NUMERICAL STUDIES OF NON-LINEAR DYNAMICS IN BEP

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Abstract

An analysis of the dependence of experimental captured positron current data from the booster storage ring BEP (VEPP-2000 facility, BINP, Russia) on the working point position on the frequency map has uncovered a great number of different non-linear resonances. The number of captured positrons after a single injection is observed to be much less than the expected value. It is anticipated that the high degree of symmetry in the magnet system of BEP, however, should lead to the suppression of such resonances. To study this discrepancy, numerical simulations of positron beam movement under different perturbations to account for potential errors in magnetic field gradient of non-linear elements and errors in their angular location are used. The findings of this research provide qualitative explanations of the experimental work diagram and answers to two main questions, specifically “Why in the absence of skew-sextupoles in structure and small coupling are strong skew-sextupole resonances observed?” and “Why skew-sextupole resonances are stronger than sextupole ones of the same harmonic?”.

INTRODUCTION

The VEPP-2000 collider (BINP SB RAS, Russia) with $2 \times 1$ GeV design energy and $10^{32}$ cm$^2$sec$^{-1}$ design luminosity has been built for the studying of a hadron production in $e^+e^-$ collisions. Besides particle physics experiments, it has examined the round beam concept which ensures the required luminosity.

At this time the beam reproduction is realized by the scheme presented in a Fig. 1. The conversion target is placed in B-3M — BEP bypass channel for positron production. The theoretical value of the captured current of an $e^+$ beam should be around 600 $\mu$A with 1 A current of the $e^-$ beam ejected from B-3M. In reality the single capture doesn’t exceed 100 $\mu$A, which makes the operation of whole facility more complicated. In the future the injection will be from VEPP-5 facility which is under construction.

Research of the positron capture efficiency depending on beam’s position on a tune diagram shows the presence of a large number of non-linear resonances the source of which is not clear (e.g. skew-sextupole resonances). The results of this paper explain the source of this resonances and their strength.

BEP OPTICS

The booster synchrotron BEP consists of 12 the same identical periods composed of $30^\circ$ bending magnet and doublet of D- and F-lens. Structural functions and the arrangement of magnet optics are shown in the Fig. 2.

![Figure 2: Structural functions and the magnet optics arrangement of one period of BEP.](image)

Table 1: BEP Parameters.

<table>
<thead>
<tr>
<th>Physical Quantity</th>
<th>Value</th>
</tr>
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<tbody>
<tr>
<td>$\Pi$, perimeter</td>
<td>22.35 m</td>
</tr>
<tr>
<td>$H$, field in bending magnet</td>
<td>0.313 T</td>
</tr>
<tr>
<td>$G_D$, gradient in D-lens</td>
<td>-63 mT/cm</td>
</tr>
<tr>
<td>$G_F$, gradient in F-lens</td>
<td>+41 mT/cm</td>
</tr>
<tr>
<td>$U_0$, RF voltage</td>
<td>30 kV</td>
</tr>
<tr>
<td>$\omega_0$, RF frequency</td>
<td>26.83 MHz</td>
</tr>
<tr>
<td>$q$, multiplicity of RF harmonic</td>
<td>2</td>
</tr>
<tr>
<td>$\alpha$, orbit compaction factor</td>
<td>0.05</td>
</tr>
<tr>
<td>$E$, injection energy</td>
<td>120 MeV</td>
</tr>
<tr>
<td>$\nu_x$, x-betatron frequency</td>
<td>3.46</td>
</tr>
<tr>
<td>$\nu_z$, z-betatron frequency</td>
<td>3.21</td>
</tr>
<tr>
<td>$\nu_s$, synchrotron frequency</td>
<td>0.001</td>
</tr>
<tr>
<td>$W$, radiation losses</td>
<td>15 eV/turn</td>
</tr>
<tr>
<td>$\tau_x$, time of damping</td>
<td>1.3 sec</td>
</tr>
<tr>
<td>$\tau_z$, time of damping</td>
<td>1.2 sec</td>
</tr>
<tr>
<td>$\delta E/E$, energy deviation</td>
<td>$\pm 3%$</td>
</tr>
</tbody>
</table>

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Further to main optics the following corrections are added to the structure of BEP:

- Additional winding in D- and F-lenses for the quadrupole gradient correction of betatron frequencies variation.
- Dipole correctors in D-lenses and bending magnet for the orbit correction.
- Sextupole corrections for chromatism compensation; main sextupoles \( S_D \) and \( S_F \) are contained in quadrupole lenses, and additional sextupoles \( S_X \) and \( S_Z \) are located after F- and D-lenses respectively.

**EXPERIMENTAL DATA**

**Tune Diagram and Resonances**

In addition to integer and linear various non-linear resonances can occur in a system due to multipole fields. In general they are given by:

\[
m_x \nu_x + m_z \nu_z = n
\]

where \( m = |m_x| + |m_z| \) — resonance order and \( n \) — perturbation harmonic \( \# \).

![Figure 3: Tune diagram and resonance lines.](image)

The observed resonances are presented below in descending order of its strength:

- Sum and parametric skew-sextupole resonances relating to 10-th harmonic of perturbation:
  \[ 2\nu_x + \nu_z = 10 \text{ and } 3\nu_z = 10. \]
- Sextupole resonances on the same harmonic:
  \[ \nu_x + 2\nu_z = 10 \text{ and } 3\nu_z = 10. \]
- Difference sextupole resonance relating to \( n = 3 \):
  \[ -\nu_x + 2\nu_z = 3. \]
- Skew-sextupole resonances in next order of perturbation theory \( (r = 2) \):
  \[ 3\nu_x + \nu_z = 13, \quad 3\nu_z - \nu_x = 7, \quad \nu_x + 3\nu_z = 13, \quad 3\nu_z - \nu_x = 6. \]

Therefore two main questions appear: “Why in the absence of skew-sextupoles in structure and small coupling are strong skew-sextupole resonances observed?” and “Why skew-sextupole resonances are stronger than sextupole ones of the same harmonic?”

**RESULTS AND DISCUSSION**

In an attempt to explain experimental data several simulations of \( e^+ \) capture efficiency has been done. In the first approximation for unperturbed system the distributed sextupole gradient can be represented as a Dirac comb:

\[
G = \sum_{k=0}^{11} \left[ G_F \delta \left( \frac{\theta - 2\pi k}{12} \right) - G_D \delta \left( \frac{\theta - 2\pi k - \theta_{FD}}{12} \right) \right]
\]
of skew-sextupole resonances should be greater than sextupole ones for the harmonics of the perturbation with # different from 12k.

The only errors in magnetic field gradient $\sigma_{F,D}$ can not produce skew-sextupole resonances, but they can contribute to them due to additional gradient variation in a case with $\sigma_\varphi \neq 0$ (Fig. 7).

The coupling inclusion leads to the amplification of straight- and skew-sextupole resonances, parametric half-integer resonance $2\nu_\varphi = 7$ appearance (which presents in experimental data) and several stop-bands formation (Fig. 8).

REFERENCES
