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# Electron Beam Polarimetry at Jefferson Lab

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1. Motivation: Why do we care so much about polarimetry?
2. Overview of JLab polarimeters
3. Cross-polarimeter comparisons
4. Future directions

# Precision Polarimetry

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- A major motivation for precision polarimetry at JLab are for experiments that test the Standard Model
- The Standard Model is remarkably successful – but there is something “missing” → too many free parameters in the theory
- To search for physics beyond the Standard Model we either need to make measurements at:
  - Higher energies → LHC or,
  - Higher precision → **JLAB**
- Knowledge of beam polarization is a limiting systematic in precision Standard Model tests ( $Q_{\text{Weak}}$ , parity violation in Deep Inelastic Scattering )
  - **Such experiments require 1% (or better) polarimetry**
- Other, demanding nuclear physics experiments (strange quarks in the nucleon, neutron skin in nuclei) also benefit from precise measurements of beam polarization

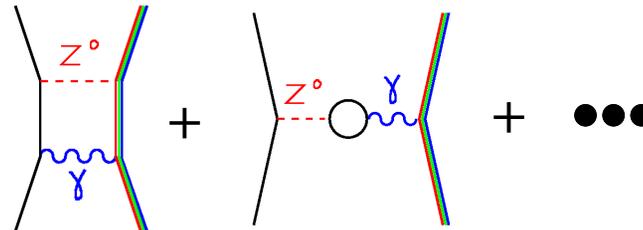
# Testing the Standard Model

Weinberg-Salam theory of electro-weak interactions

$$L = \frac{g}{\sqrt{2}} (J_{\mu}^{-} W_{\mu}^{+} + J_{\mu}^{+} W_{\mu}^{-}) + \frac{g}{\cos \theta_W} (J_{\mu}^{(3)} - \sin^2 \theta_W J_{\mu}^{e.m.}) Z_{\mu} + g \sin \theta_W J_{\mu}^{e.m.} A_{\mu}$$

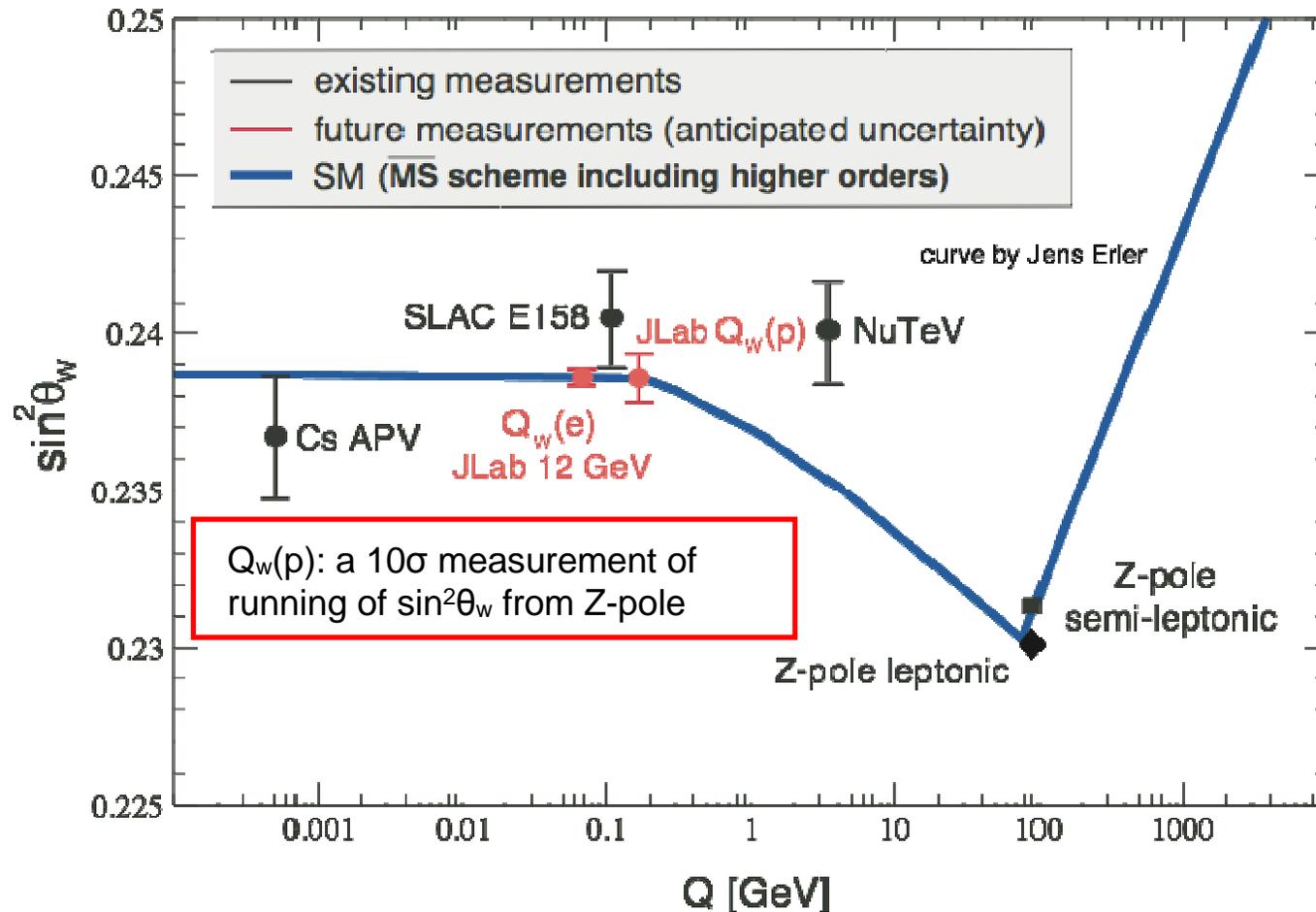
$\theta_W$  = weak mixing angle

The evolution of  $\sin^2 \theta_W$  with  $Q^2$  can be predicted in the context of the Standard Model

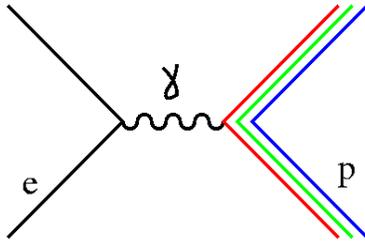


→ deviations from this prediction are the signature of new physics, beyond the Standard Model (super-symmetry, Leptoquarks, etc.)

# Running of $\sin^2\theta_w$

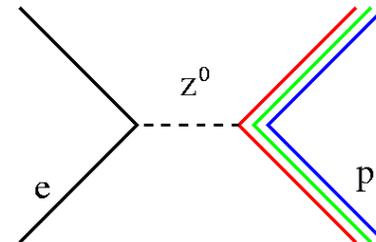


# $Q_{\text{Weak}}$ in Hall C



measures  $Q^p$  - proton's electric charge

As  $Q^2 \rightarrow 0$



measures  $Q_{\text{weak}}^p$  - proton's weak charge

The  $Q_{\text{weak}}$  experiment measures the parity-violating analyzing power  $A_z$

$$A_z \xrightarrow[Q^2 \rightarrow 0]{\theta \rightarrow 0} \frac{-G_F}{4\pi\alpha\sqrt{2}} [Q^2 Q_{\text{weak}}^p + Q^4 B(Q^2)]$$

Contains  $G_{E,M}^Y$  and  $G_{E,M}^Z$ ,  
Extracted using global fit  
of existing PVES experiments!

$$Q_{\text{weak}}^p = 1 - 4 \sin^2 \theta_W \sim 0.072 \text{ (at tree level)}$$

# Projected Uncertainties

Uncertainty	$\Delta A_z/A_z$	$\Delta Q_w/Q_w$
Statistical (2,544 hours at 180 $\mu$ A)	2.1%	3.2%
<b>Systematic:</b>		<b>2.6%</b>
Hadronic structure uncertainties	---	1.5%
Beam polarimetry	1.0%	1.5%
Absolute $Q^2$ determination	0.5%	1.0%
Backgrounds	0.5%	0.7%
Helicity correlated beam properties	0.5%	0.7%
<b>Total:</b>	<b>2.5%</b>	<b>4.1%</b>

Nominal error magnification due to the 33% hadronic dilution  $\times$  1.49. The enhancement for the  $Q^2$  term is somewhat larger.

→ ***Polarimetry approaching limiting systematic uncertainty***

## Experiments Requiring Precision Polarimetry

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$Q_{\text{Weak}}$  ( $Q^2$  evolution of weak mixing angle)

→ 1% at 1 GeV

PREX (neutron distribution in Pb)

→ 1% at 850 MeV

Møller Scattering at 12 GeV\* (more weak mixing angle)

→ 1% at 11 GeV

Parity Violation in Deep Inelastic Scattering (DIS-Parity)

→ 0.5% at 11 GeV

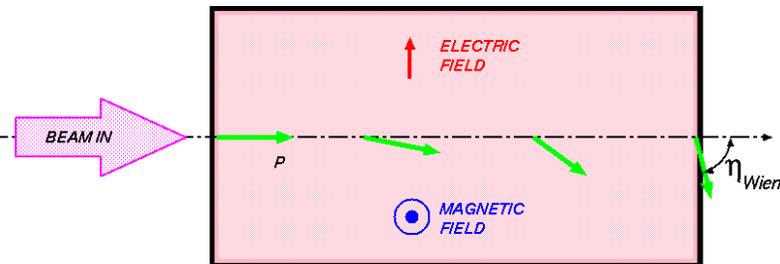
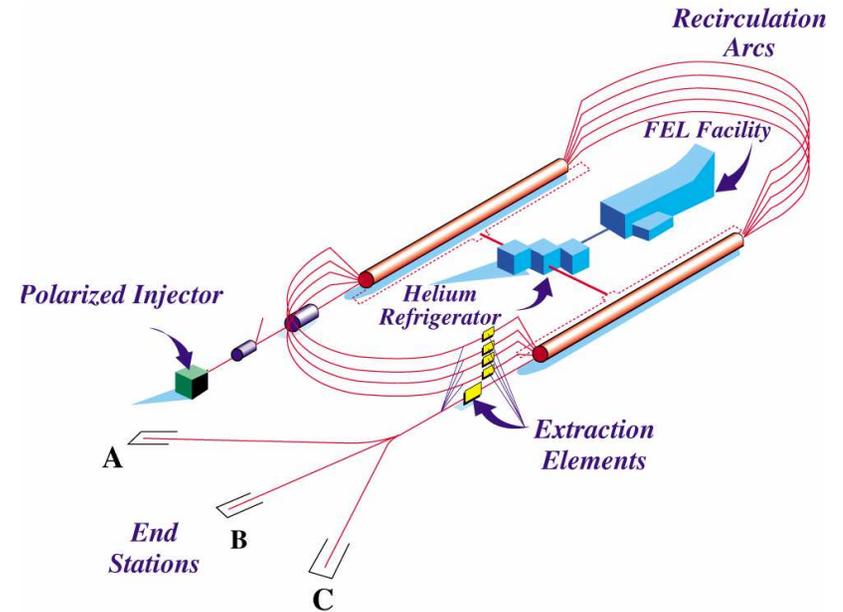
\* = not yet officially proposed (but it will be)

# Polarized Electrons at Jefferson Lab

- Polarized electrons generated “at the source” using Superlattice GaAs photocathode
- Electrons polarized in the plane of the accelerator
  - spin direction precesses as beam circulates (up to 5 times) through machine

$$\phi_{spin} = \frac{g - 2}{2} \frac{E_{beam}}{m_e} \theta_{bend}$$

- Spin direction manipulated at source using Wien filter to get long. Polarization in Halls
- JLab now routinely provides electron beam polarizations >80% to experimental halls



# JLab Polarimetry Techniques

- Three different processes used to measure electron beam polarization at JLab
  - Møller scattering:  $\vec{e} + \vec{e} \rightarrow e + e$  , atomic electrons in Fe (or Fe-alloy) polarized using by external magnetic field
  - Compton scattering:  $\vec{e} + \vec{\gamma} \rightarrow e + \gamma$  , laser photons scatter from electron beam
  - Mott scattering:  $\vec{e} + Z \rightarrow e$  , spin-orbit coupling of electron spin with (large Z) target nucleus
- Each has advantages and disadvantages in JLab environment

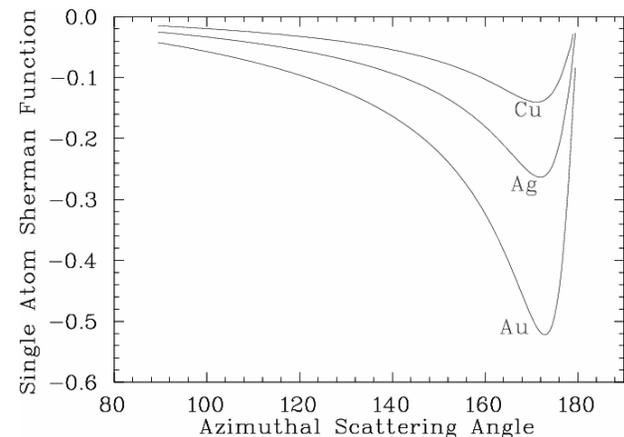
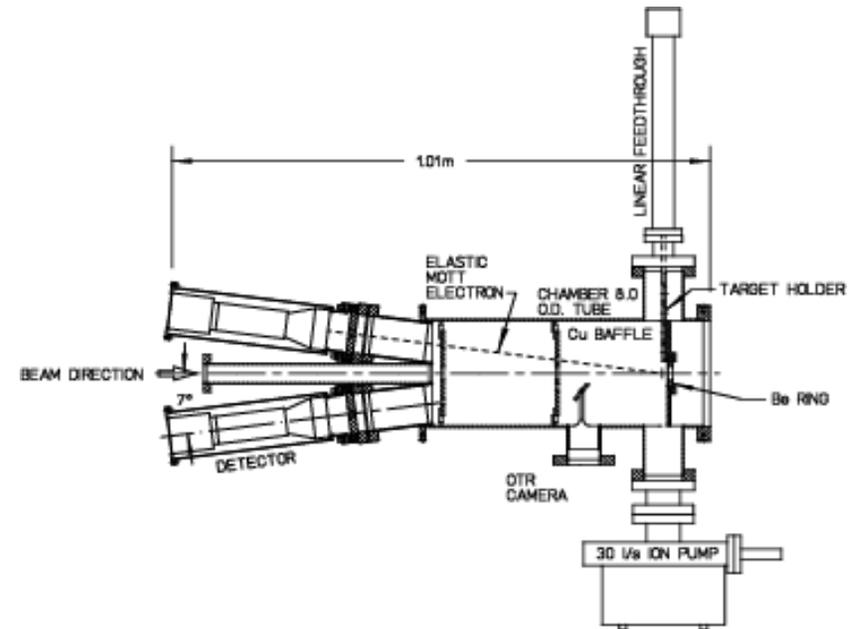
Method	Advantage	Disadvantage
Compton	Non-destructive	Can be time consuming, systematics energy dependent
Møller	Rapid, precise measurements	Destructive, low current only
Mott	Rapid, precise measurements	Does not measure polarization at the experiment

