

Electromagnetism I

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Three Dimensions

Define 3-D delta function by multiplication

$$
\delta^{3}(\mathbf{r}) = \delta(x)\delta(y)\delta(z) \quad \mathbf{r} = x\hat{x} + y\hat{y} + z\hat{z}
$$

By repeated integrations

egrations

\n
$$
\iiint \delta^3(\mathbf{r})dxdydz = \iiint \delta(x)\delta(y)\delta(z)dxdydz
$$
\n
$$
= 1 \iiint \delta(x)\delta(y)dxdy = 1 \iiint \delta(x)dx = 1
$$

Main properties

$$
\iiint f(x, y, z) \delta^{3}(x - a) dx dy dz = f(a_{x}, a_{y}, a_{z})
$$

space containing

Three Dimensions
\n3-D delta function by multiplication
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$$
\delta^3(r) = \delta(x)\delta(y)\delta(z) \quad r = x\hat{x} + y\hat{y} + z\hat{z}
$$
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\neated integrations
\n
$$
\iiint_{\delta} \delta^3(r) dxdydz = \iiint_{\delta} \delta(x)\delta(y)\delta(z) dxdydz
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= i\iiint_{\delta} \delta(x)\delta(y) dxdy = i\int_{\delta} \delta(x)dx = 1
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$$
\iiint_{\delta(\text{conmiting})} f(x, y, z) \delta^3(x-a) dxdydz = f(a_x, a_y, a_z)
$$
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$$
\nabla \cdot \frac{x-a}{|x-a|^3} = 4\pi \delta^3(x-a) \qquad \nabla^2 \left(\frac{1}{|x-a|}\right) = -4\pi \delta^3(x-a)
$$
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\text{Thomas Jefferson National Accelerator Facility}
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\text{Thus, a 25525 Tall 2024}
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\nThus, a 25525 Tall 2024

Helmholtz Theorem

Any vector field with appropriate boundary conditions (vanishing fast enough at infinity works) can be uniquely decomposed into a part with no curl (irrotational) and a part with no divergence (solenoidal). Conversely, if the divergence and curl of a vector field are given, along with boundary conditions, the vector field can be uniquely found.

In a single simply connected region of space, the irrotational part can be described by a scalar potential function ϕ and the solenoidal part by a vector potential function *A*

 $v = -\nabla \phi + \nabla \times A$,

a deep result known as Poincare's lemma.

Finding the Potentials

The scalar potential is found simply by performing line integrals, as in mechanics potential energy functions for conservative fields *x*

$$
\phi(x) = -\int_a^b v \cdot dl
$$

(*a* is a reference location where the potential vanishes).

Here is a method that allows you to find a vector potential for simple solenoidal vector fields (from a proof of Poincare's lemma) $\phi(x) = -\int_a^b v \cdot dl$
reference location where the potential vanishes).
a method that allows you to find a vector potential for
solenoidal vector fields (from a proof of Poincare's
)
 $\int_0^1 v_x(tx, ty, tz) t dt, V_y = \int_0^1 v_y(tx, ty, tz) t dt, V_z = \$

nma)
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$$
V_x = \int_0^1 v_x \left(tx, ty, tz \right) t dt, V_y = \int_0^1 v_y \left(tx, ty, tz \right) t dt, V_z = \int_0^1 v_z \left(tx, ty, tz \right) t dt
$$
\n
$$
A_x = V_y z - V_z y, A_y = V_z x - V_x y, A_z = V_x y - V_y x
$$

