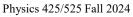


Electromagnetism I

G. A. Krafft, V. Ziemann Jefferson Lab Old Dominion University Lecture 23









Electrodynamic Equations till now

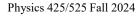
• So far

$$\nabla \cdot \boldsymbol{E} = \rho / \varepsilon_0$$
$$\nabla \times \boldsymbol{E} = -\partial \boldsymbol{B} / \partial t$$
$$\nabla \cdot \boldsymbol{B} = 0$$
$$\nabla \times \boldsymbol{B} = \mu_0 \boldsymbol{J}$$

• Take divergence of curl equations

 $\nabla \cdot \nabla \times \boldsymbol{E} = \nabla \cdot (-\partial \boldsymbol{B} / \partial t) = -\partial (\nabla \cdot \boldsymbol{B}) / \partial t$ $0 = 0 \qquad \text{OK!}$ $\nabla \cdot \nabla \times \boldsymbol{B} = \nabla \cdot (\mu_0 \boldsymbol{J}) = \mu_0 \nabla \cdot \boldsymbol{J}$ $0 = \nabla \cdot \boldsymbol{J} \qquad \text{Inconsistent with } \nabla \cdot \boldsymbol{J} + \partial \rho / \partial t = 0!$







Displacement Current (Density)

• Note

$$\frac{\partial}{\partial t} \left(\nabla \cdot \boldsymbol{E} \right) = \nabla \cdot \frac{\partial \boldsymbol{E}}{\partial t} = \frac{1}{\varepsilon_0} \frac{\partial \rho}{\partial t}$$

• To fix, need to add something to the **B** curl equation

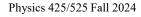
$$\nabla \times \boldsymbol{B} = \mu_0 \boldsymbol{J} + \mu_0 \varepsilon_0 \frac{\partial \boldsymbol{E}}{\partial t}$$
$$\nabla \cdot \nabla \times \boldsymbol{B} = \mu_0 \nabla \cdot \boldsymbol{J} + \mu_0 \varepsilon_0 \frac{\partial (\nabla \cdot \boldsymbol{E})}{\partial t} = 0$$

• In the added term,

$\varepsilon_0 \frac{\partial E}{\partial t}$

was called by Maxwell the displacement current (density)







Maxwell's Electrodynamics

• Equations, now including all time dependences

$$\nabla \cdot \boldsymbol{E} = \rho / \varepsilon_0$$

$$\nabla \times \boldsymbol{E} = -\partial \boldsymbol{B} / \partial t$$

$$\nabla \cdot \boldsymbol{B} = 0$$

$$\nabla \times \boldsymbol{B} = \mu_0 \boldsymbol{J} + \mu_0 \varepsilon_0 \partial \boldsymbol{E} / \partial t$$

• By Faraday's Law, a changing magnetic field generates and electric field. Now, due to the displacement current, *a changing electric field generates a magnetic field*. Next semester will show electromagnetic waves, including light, are possible solutions!





Magnetic Charges



• If there were magnetic charges, Maxwell's equations would have to be

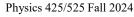
 $\nabla \cdot \boldsymbol{E} = \rho_e / \varepsilon_0$ $\nabla \times \boldsymbol{E} = -\mu_0 \boldsymbol{J}_m - \partial \boldsymbol{B} / \partial t$ $\nabla \cdot \boldsymbol{B} = \mu_0 \rho_m$ $\nabla \times \boldsymbol{B} = \mu_0 \boldsymbol{J}_e + \mu_0 \varepsilon_0 \partial \boldsymbol{E} / \partial t$

• Continuity for electric and magnetic charges

$$\nabla \cdot \boldsymbol{J}_{e} = -\partial \rho_{e} / \partial t \qquad \nabla \cdot \boldsymbol{J}_{m} = -\partial \rho_{m} / \partial t$$

• Sign of J_m term must be negative to ensure continuity







Force Law (Problem 7.40)

• Magnetic analogue of "Coulomb's Law" must be

$$F = \frac{\mu_0}{4\pi} \frac{q_{m1}q_{m2}}{|\mathbf{z}|^2} \hat{\mathbf{x}}$$
$$B = \frac{\mu_0}{4\pi} \frac{q_m}{|\mathbf{z}|^2} \hat{\mathbf{x}}$$

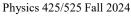
$$\boldsymbol{B} = \frac{\boldsymbol{\mu}_0}{4\pi} \frac{\boldsymbol{q}_m}{\left|\boldsymbol{z}\right|^2} \, \hat{\boldsymbol{z}}$$

$$\nabla \cdot \boldsymbol{B} = \frac{\mu_0}{4\pi} q_m \nabla \cdot \frac{\hat{\boldsymbol{x}}}{\left|\boldsymbol{z}\right|^2} = \mu_0 q_m \delta^3 \left(\boldsymbol{r} - \boldsymbol{r}'\right)$$