# 5 MeV Mott Polarimeter

- Mott polarimeter located in the 5 MeV region of the CEBAF injector
- Target must be thin, large Z material → 1 μm Au foil
- Asymmetry maximized near 172°, given by

$$A = \frac{N_r - N_l}{N_r + N_l} = P_{beam} S(\theta)$$

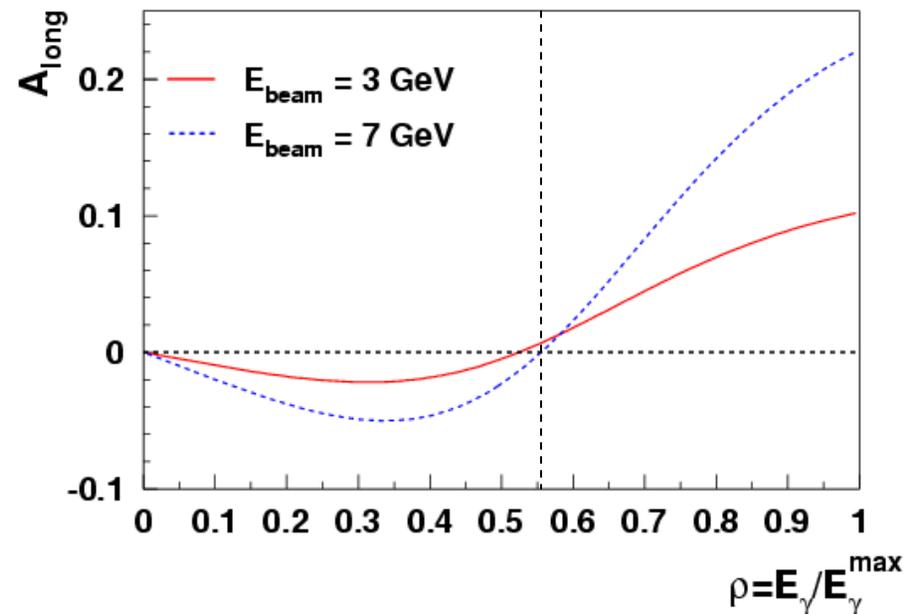
- $S(\theta)$  is the Sherman function → must be calculated from e-nucleus cross section
- Knowledge of Sherman function dominant systematic uncertainty ~ **1.0%**



# Compton Polarimetry at JLab

Two main challenges for Compton polarimetry at JLab

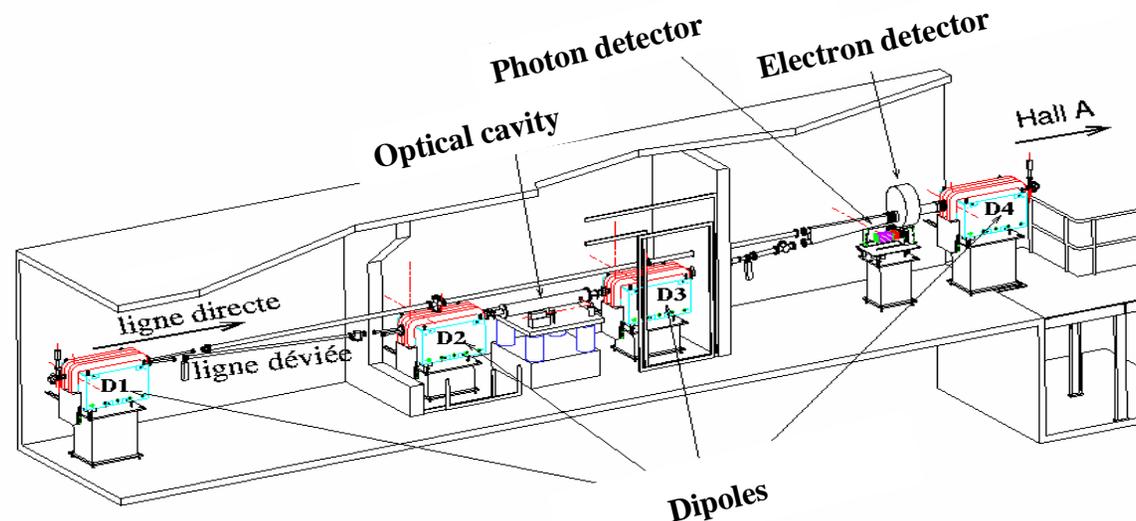
- Low beam currents ( $\sim 100 \mu\text{A}$ )
  - Measurements can take on the order of hours
  - Makes systematic studies difficult
- Relatively small asymmetries
  - Smaller asymmetries lead to harder-to-control systematics



- Strong dependence of asymmetry on  $E_{\gamma}$  is a challenge  
→ Understanding of detector response crucial

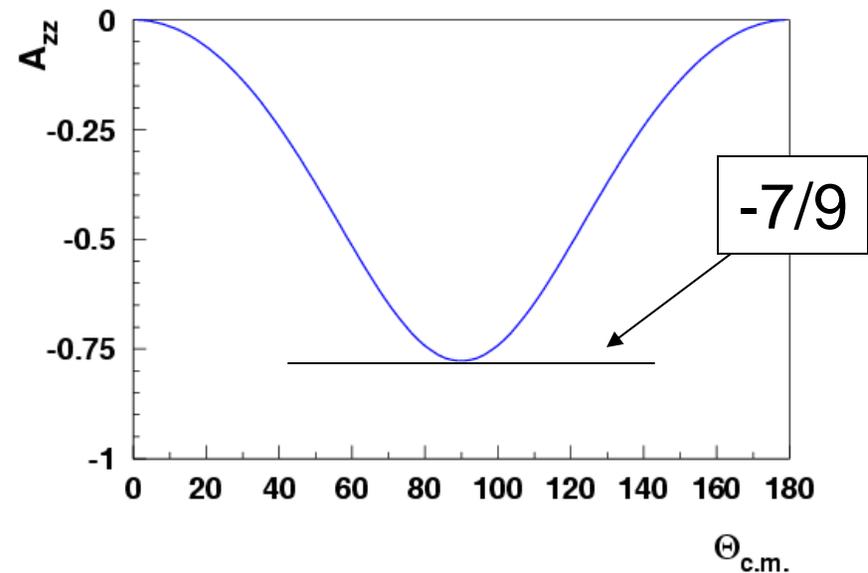
# Hall A Compton Polarimeter

- Hall A Compton polarimeter uses high gain Fabry-Perot cavity to create  $\sim 1$  kW of laser power in IR (1064 nm)
- Detects *both* scattered electron and backscattered  $\gamma \rightarrow 2$  independent measurements, coincidences used to calibrate  $\gamma$  detector
- Systematic errors quoted at **1%** level for recent HAPPEX experiments @ 3 GeV [PRL 98 (2007) 032301]



# Møller Polarimetry at JLab

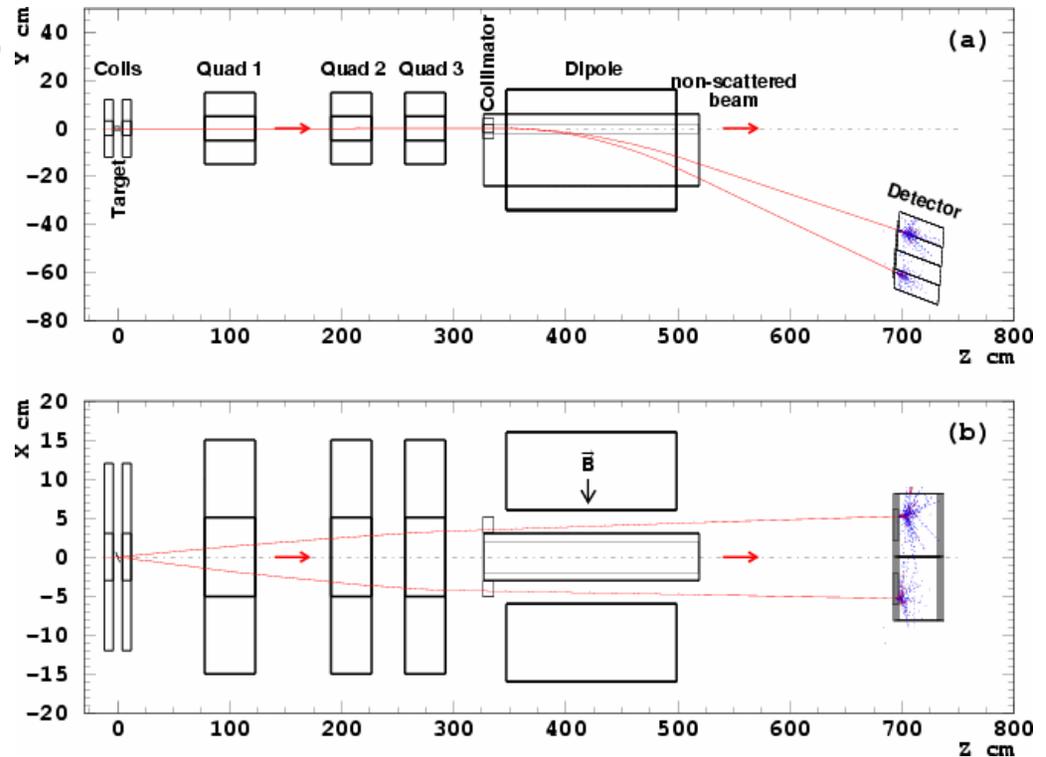
- Møller polarimetry benefits from large long. asymmetry  
→  $-7/9$ 
  - Asymmetry independent of energy
  - Relatively slowly varying near  $\theta_{\text{cm}}=90^\circ$
  - Large asymmetry diluted by need to use iron foils to create polarized electrons
    - $P_e \sim 8\%$
- Rates are large, so rapid measurements are easy
- Need to use Fe or Fe-alloy foils means measurement must be destructive



- Making measurements at high beam currents challenging

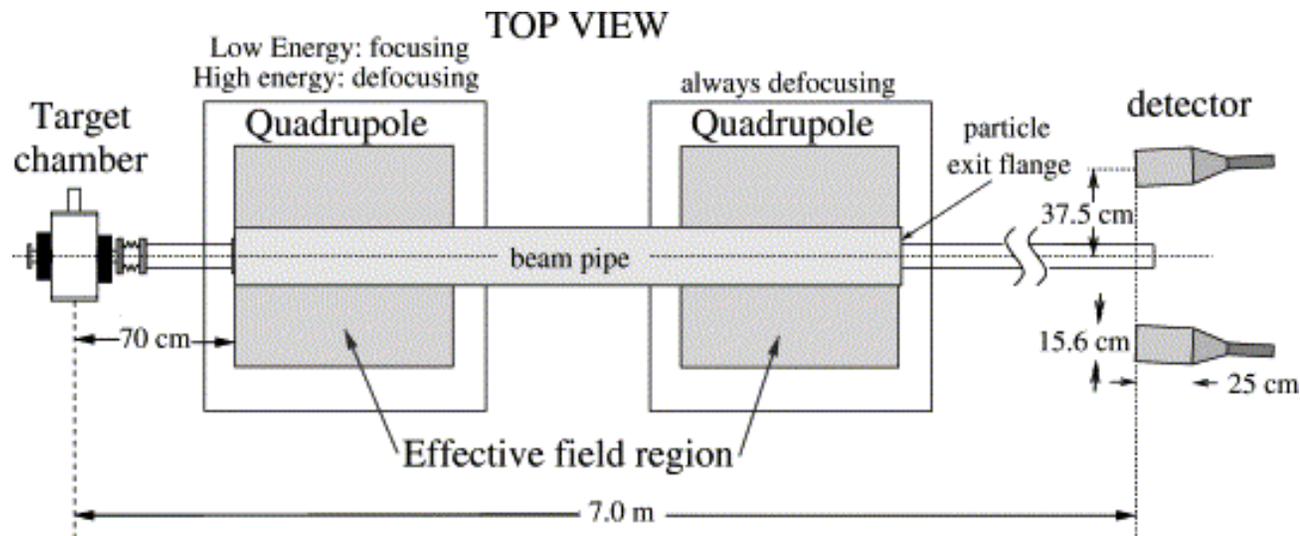
# Hall A Møller Polarimeter

- Target =supermendeur foil, polarized in-plane
  - Low field applied (240 G)
  - Tilted  $20^\circ$  relative to beam direction
  - Target polarization known to  $\sim 2\%$   $\rightarrow$  *this will improve*
- Large acceptance of detectors mitigates potentially large systematic unc. from **Levchuk** effect (atomic Fermi motion of bound electrons)
- Large acceptance also leads to large rates - dead time corrections cannot be ignored, but are tractable



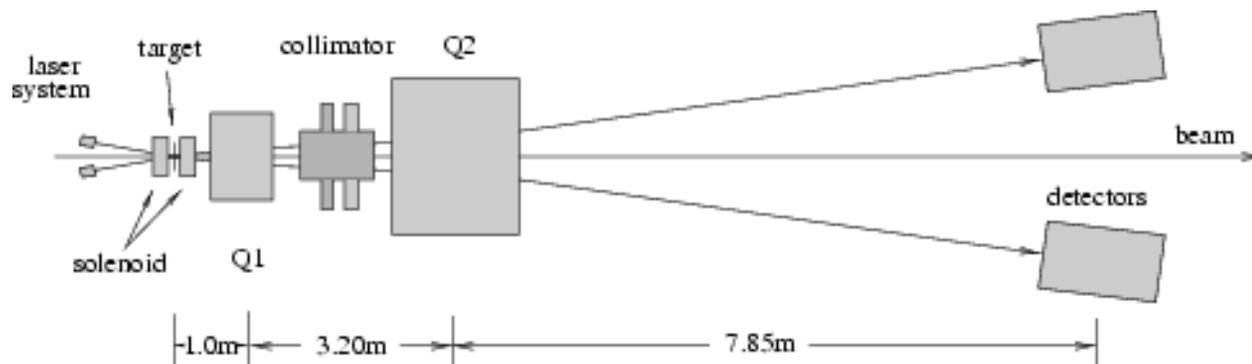
# Hall B Møller Polarimeter

- Hall B Møller uses similar target design as Hall A → Fe alloy in weak magnetic field
- Two-quadrupole system rather than QQQD
- Detector acceptance not as large – Levchuk effect corrections important
- Dominant systematics [NIM A 503 (2003) 513]
  - Target polarization ~ **1.4%**
  - Levchuk effect ~ **0.8%**



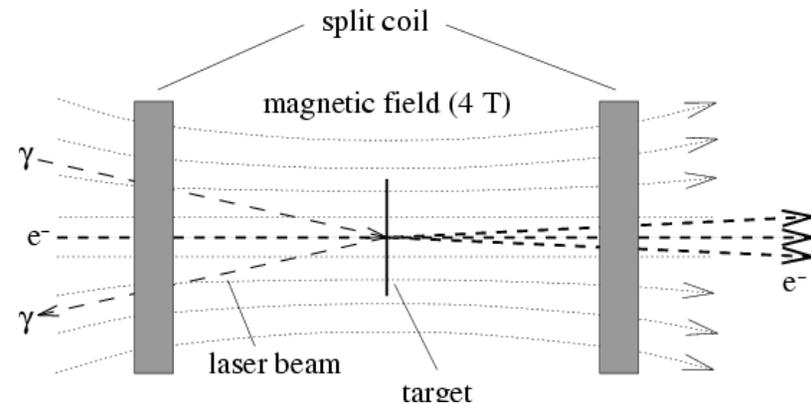
# Hall C Møller Polarimeter

- 2 quadrupole optics maintains constant tune at detector plane
- “Moderate” (compared to Hall A) acceptance mitigates Levchuk effect → still a non-trivial source of uncertainty
- Target = pure Fe foil, brute-force polarized out of plane with 3-4 T superconducting magnet
- Total systematic uncertainty = **0.47%** [NIM A 462 (2001) 382]



