

Electromagnetism I

G. A. Krafft, V. Ziemann Jefferson Lab Old Dominion University Lecture 4









Three Dimensions

Define 3-D delta function by multiplication

$$\delta^{3}(\mathbf{r}) = \delta(x)\delta(y)\delta(z) \quad \mathbf{r} = x\hat{x} + y\hat{y} + z\hat{z}$$

By repeated integrations

$$\iiint \delta^{3}(\mathbf{r}) dx dy dz = \iiint \delta(x) \delta(y) \delta(z) dx dy dz$$
$$= 1 \iint \delta(x) \delta(y) dx dy = 1 \iint \delta(x) dx = 1$$

Main properties

$$\iiint f(x, y, z)\delta^{3}(\boldsymbol{x} - \boldsymbol{a})dxdydz = f(a_{x}, a_{y}, a_{z})$$

space containing **a**

$$\nabla \cdot \frac{\boldsymbol{x} - \boldsymbol{a}}{|\boldsymbol{x} - \boldsymbol{a}|^3} = 4\pi\delta^3(\boldsymbol{x} - \boldsymbol{a}) \qquad \nabla^2 \left(\frac{1}{|\boldsymbol{x} - \boldsymbol{a}|}\right) = -4\pi\delta^3(\boldsymbol{x} - \boldsymbol{a})$$



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Helmholtz Theorem

Any vector field with appropriate boundary conditions (vanishing fast enough at infinity works) can be uniquely decomposed into a part with no curl (irrotational) and a part with no divergence (solenoidal). Conversely, if the divergence and curl of a vector field are given, along with boundary conditions, the vector field can be uniquely found.

In a single simply connected region of space, the irrotational part can be described by a scalar potential function ϕ and the solenoidal part by a vector potential function A

 $\mathbf{v} = -\nabla \phi + \nabla \times \mathbf{A},$

a deep result known as Poincare's lemma.





Finding the Potentials

The scalar potential is found simply by performing line integrals, as in mechanics potential energy functions for conservative fields $\phi(x) = -\int_{a}^{x} v \cdot dl$

(*a* is a reference location where the potential vanishes).

Here is a method that allows you to find a vector potential for simple solenoidal vector fields (from a proof of Poincare's lemma)

$$V_{x} = \int_{0}^{1} v_{x} (tx, ty, tz) t dt, V_{y} = \int_{0}^{1} v_{y} (tx, ty, tz) t dt, V_{z} = \int_{0}^{1} v_{z} (tx, ty, tz) t dt$$
$$A_{x} = V_{y} z - V_{z} y, A_{y} = V_{z} x - V_{x} y, A_{z} = V_{x} y - V_{y} x$$