• And force law is (- sign from relativity argument!)

$$\boldsymbol{F} = q_m \left(\boldsymbol{B} - \frac{\boldsymbol{v} \times \boldsymbol{E}}{c^2} \right)$$









Maxwell's Equations with Matter

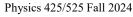
• When allow time-dependent fields, need to allow polarization and magnetization (and the bound charges and currents) to be time dependent

$$\rho_{b} = -\nabla \cdot \boldsymbol{P}$$
$$\frac{\partial \rho_{b}}{\partial t} = -\nabla \cdot \boldsymbol{J}_{p} = -\nabla \cdot \frac{\partial \boldsymbol{P}}{\partial t}$$
$$\boldsymbol{J}_{p} = \frac{\partial \boldsymbol{P}}{\partial t}$$

• A time dependent magnetization *M* adds no new source term to Maxwell Equations. Still

$$\boldsymbol{J}_b = \nabla \times \boldsymbol{M}$$









Gauss's Law (with materials)

• As before

$$\rho = \rho_f + \rho_b$$
$$= \rho_f - \nabla \cdot \boldsymbol{P}$$

• Gauss's Law

$$\nabla \cdot \boldsymbol{\varepsilon}_0 \boldsymbol{E} = \boldsymbol{\rho}_f - \nabla \cdot \boldsymbol{P}$$

• In terms of the electric displacement **D**

$$\nabla \cdot \boldsymbol{D} = \rho_f$$

• And for linear materials

$$\boldsymbol{D} = \boldsymbol{\varepsilon} \boldsymbol{E} = \boldsymbol{\varepsilon}_0 \left(1 + \boldsymbol{\chi}_e \right) \boldsymbol{E}$$





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Ampere's Law (with materials)

• Now

$$\boldsymbol{J} = \boldsymbol{J}_{f} + \boldsymbol{J}_{b} + \boldsymbol{J}_{p} = \boldsymbol{J}_{f} + \nabla \times \boldsymbol{M} + \frac{\partial \boldsymbol{P}}{\partial t}$$

• Ampere's Law is

$$\nabla \times \boldsymbol{B} = \mu_0 \boldsymbol{J}_f + \mu_0 \nabla \times \boldsymbol{M} + \mu_0 \frac{\partial \boldsymbol{P}}{\partial t} + \mu_0 \varepsilon_0 \frac{\partial \boldsymbol{E}}{\partial t}$$

• In terms of *H*

$$\nabla \times \left(\frac{\boldsymbol{B}}{\mu_0} - \boldsymbol{M}\right) = \nabla \times \boldsymbol{H} = \boldsymbol{J}_f + \frac{\partial \boldsymbol{D}}{\partial t}$$

• Linear materials

$$\boldsymbol{H} + \boldsymbol{\chi}_{m}\boldsymbol{H} = \frac{\boldsymbol{B}}{\mu_{0}} \rightarrow \boldsymbol{H} = \frac{\boldsymbol{B}}{\mu_{0}\left(1 + \boldsymbol{\chi}_{m}\right)} = \frac{\boldsymbol{B}}{\mu}$$



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Maxwell's Equations with Matter

• Full set

$$\nabla \cdot \boldsymbol{D} = \rho_f$$
$$\nabla \times \boldsymbol{E} = -\partial \boldsymbol{B} / \partial t$$
$$\nabla \cdot \boldsymbol{B} = 0$$
$$\nabla \times \boldsymbol{H} = \boldsymbol{J}_f + \partial \boldsymbol{D} / \partial t$$

- In general need a relationship (could be linear or nonlinear) relating **B** and **H** and **E** and **D** (constitutive relations)
- Need to apply proper boundary conditions to make solutions!

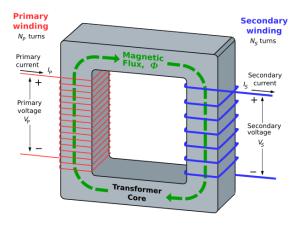




Example: Transformer



• Iron core linking different turns (ideally, no flux lost)



• Currents are sinusoids

$$\frac{\mathbf{Z}_{1}}{\mathbf{Z}_{2}} = \frac{n_{1}\dot{\Phi}_{1}}{n_{2}\dot{\Phi}_{2}} = \frac{n_{1}\dot{\Phi}_{1}}{n_{2}\dot{\Phi}_{1}} = \frac{n_{1}}{n_{2}}$$

• No power lost (high voltage, low current→low voltage high current)

$$\boldsymbol{\mathcal{E}}_{1}I_{1} = \boldsymbol{\mathcal{E}}_{2}I_{2} \rightarrow \frac{I_{1}}{I_{2}} = \frac{n_{2}}{n_{1}}$$