# Hall C Møller Target

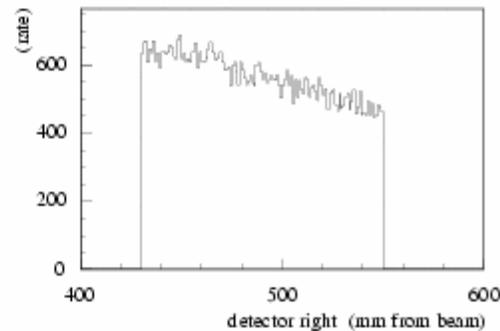
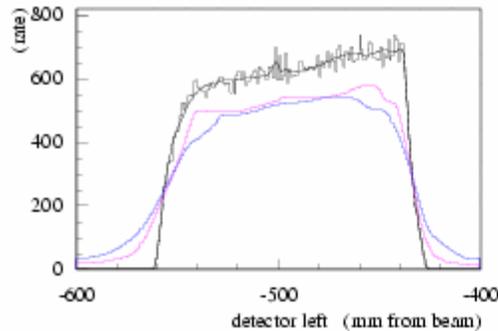
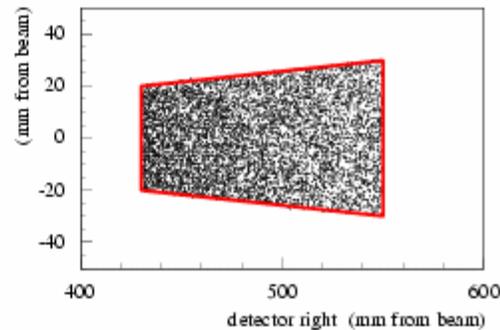
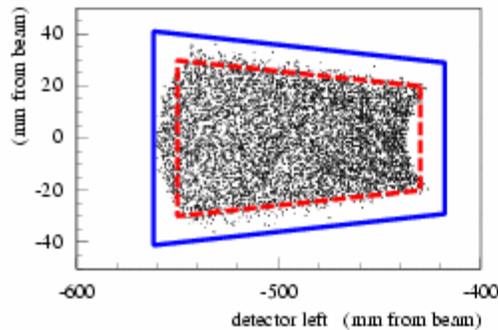
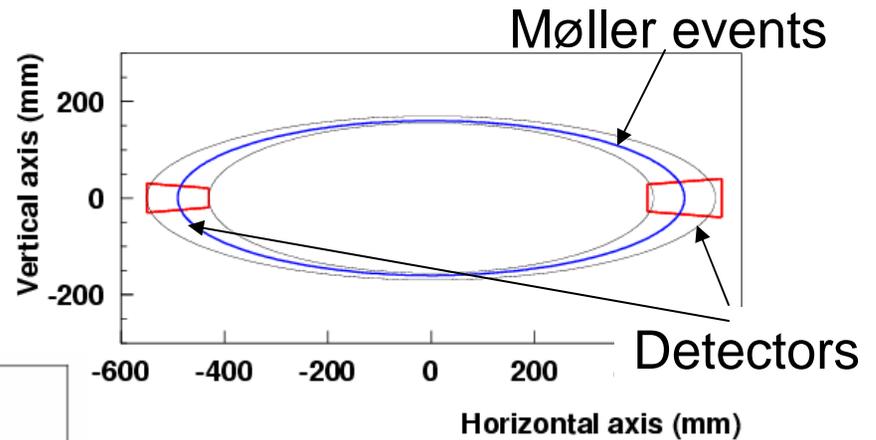
- Fe-alloy, in plane polarized targets typically result in systematic errors of 2-3%
  - Require careful measurement magnetization of foil
- Pure Fe saturated in 4 T field
  - Spin polarization well known  $\rightarrow$  0.25%
  - Temperature dependence well known
  - No need to directly measure foil polarization



Effect	$M_s [\mu_B]$	error
Saturation magnetization ( $T \rightarrow 0$ K, $B \rightarrow 0$ T)	2.2160	$\pm 0.0008$
Saturation magnetization ( $T=294$ K, $B=1$ T)	2.177	$\pm 0.002$
Corrections for $B=1 \rightarrow 4$ T	0.0059	$\pm 0.0002$
Total magnetization	2.183	$\pm 0.002$
Magnetization from orbital motion	0.0918	$\pm 0.0033$
Magnetization from spin	2.0911	$\pm 0.004$
Target electron polarization ( $T=294$ K, $B=4$ T)	0.08043	$\pm 0.00015$

# Hall C Møller Acceptance

Optics designed to maintain similar acceptance at detectors independent of beam energy



Collimators in front of Pb:Glass detectors define acceptance

One slightly larger to reduce sensitivity to Levchuk effect

# Hall C Møller Systematics (I)

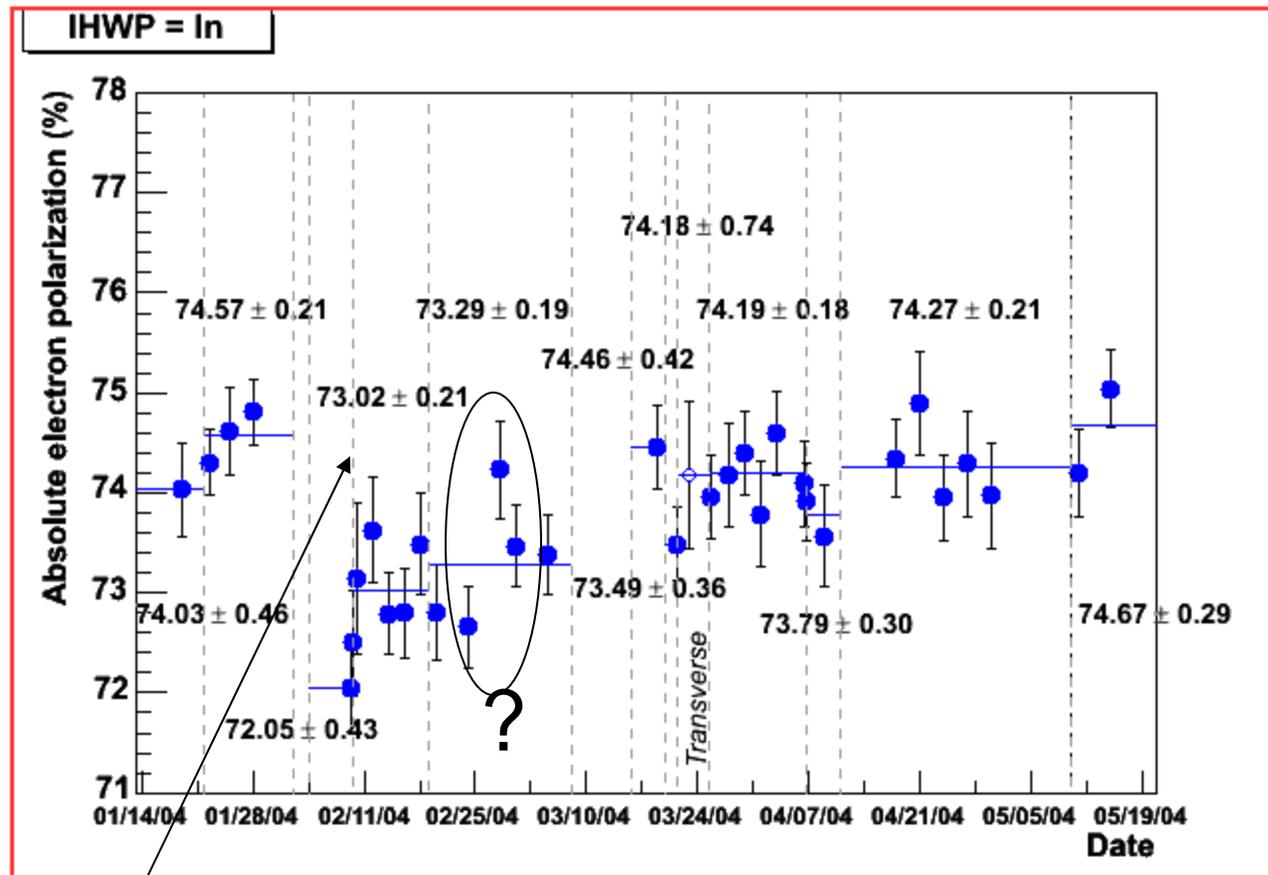
Systematic error budget from NIM article

Idealized?

A particular measurement under specific conditions may approach this level of uncertainty – ***but we want the polarization for the same conditions as the experiment!***

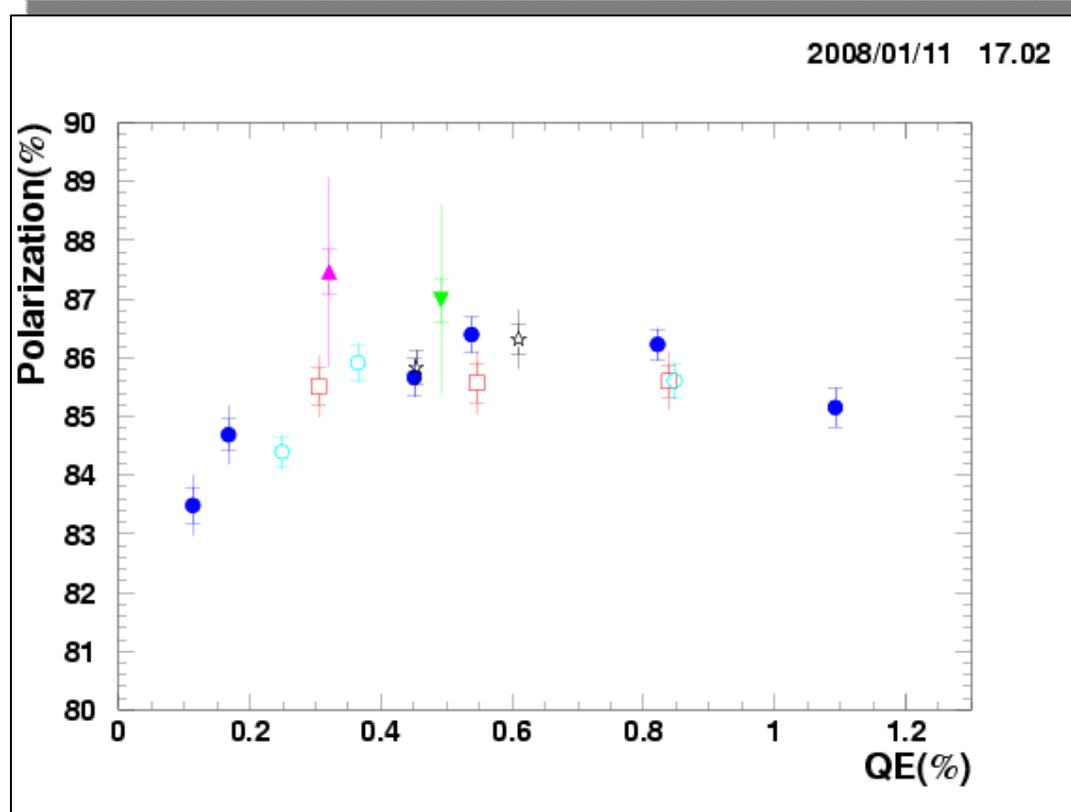
Source	Uncertainty	dAsy./Asy. (%)
Beam position x	0.5 mm	0.15
Beam position y	0.5 mm	0.03
Beam direction x	0.15 mr	0.04
Beam direction y	0.15 mr	0.04
Q1 current	2%	0.10
Q2 current	1%	0.07
Q2 position	1 mm	0.02
Multiple Scattering	10%	0.12
Levchuk effect	10%	0.30
Collimator positions	0.5 mm	0.06
Target temperature	50%	0.05
B-field direction	2°	0.06
B-field strength	5%	0.03
Spin polarization in Fe		0.25
<b>Total</b>		<b>0.47</b>

# Hall C Møller during G0 Forward Angle



Each dashed line corresponds to an “event” that *may* have impacted the polarization in machine

# Hall C Møller during GEp



Some hint of dependence of polarization on QE!

# Hall C Møller Systematics (II)

Systematic error budget from G0 Forward Angle expt.

$$dP/P = 1.32\%$$

Source	Uncertainty	dAsy./Asy. (%)
Beam position x	0.5 mm	0.15
Beam position y	0.5 mm	0.03
Beam direction x	0.15 mr	0.04
Beam direction y	0.15 mr	0.04
Q1 current	2%	0.10
Q2 current	1%	0.07
Q2 position	1 mm	0.02
Multiple Scattering	10%	0.12
Levchuk effect	10%	0.30
Collimator positions	0.5 mm	0.06
Target temperature	50%	0.05
B-field direction	2°	<del>0.06</del> 0.37
B-field strength	5%	0.03
Spin polarization in Fe		0.25
Leakage	30 nA	0.2
High current extrap.	1%/40 uA	1.0
Solenoid focusing	100%	0.1
Elec. DT.	100%	0.04
Charge measurement		0.02
Monte Carlo Statistics		0.28
Unknown accelerator changes		0.5
<b>Total</b>		<b>1.32</b>

# JLab Polarimeter Roundup

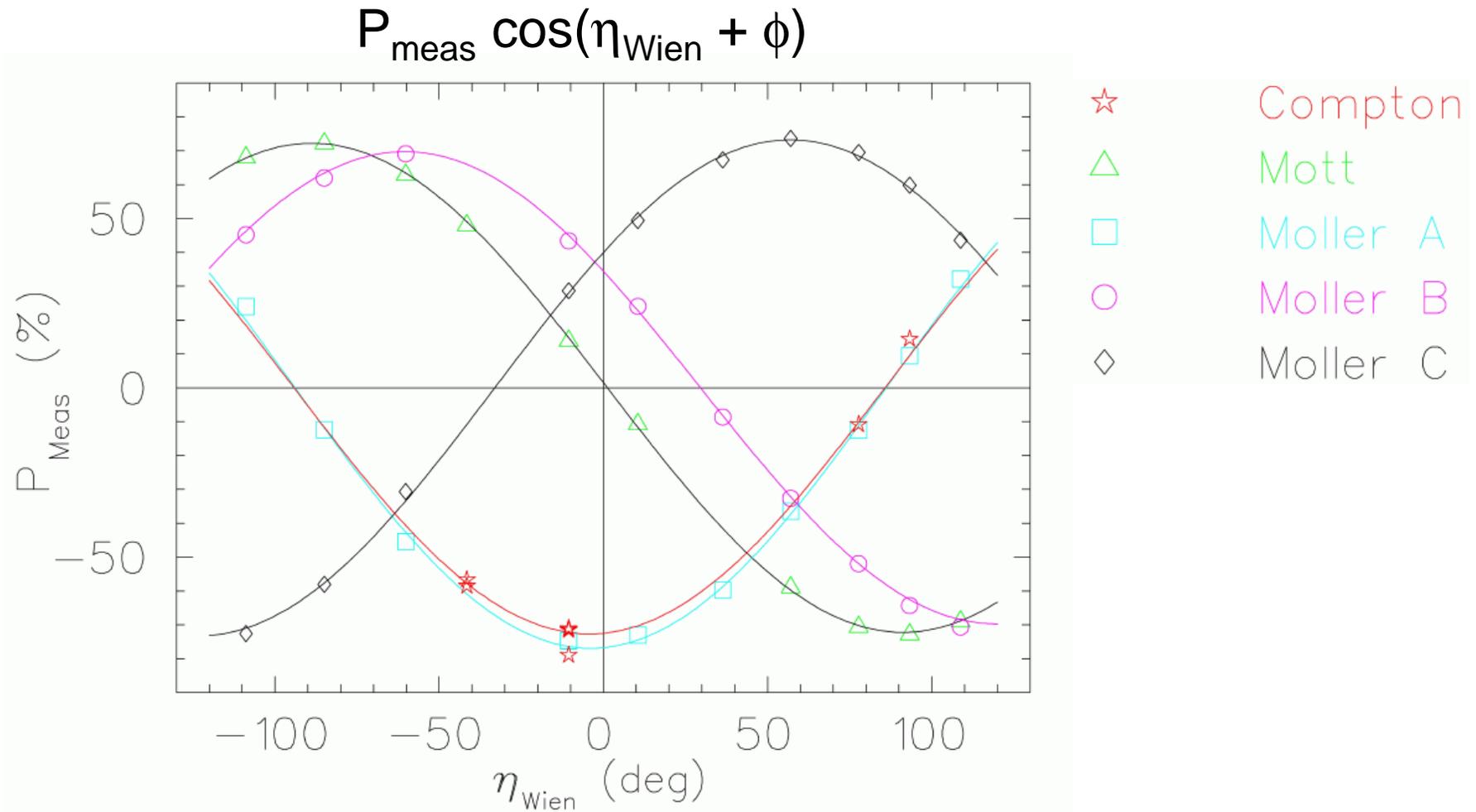
Polarimeter	Relative precision	Limiting systematic
5 MeV Mott	~1%	Sherman function
Hall A Møller	~2-3%	Target polarization
Hall B Møller	1.6% (?)	Target polarization, Levchuk effect
Hall C Møller	0.5% ( $\rightarrow$ 1.3%)	Target polarization, Levchuk effect, high current extrapolation
Hall A Compton	1% (@ > 3 GeV)	Detector acceptance + response

