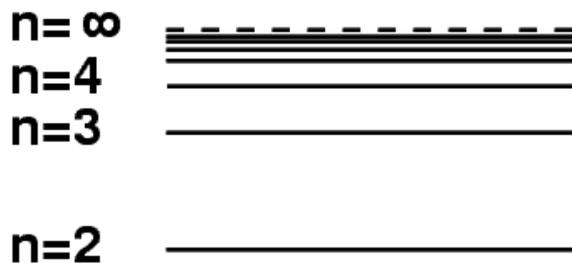
No field



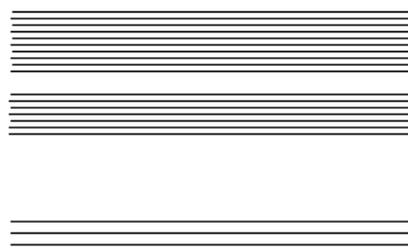
Weak field



Strong field



splitting



broadening



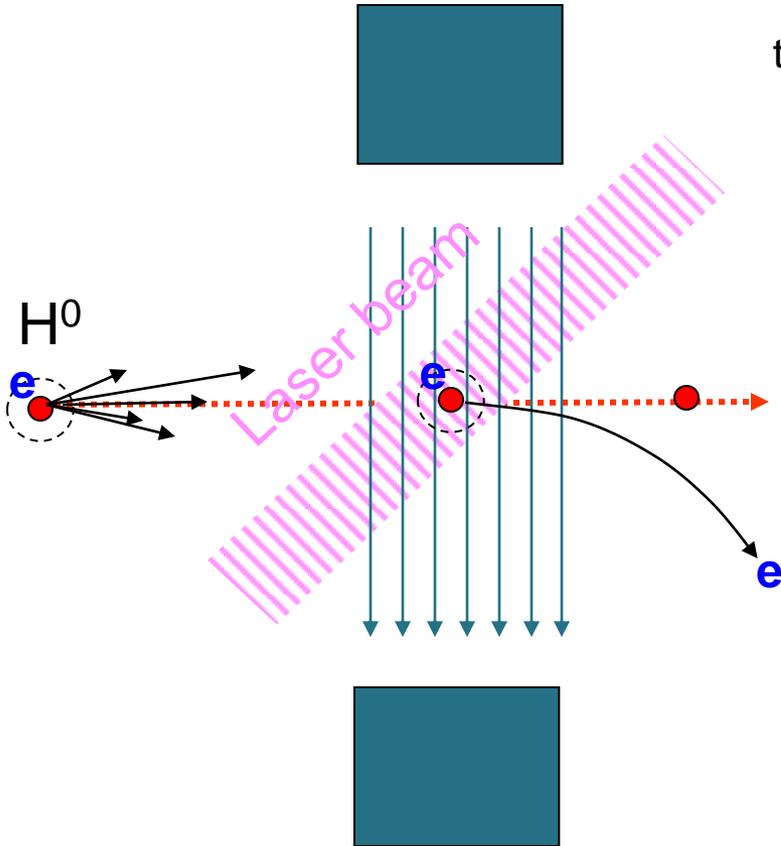
Continuous spectrum

$$E = -\frac{1}{2n^2}$$

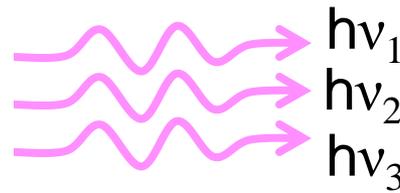


Third way to compensate energy spread

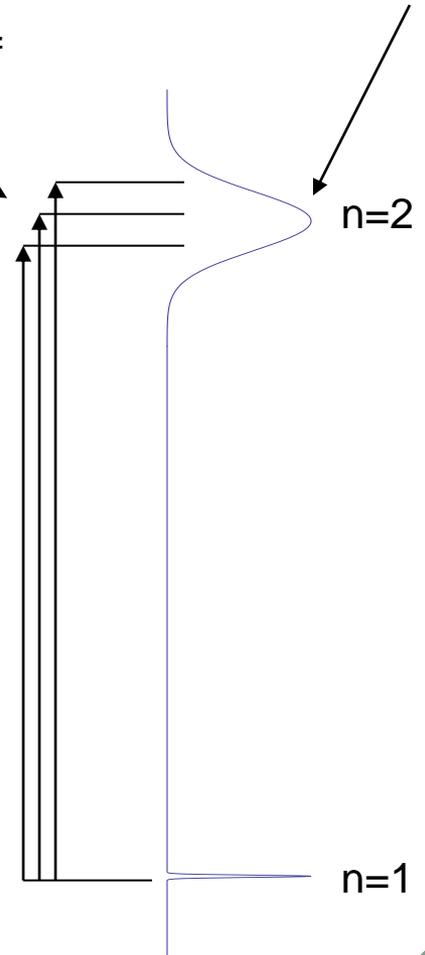
High-field Dipole Magnet



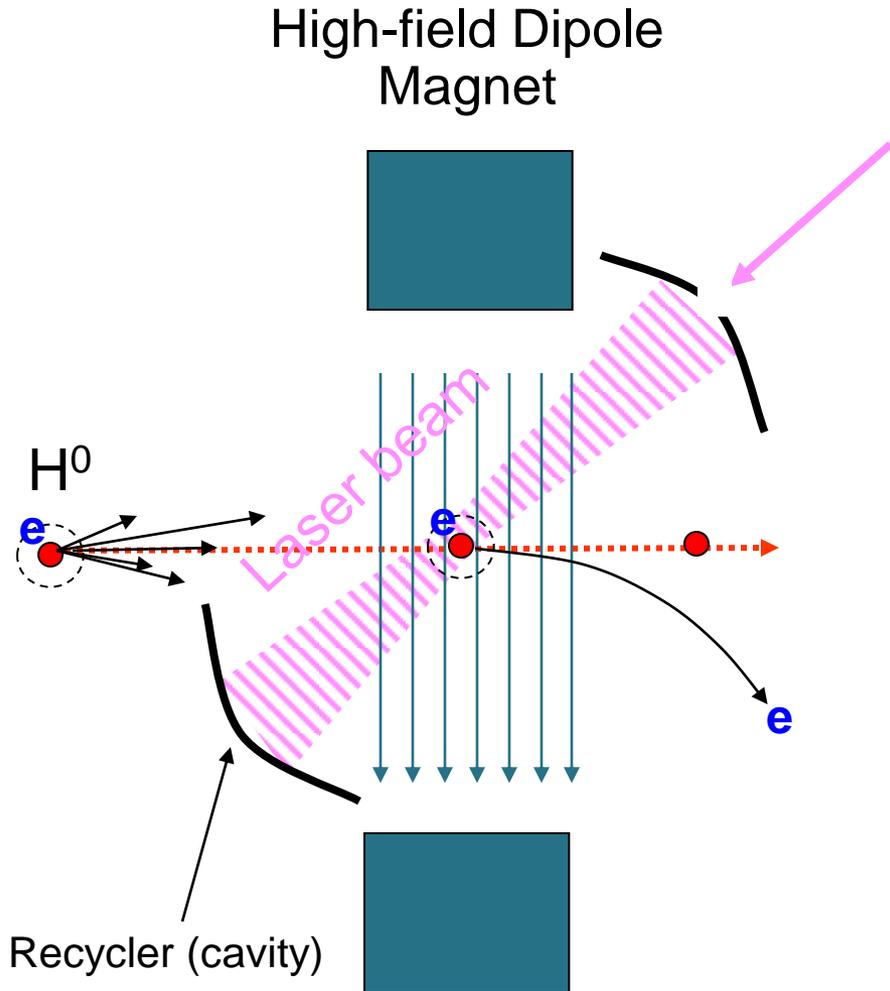
Different particles will see the different frequencies of laser



But they will all be excited because of the broad shape resonance



Advantages of the scheme



- Possibility to recycle laser beam
- No spontaneous transition losses
- Small emittance growth
- Using just one stripping dipole magnet (the scheme is more compact)
- Can be applied for energy 4 GeV and more (considered as disadvantage)

Selected references for more details

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- V. Danilov, A. Aleksandrov, S. Assadi, S. Henderson, N. Holtkamp, T. Shea, A. Shishlo, Y. Braiman, Y. Liu, J. Barhen, and T. Zacharia, Phys. Rev. ST Accel. Beams 6, 053501 (2003).
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- T. Bergeman, C. Harvey, K. B. Butterfield, H. C. Bryant, D. A. Clark, P. A. M. Gram, D. MacArthur, M. Davis, J. B. Donahue, J. Dayton, and W.W. Smith, Phys. Rev. Lett. 53, 775 (1984).
- T. Gorlov and V. Danilov, Phys. Rev. ST Accel. Beams 13, 074002 (2010).