# Spin Dance 2000

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- In July 2000, a multi-hall “Spin Dance” was performed at JLab
- Wien filter in the injector was varied from  $-110^\circ$  to  $110^\circ$ , thus varying degree of longitudinal polarization in each hall
- Purpose was 2-fold
  - Allow precise cross-comparison of JLab polarimeters
  - Extract measurement of beam energy using spin precession through machine
- Results can be found in: *Phys. Rev. ST Accel. Beams* **7**, 042802 (2004)

# Spin Dance 2000 Data



# Polarization Results

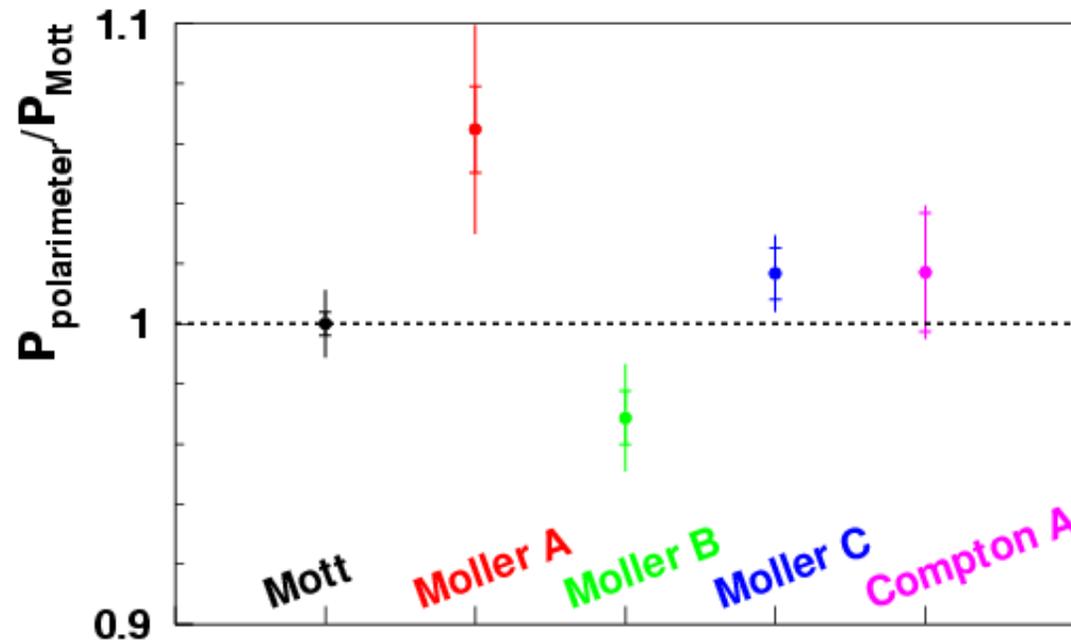
Results shown include statistical errors only

→some amplification to account for non-sinusoidal behavior

Statistically significant

Systematics shown:

Mott	}	1%
Møller C		
Compton		
Møller B	1.6%	
Møller A	3%	



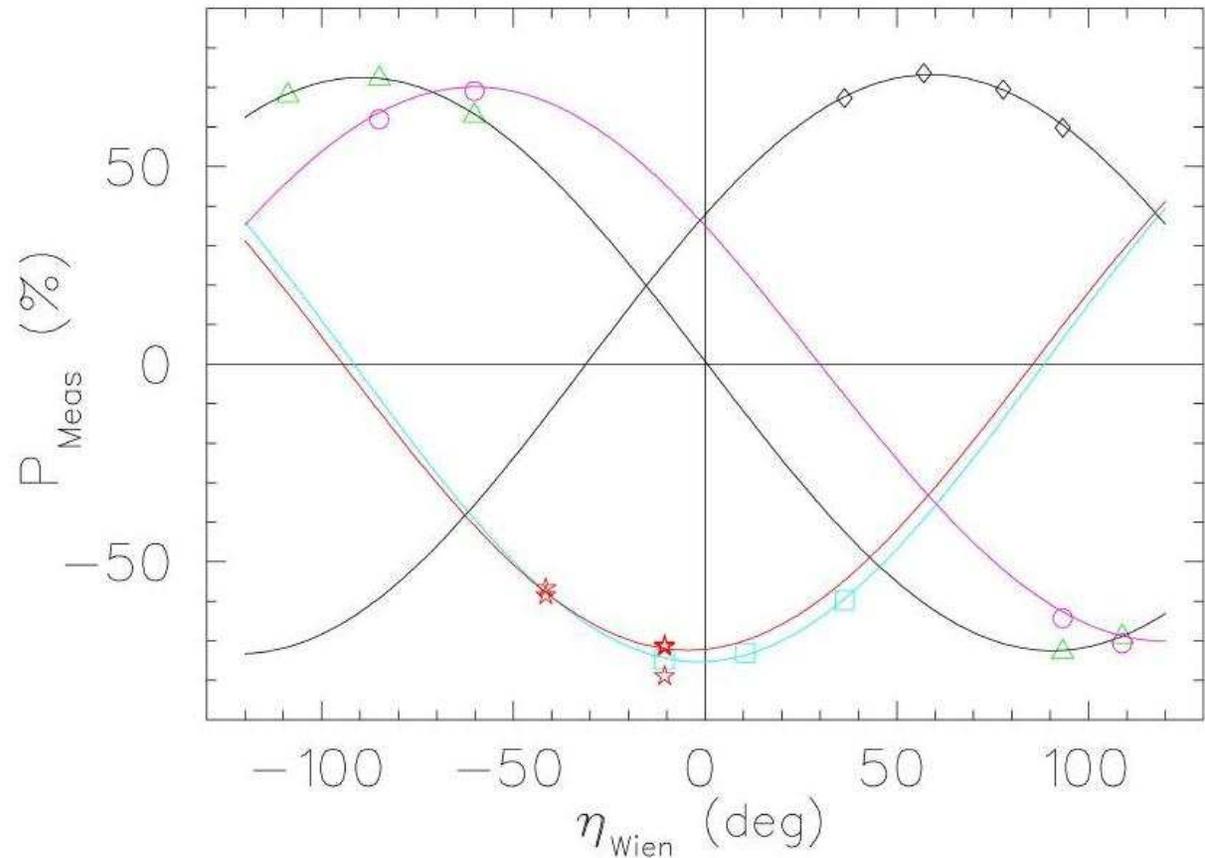
Even including systematic errors, discrepancy still significant

# Reduced Data Set

Hall A, B Møllers  
sensitive to transverse  
components of beam  
polarization

Normally – these  
components eliminated  
via measurements with  
foil tilt reversed, but some  
systematic effects may  
remain

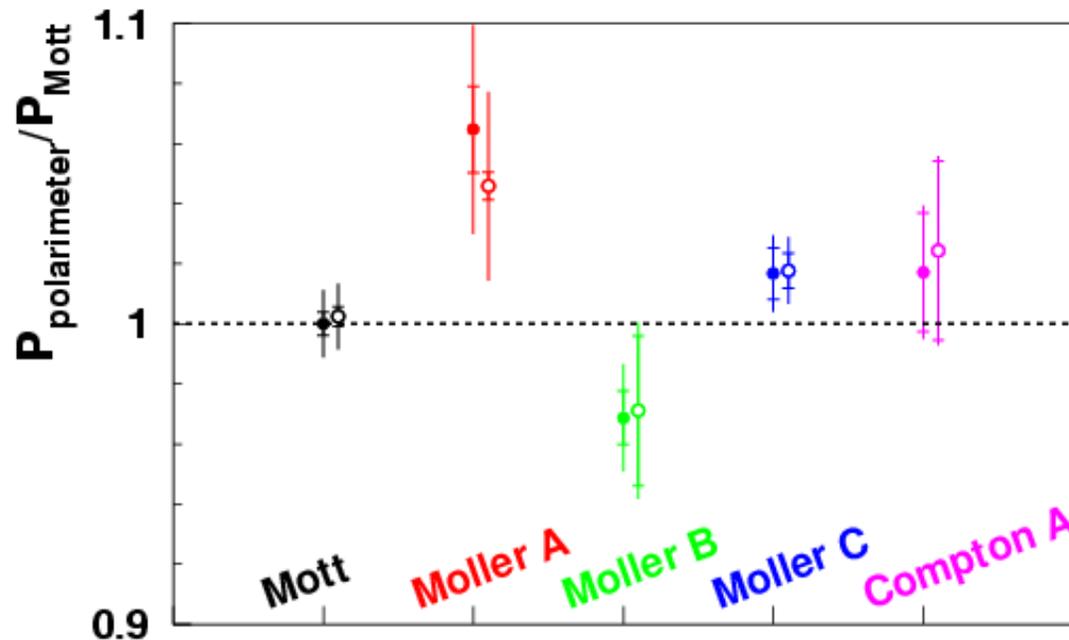
Fit was redone with data  
only within 20% of total  
polarization



# Polarization Results – Reduced Data Set

Agreement improves, but still statistically significant deviations  
→ when systematics included, discrepancy less significant

Further study in Hall A suggests measured foil polarization too big by 1-1.5%



closed circles = full data set  
open circles = reduced data set

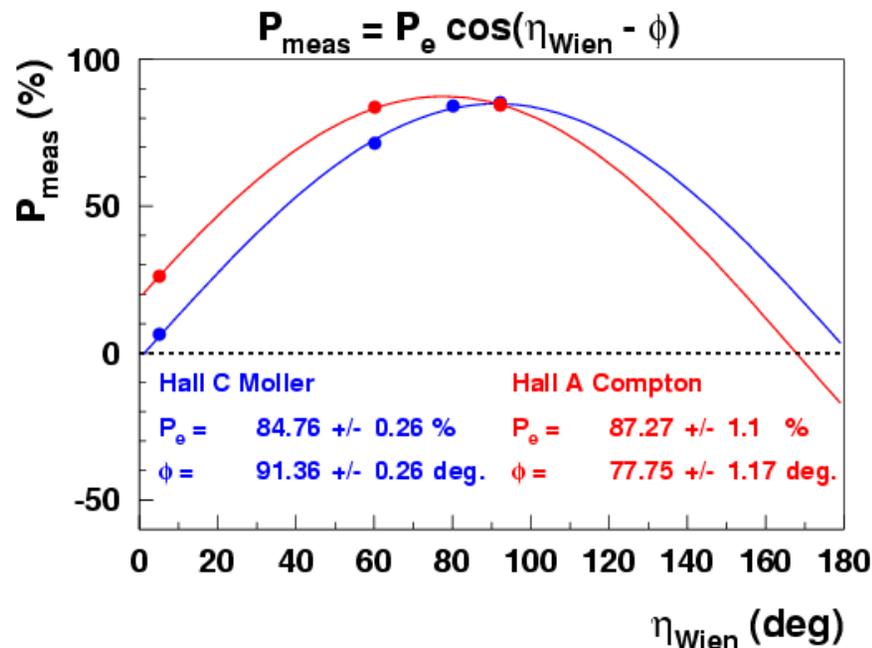
# Spin Dance 200X?

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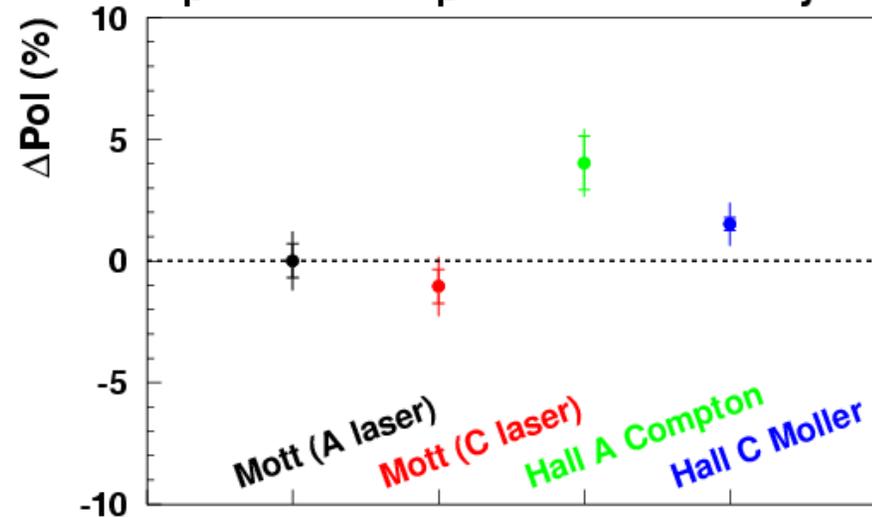
- Since Spin Dance 2000, there has not been another full-blown, cross-hall, polarimeter comparison
- Dedicated time for these measurements difficult to obtain – beam time is precious and there is enormous pressure to complete as much of the physics program as possible
- There are sometimes opportunities for multi-hall comparisons, but usually only when experiments are using polarized beam and polarimeters are already commissioned
- PREX and  $Q_{\text{Weak}}$  will be running in the next few years → this may be an excellent opportunity to push for further studies
- In particular, Hall A Møller implementing Hall C style target
  - Systematics due to target polarization *identical*
  - Comparison (if done carefully) would isolate **instrumental** effects

# Additional Cross-Hall Comparisons

- During G0 Backangle, performed “mini-spin dance” to ensure purely longitudinal polarization in Hall C
- Hall A Compton was also in use, so they participated as well



April 13 2006 Spin Dance Summary



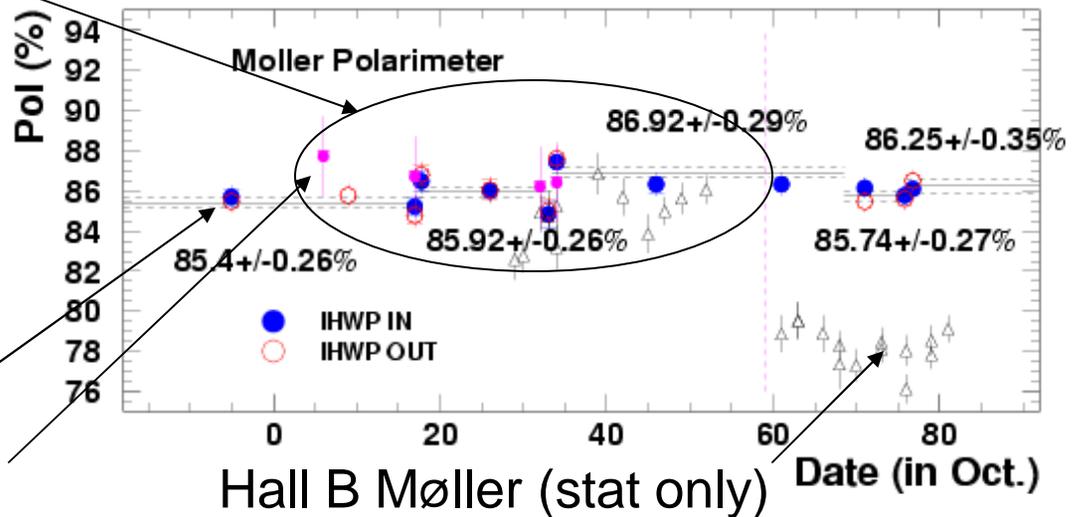
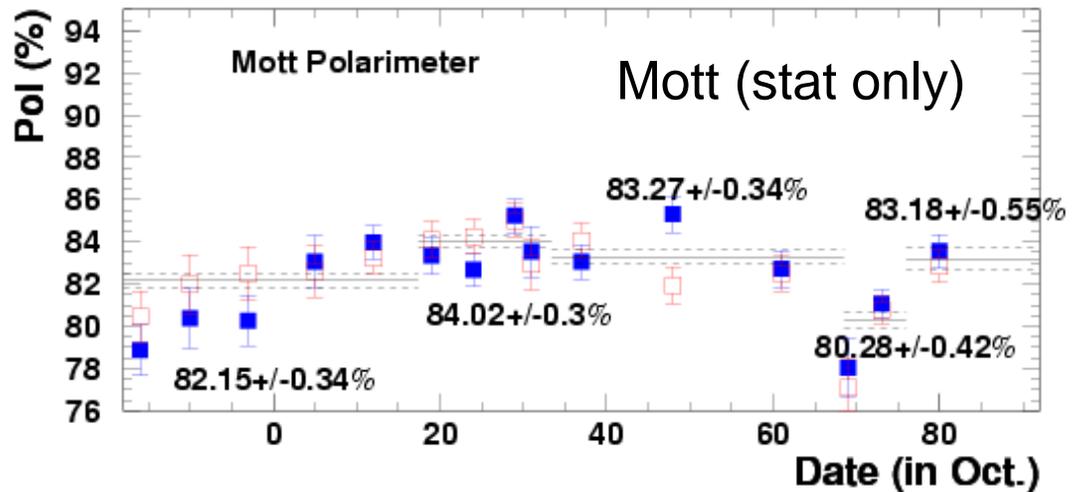
- Relatively good agreement between Hall C Møller and Mott and between Hall C Møller and Compton
- Hall A results are “online” only even though I show 1% syst.  
→ Compton takes significant offline analysis

# Comparisons During Fall 2006

Fall 2006, CEBAF was at “magic energy”  
 → allows longitudinal polarization to all halls for some passes

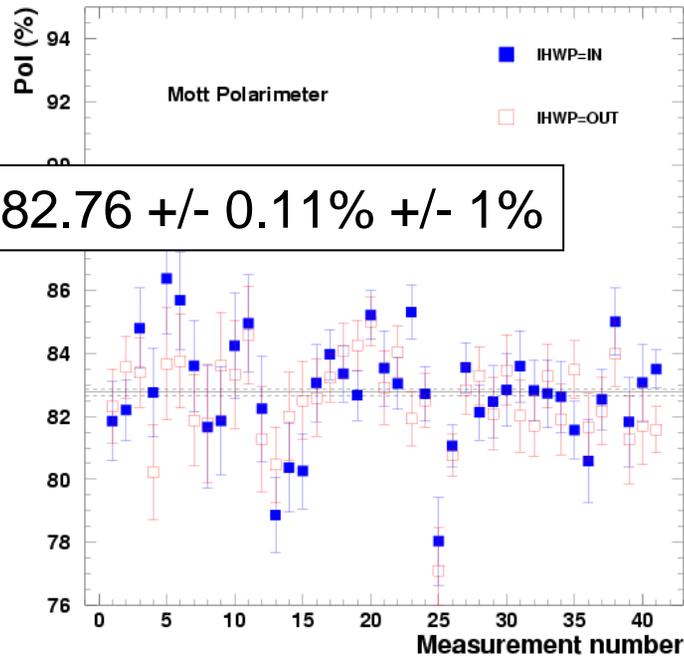
During this period A, B, C agreement quite good  
 → some unexplained variation in B measurements

Hall C Møller (stat only)  
 Hall A Møller (+2% sys.)

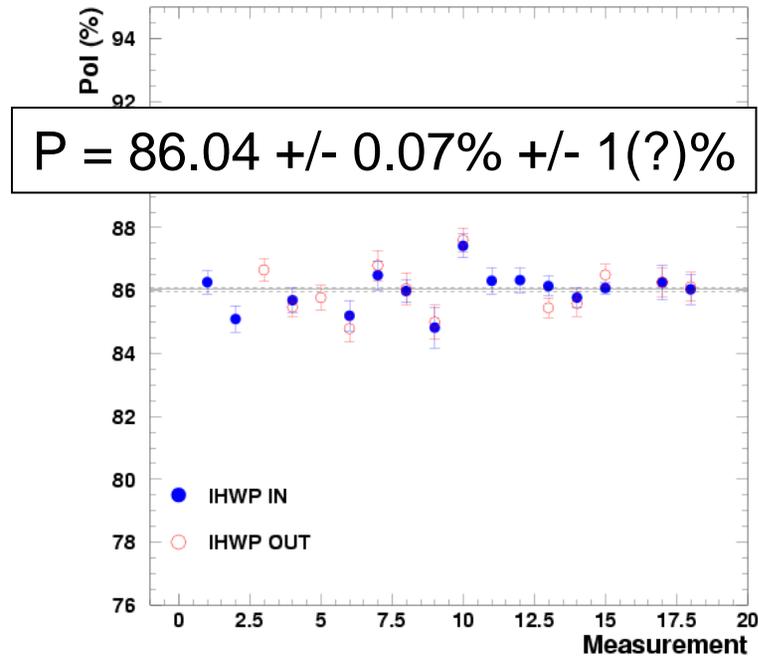


# Hall C Møller and Mott Discrepancy

## Mott Polarimeter



## Hall C Møller



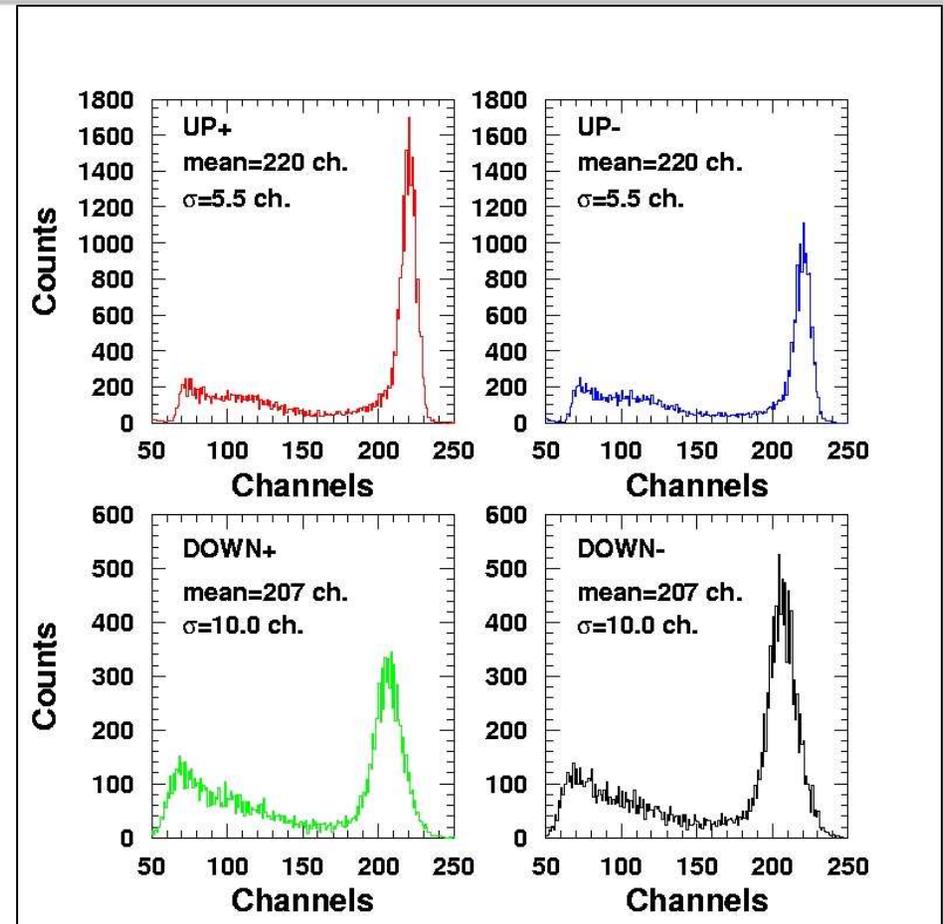
- Historically, Hall C Møller and Mott have agreed to 1.5-2%
- Measurements made during G0 Backangle indicate this difference has grown  $\rightarrow$  4%!

# Mott Detector Resolution

Resolution of DOWN detector in Mott polarimeter has worsened

→ Background under elastic peak may no longer be identical between UP and DOWN detectors

→ Simulation of background + resolution effects increases polarization extracted by Mott



Mott/Møller discrepancy: 4% → 2.5%

# Polarimetry Goals

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- Halls A and C are the primary drivers for precision polarimetry
- Hall C has “highest precision” polarimeter at JLab
  - $dP/P \sim 0.5\%^*$
  - Limited to low currents, destructive measurement
- Hall A the only experimental Hall with 2 electron beam polarimeters
  - Hall A Compton:  $dP/P = 1\%$  for  $E_{\text{beam}} > 3 \text{ GeV}$
  - Hall A Møller:  $dP/P = 2\text{-}3\%$ , limited by knowledge of target polarization
- ***Both Halls would like 2 polarimeters capable of polarization measurements yielding  $dP/P=1\%$***

# Polarimetry Upgrade Plans

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- Hall A
  - Upgrade Møller to use saturated foil target like Hall C
  - *Pursue extending reach of Møller to high currents*
  - Upgrade Compton to achieve high precision at low energy
- Hall C
  - *Pursue extending reach of Møller to high currents*
  - Build new Compton polarimeter to allow continuous non-destructive polarization measurements

# Møller Polarimetry Improvements

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- Primary limitation of Hall A Møller polarimeter is knowledge of target foil polarization
  - For PREX, Hall A will use old Hall C superconducting solenoid to make use of saturated foil technique
  - Identical targets in Halls A and C will allow direct comparison of *other* polarimeter systematics (Levchuk effect, etc.)
- Limitation in both Halls A *and* C is maximum current on foil target

$$Is P_e @ 2 \mu A = P_e @ 100 \mu A ?$$

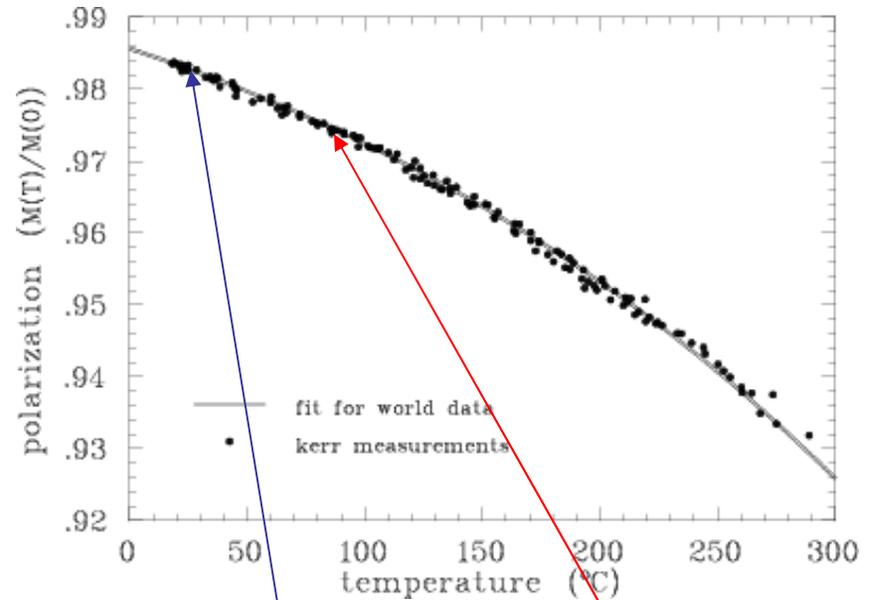
# Hall C Møller at High Beam Currents

In general Møller limited to low beam currents to avoid foil depolarization

Since 2003, have been pursuing tests with a fast “kicker” magnet and various targets

Minimizing time of beam on target and allowing enough time between “kicks” will mitigate foil heating

Fe Foil Depolarization

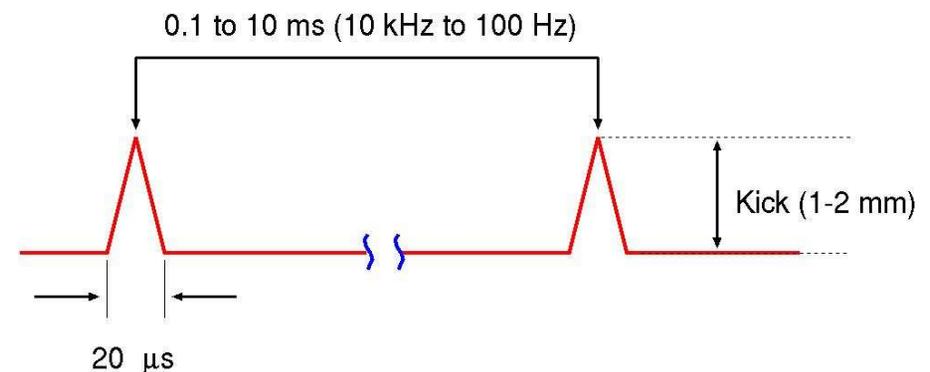
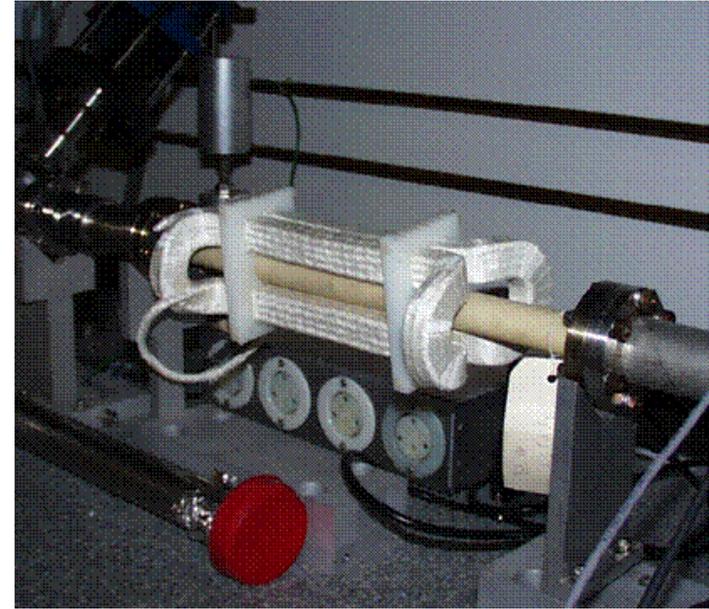


$\Delta P \sim 1\%$  for  
 $\Delta T \sim 60-70$  deg.

Operating Temp.

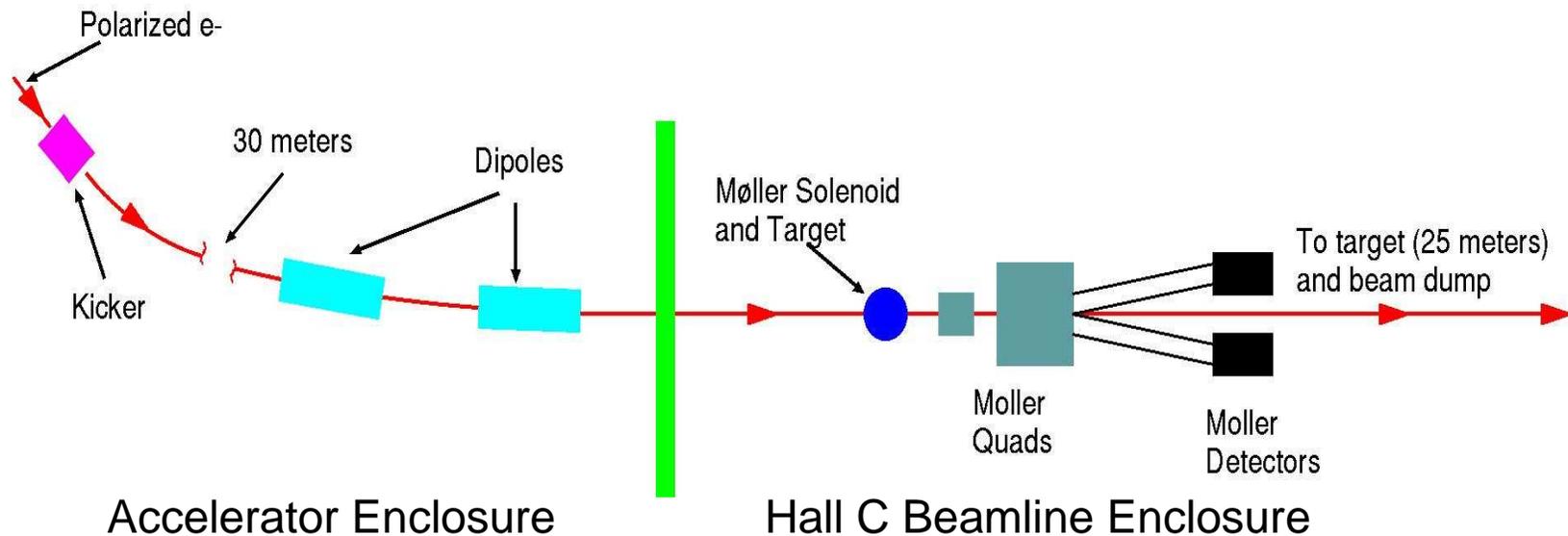
# Kicker Magnet for High Current Møller Polarimetry

- We can overcome target heating effects by using a **fast kicker magnet** to scan the electron beam across an iron wire or strip target
- Kicker needs to move beam quickly and at low duty cycle to **minimize** time on iron target and **beam heating**
- First generation kicker was installed in Fall 2003 (built by Chen Yan, Hall C)



# Kicker + Møller Layout

- Kicker located upstream of Møller target in Hall C beam transport arc
- Beam excursion ~ 1-2 mm at target
- The kick angle is small and the beam optics are configured to allow beam to continue cleanly to the dump



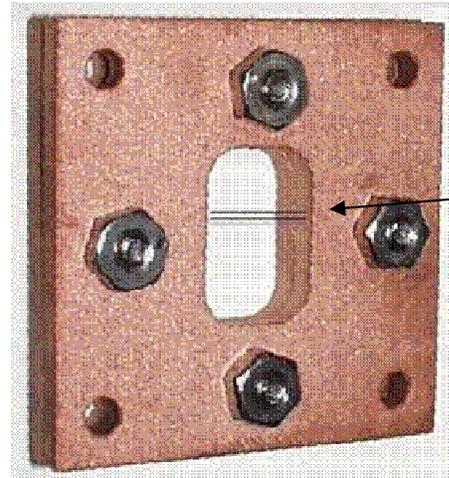
# Kicker and Iron Wire Target

- Initial tests with kicker and an **iron wire target** were performed in Dec. 2003
- Many useful lessons learned
  - 25  $\mu\text{m}$  wires too thick
  - Large instantaneous rate gave large rate of random coincidences

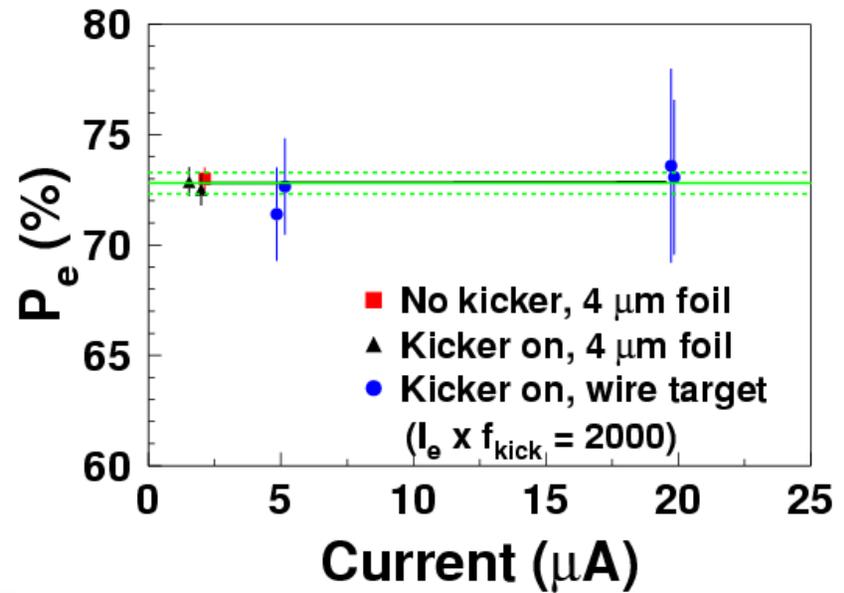
$N_{\text{coincidence}} \sim \text{target thickness}$

$N_{\text{random}} \sim (\text{target thickness})^2$

- Nonetheless, we were able to make measurements at currents up to 20  $\mu\text{A}$  (large uncertainties from large random rates)

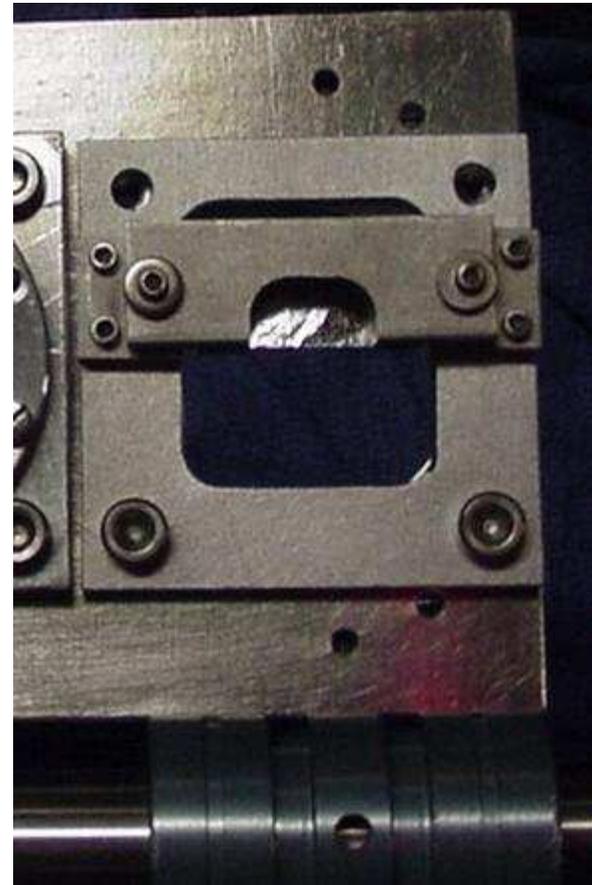


Target built by  
Dave Meekins  
JLab Target Group



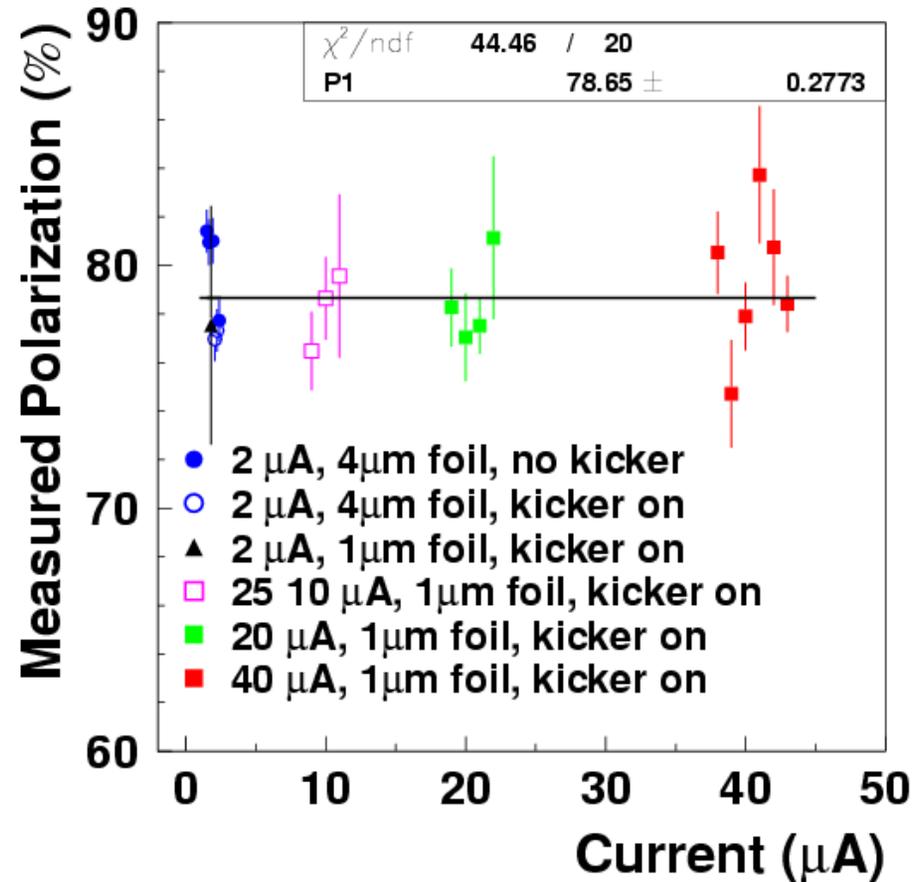
# Tests With a 1 $\mu\text{m}$ “Strip” Target

- The only way to keep random coincidences at an acceptable level is to reduce the instantaneous rate
- This can be achieved with a 1 mm foil
  - $N_{\text{real}}/N_{\text{random}} \approx 10$  at 200  $\mu\text{A}$  at a few GeV
- Replaced iron wire target with a 1  $\mu\text{m}$  thick iron “strip” target
- Conducted more tests with this target and slightly upgraded kicker in December 2004



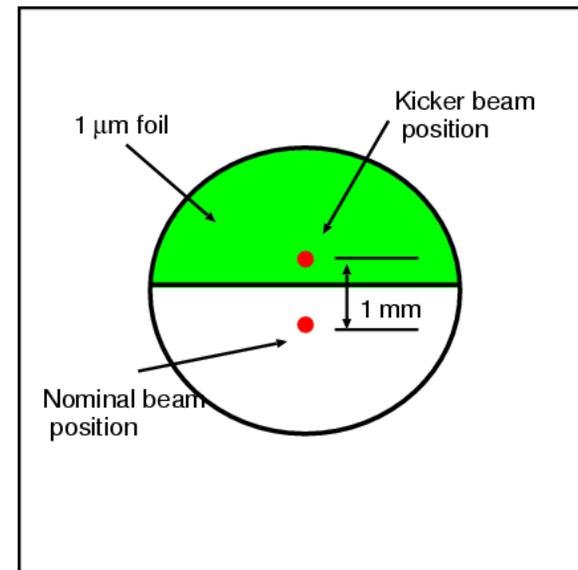
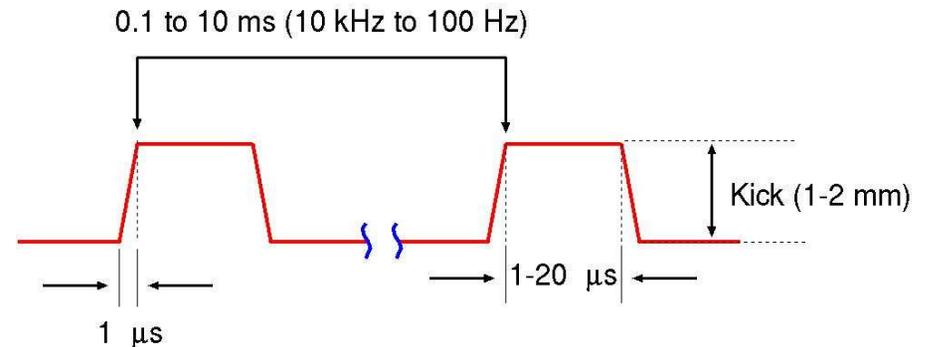
# December 2004 Kicker Test Results

- Short test – no time to optimize polarized source
  - Tests cannot be used to prove 1% precision
- Took measurements up to 40  $\mu\text{A}$ 
  - Ion chamber trips prevented us from running at higher currents
  - Lesson learned: need a beam tune that includes focus at Møller target AND downstream
- Demonstrated ability to make measurements at high currents – good proof of principle



# Optimized Kicker with “Half-Target”

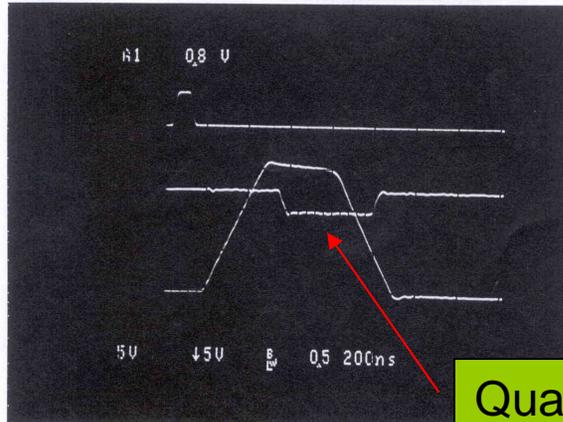
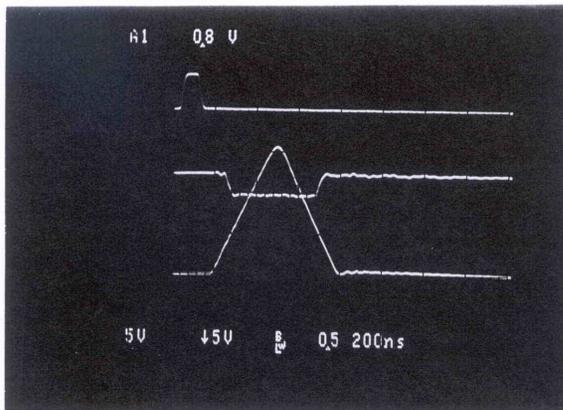
- The ideal kicker would allow the beam to dwell at a certain point on the target for a few ms rather than continuously move across the foil
- To reach the very highest currents, the kick duration must be as small as  $2 \mu\text{s}$  to keep target heating effects small
- The  $1 \mu\text{m}$  target is crucial – we need to improve the mounting scheme to avoid wrinkles and deformations



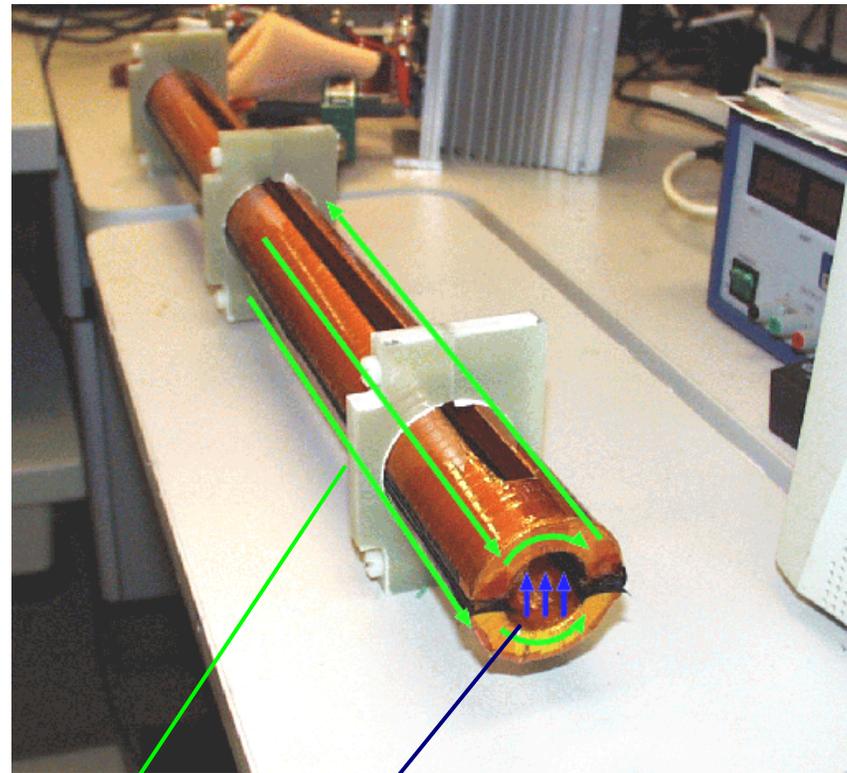
# Kicker R&D

## Kicker Current Waveform from Pearson Probe

Base width ~ 600 ns, Reputation ~ 5 kHz,  $I_{\text{peak}} \sim 100$  A



Quasi-flat top  
kicker interval



Current flow

Magnetic field

“Two turn” kicker –  
2  $\mu$ s total dwell time!

# Møller Polarimetry Using “Pulsed” Beam

- Use electron beam in “pulsed” mode
  - 0.1-1  $\mu\text{s}$  pulses at 30 to 120 Hz
  - Low average current, but for the duration of the pulse, same current as experiment conditions (10s of  $\mu\text{A}$ )
- Using a raster (25 kHz) to blow up the effective beam size, target heating can be kept at acceptable levels

Target Heating vs. Time for one beam pulse

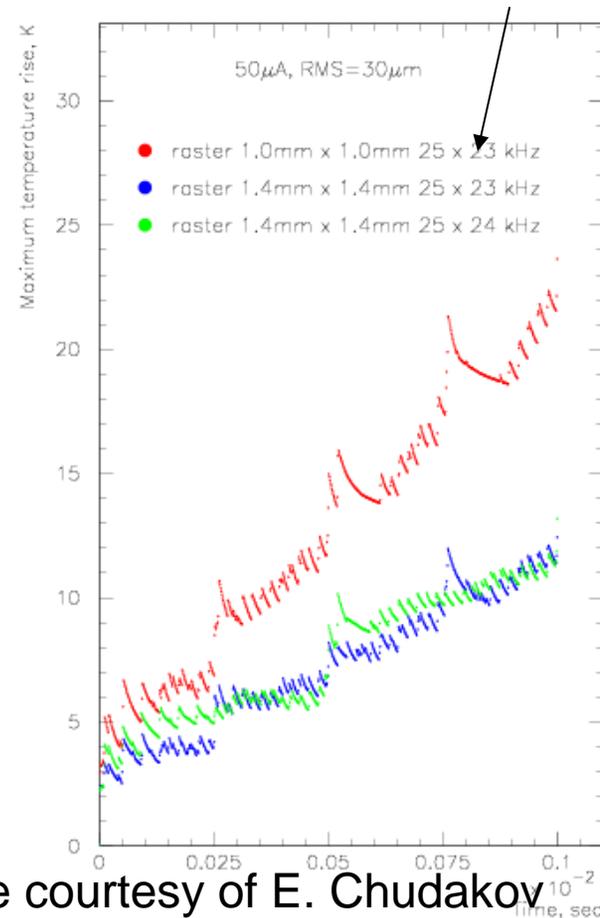


Figure courtesy of E. Chudakov

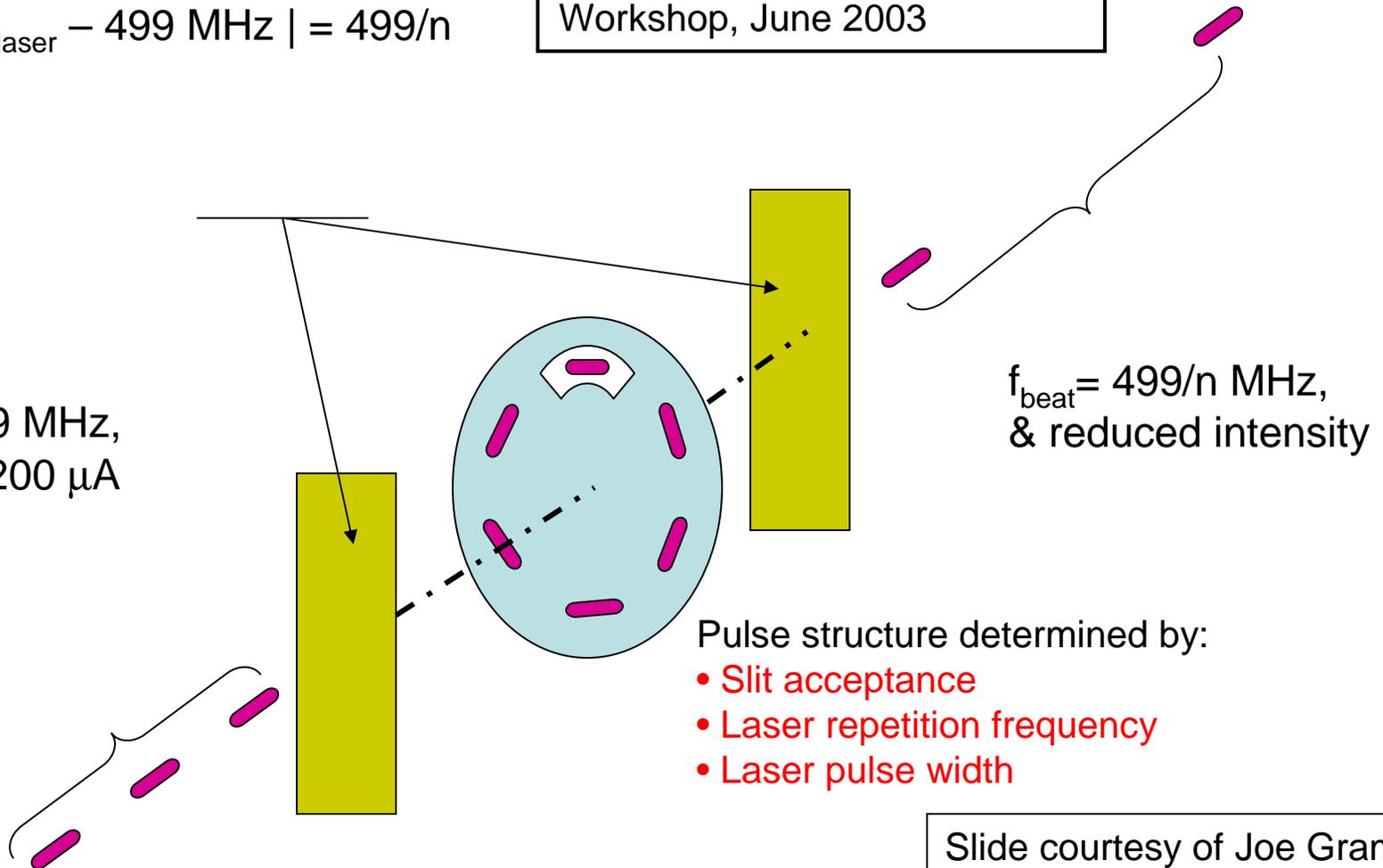
# Generation of 31 MHz Beam Structure

## RF beam & chopping

$$f_{\text{beat}} = |f_{\text{laser}} - 499 \text{ MHz}| = 499/n$$

Proposed at JLAB Precision  
Electron Beam Polarimetry  
Workshop, June 2003

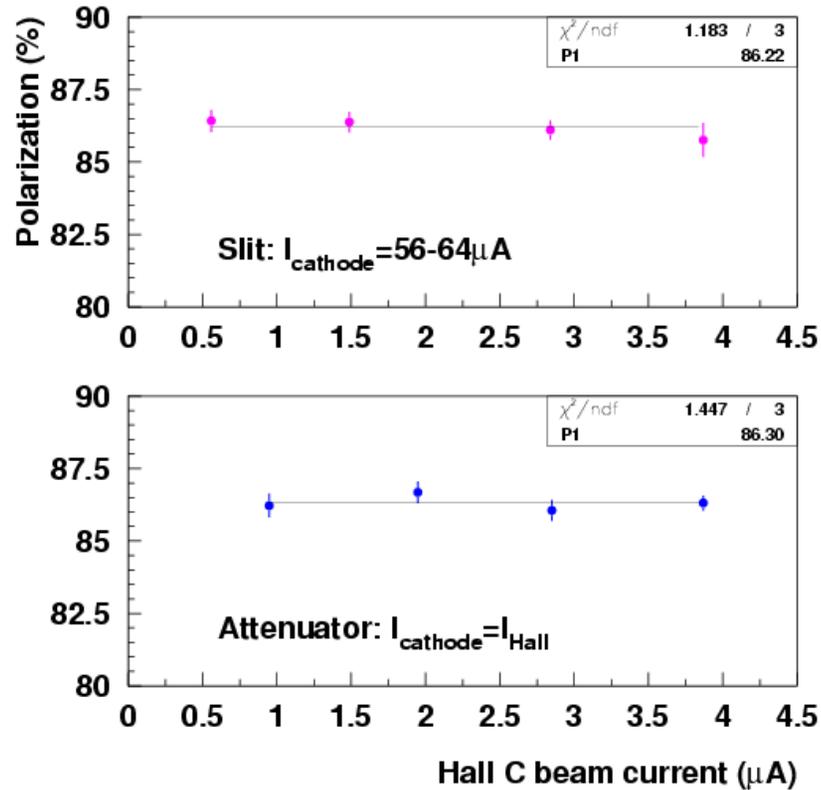
$f_{\text{laser}} \neq 499 \text{ MHz}$ ,  
but  $I_{\text{ave}} = 200 \mu\text{A}$



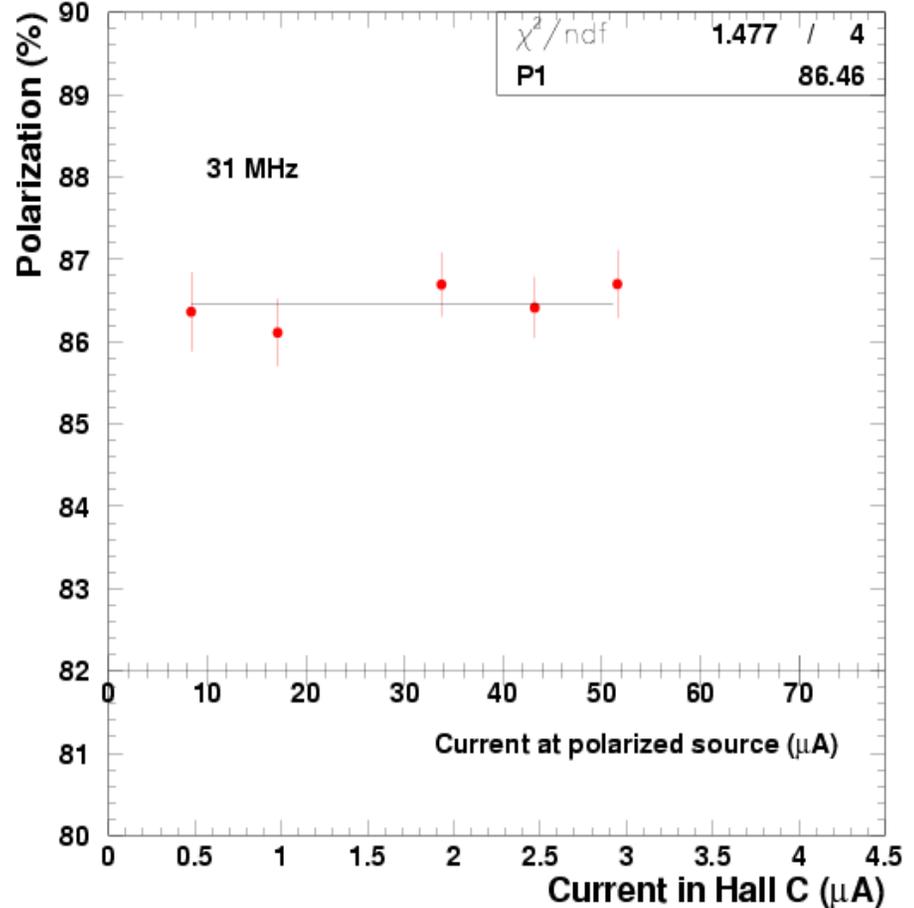
Slide courtesy of Joe Grames

# 31 MHz Test in Hall C

499 MHz – current reduced using slit or attenuator



31 MHz



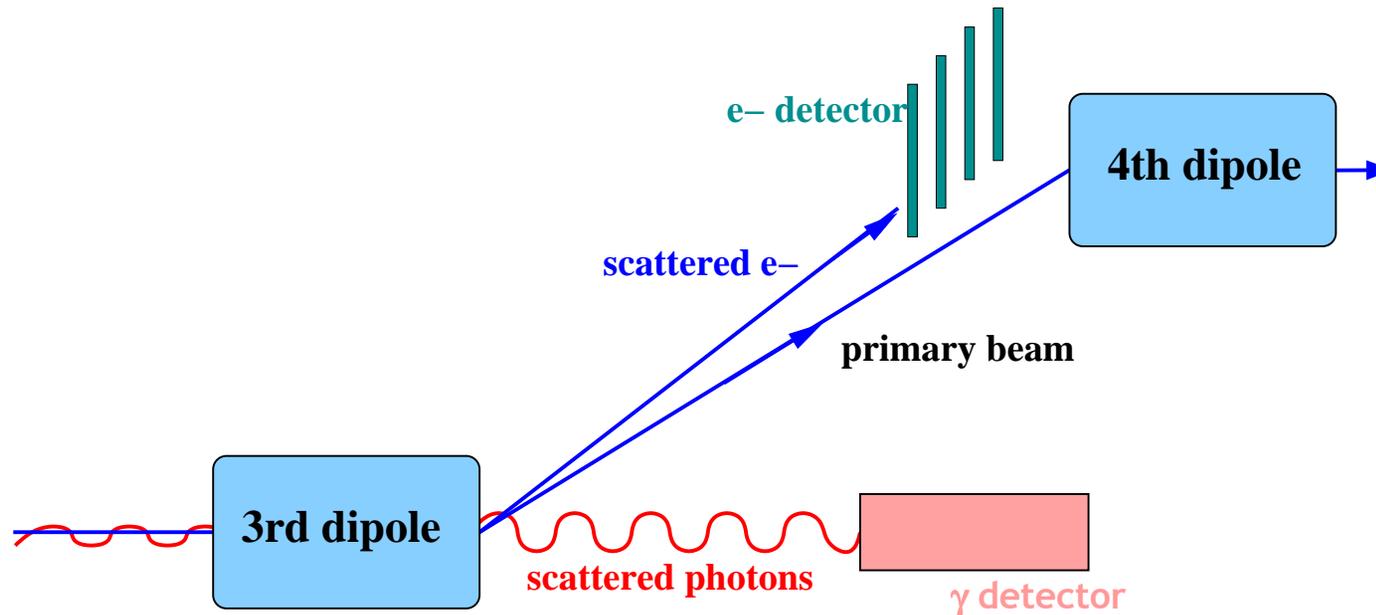
Error bars statistical only

# Compton Polarimetry Plans

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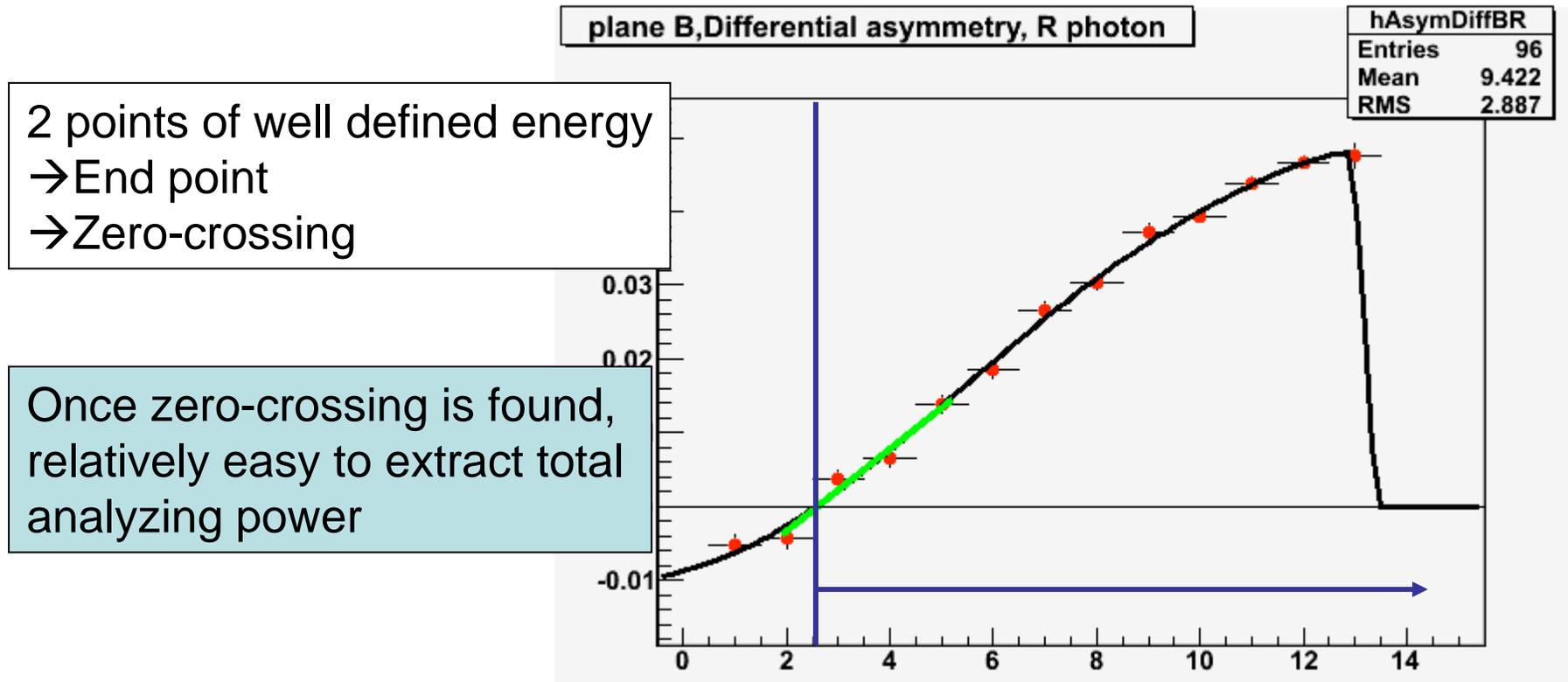
- Hall C has no Compton Polarimeter
  - Previous experience (G0, GEp) has shown us it is at best difficult to track the polarization over time using “periodic” measurements
  - To achieve  $dP/P$  of 1% *for the experiment*, a continuous measurement is required
- Hall A has a Compton Polarimeter that quotes  $dP/P=1\%$  at a few GeV
  - For PREX, 1% is required down to 850 MeV
  - There are fundamental issues that limit the systematic error of the existing device at low energy
  - Hall A will upgrade their Compton to achieve  $dP/P=1\%$  down to 850 MeV

# Hall A Compton Polarimeter



Hall A Compton detects scattered electron and backscattered photon simultaneously  
→ Uses high gain IR cavity to increase luminosity  
→ Best precision typically achieved using electron detector analysis

# Electron Detector



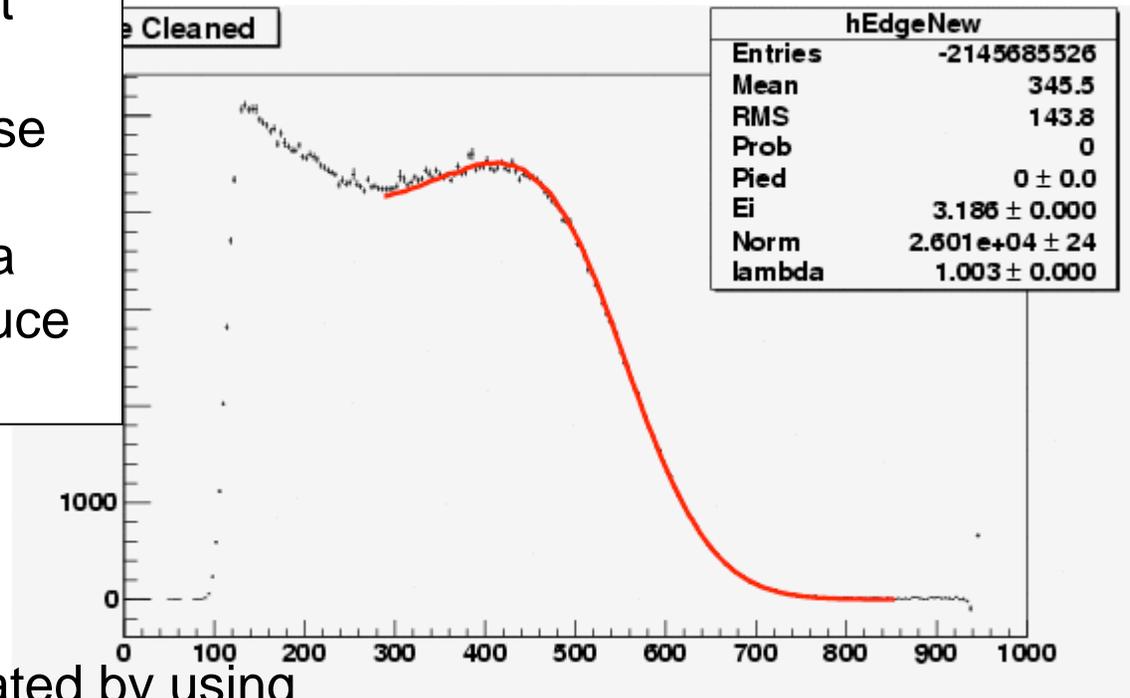
- At low energy – zero-crossing not in the acceptance
- This can be fixed by switching from IR to green laser
  - Green laser also has the advantage of yielding larger asymmetry
  - Upgrade will also increase granularity to improve resolution

# Photon Detector

Extracting analyzing power for photon detector more difficult

→ Need to know the “response function”,  $R(ADC, k)$

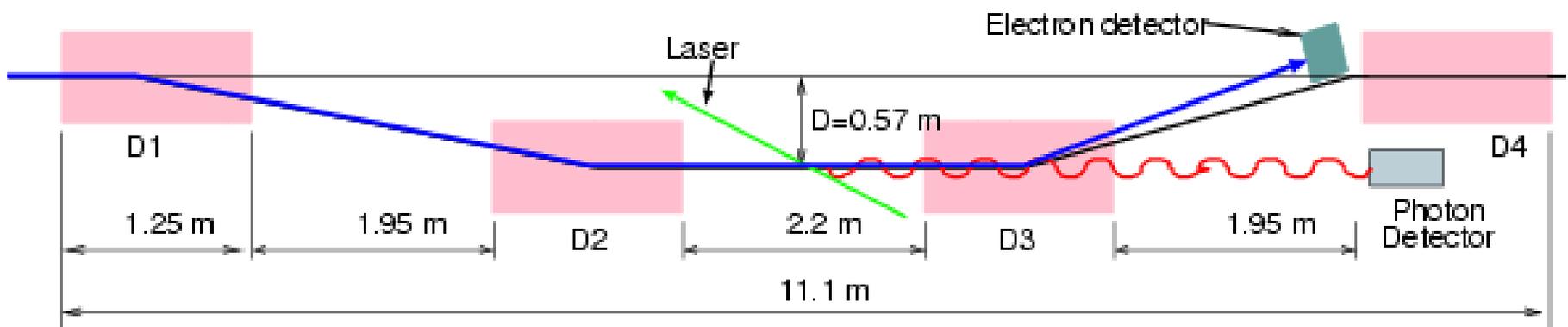
→  $R(ADC, k)$  = probability of a photon of energy “ $k$ ” to produce a signal in channel “ADC”



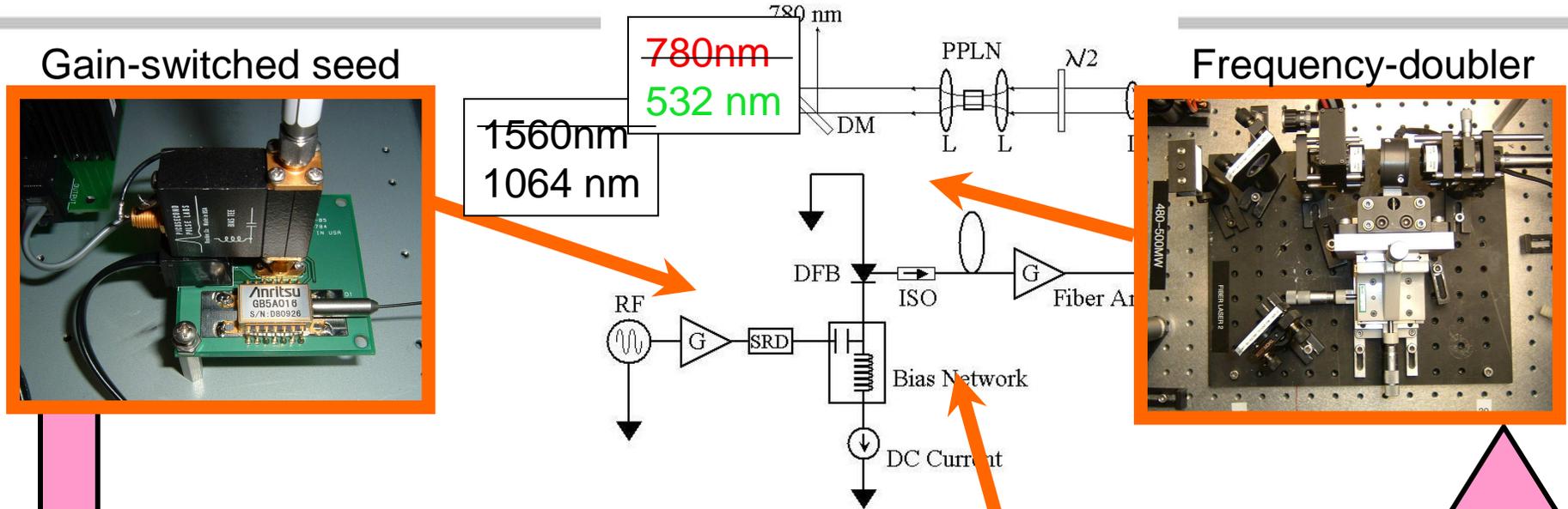
This can be somewhat alleviated by using “integrating” method → less sensitive to thresholds and absolute response  
→ Integrating method (in combination with new photon detector) will be pursued for PREX

# Hall C Compton Polarimeter

- The Hall C Compton polarimeter will be very similar to the Hall A Compton with some differences in the details
  - Chicane  $\rightarrow$  vertical drop = 57 cm (30 cm in Hall A)  $\rightarrow$  larger separation between scattered electron and beam
  - Electron detector  $\rightarrow$  diamond strip detector with 200  $\mu\text{m}$  pitch (silicon strip in Hall A)
  - Photon detector  $\rightarrow$  Lead-tungstate for now
  - RF pulsed 20 W green laser (high gain cavity in Hall A)



# Fiber-Based Drive Laser



ErYb-doped fiber amplifier

# Luminosity from Fiber Laser

- Average power from fiber laser modest (20 W) → Hall A will have high gain cavity with 1 kW of green power
- How can we live with this relatively low laser power?
  - For laser pulsed at electron beam repetition rate (499 MHz) and comparable pulse width (on the order of ps), the luminosity is increased by a factor:

$$\frac{L_{pulsed}}{L_{CW}} \approx \frac{c}{f \sqrt{2\pi}} \left( \sqrt{\sigma_{c\tau, laser}^2 + \sigma_{c\tau, e}^2 + \frac{1}{\sin^2(\alpha/2)} (\sigma_e^2 + \sigma_{laser}^2)} \right)^{-1}$$

- For typical JLab parameters, this yields about a factor of 20 improvement in luminosity for  $\alpha = 20$  mrad

**20 W RF pulsed laser ~ 400 W CW laser**

# Hall C Compton Project

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- Laser → EGG+Shukui Zhang+UVa+Hall C
- Electron Detector → U. Winnipeg + Miss. State
- Photon Detector → Yerevan + Hampton U.
- Dipole chicane → MIT-Bates
- Beamline overhaul → JLAB

***Everything must be ready at to be installed simultaneously with  $Q_{weak}$  → Fall 2009!***

Systematic goal = 1%: this will take significant effort and time. Initially we will cross-calibrate to Møller

# Polarimetry at JLab

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- Measuring the electron beam polarization at JLab is an industry with many contributors from EGG, and the experimental Halls
- Availability of multiple polarimeters using multiple techniques at JLab is a real advantage and strength
- To achieve 1% (and better) polarimetry, 2 things are required
  - More time: Experimenters want 1%, but they don't want to pay the price. It is hard to get beam time for systematic tests.
  - Better coordination between halls: As hard as it is to get time in 1 hall for polarimetry tests, it is harder to get time for 2 halls! (at the same time)