



Physics 704/804 Electromagnetic Theory II

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Radiation From Relativistic Electrons



From discussion early in the course, in the Lorenz gauge the equation for the potentials is

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right] \phi = -\frac{\rho}{\epsilon_0}$$

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right] \vec{A} = -\mu_0 \vec{J}$$

The solution, using the retarded Green Function is

$$\phi(\vec{x}, t) = \frac{1}{4\pi\epsilon_0} \int d^3x' dt' \frac{\rho(\vec{x}', t')}{|\vec{x} - \vec{x}'|} \delta(|\vec{x} - \vec{x}'|/c - t + t')$$

$$\vec{A}(\vec{x}, t) = \frac{\mu_0}{4\pi} \int d^3x' dt' \frac{\vec{J}(\vec{x}', t')}{|\vec{x} - \vec{x}'|} \delta(|\vec{x} - \vec{x}'|/c - t + t')$$

Delta Function Representation



$$\phi(\vec{x}, t) = \frac{1}{8\pi^2 \epsilon_0} \int d^3x' dt' d\omega \frac{\rho(\vec{x}', t')}{|\vec{x} - \vec{x}'|} e^{i\omega[|\vec{x} - \vec{x}'|/c - (t - t')]} \quad (1)$$

$$\vec{A}(\vec{x}, t) = \frac{\mu_0}{8\pi^2} \int d^3x' dt' d\omega \frac{\vec{J}(\vec{x}', t')}{|\vec{x} - \vec{x}'|} e^{i\omega[|\vec{x} - \vec{x}'|/c - (t - t')]} \quad (2)$$

$$\rho(\vec{x}, t) = q\delta^3(\vec{x} - \vec{r}(t)) \quad \vec{J}(\vec{x}, t) = q\vec{v}(t)\delta^3(\vec{x} - \vec{r}(t)) \quad (3)$$

$$\phi(\vec{x}, t) = \frac{q}{8\pi^2 \epsilon_0} \int dt' d\omega \frac{1}{|\vec{x} - \vec{r}(t')|} e^{i\omega[|\vec{x} - \vec{r}(t')|/c - (t - t')]} \quad (4)$$

$$\vec{A}(\vec{x}, t) = \frac{q\mu_0}{8\pi^2} \int dt' d\omega \frac{\vec{v}(t')}{|\vec{x} - \vec{r}(t')|} e^{i\omega[|\vec{x} - \vec{r}(t')|/c - (t - t')]} \quad (5)$$

Lienard-Weichert Potentials



$$\phi(\vec{x}, t) = \frac{q}{4\pi\epsilon_0} \int dt' \frac{\delta(|\vec{x} - \vec{r}(t')|/c - (t - t'))}{|\vec{x} - \vec{r}(t')|}$$

$$\vec{A}(\vec{x}, t) = \frac{q\mu_0}{4\pi} \int dt' \frac{\vec{v}(t') \delta(|\vec{x} - \vec{r}(t')|/c - (t - t'))}{|\vec{x} - \vec{r}(t')|}$$

$$\phi(\vec{x}, t) = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{|\vec{x} - \vec{r}(t')|(1 - \hat{n} \cdot \vec{\beta}(t'))} \right)_{ret}$$

$$\vec{A}(\vec{x}, t) = \frac{q\mu_0}{4\pi} \left(\frac{\vec{v}(t')}{|\vec{x} - \vec{r}(t')|(1 - \hat{n} \cdot \vec{\beta}(t'))} \right)_{ret}$$

EM Field Radiated



$$\phi(\vec{x}, t) = \frac{q}{8\pi^2 \epsilon_0} \int dt' d\omega \frac{1}{|\vec{x} - \vec{r}(t')|} e^{i\omega[|\vec{x} - \vec{r}(t')|/c - (t - t')]} \quad (1)$$

$$\vec{A}(\vec{x}, t) = \frac{q \mu_0}{8\pi^2} \int dt' d\omega \frac{\vec{v}(t')}{|\vec{x} - \vec{r}(t')|} e^{i\omega[|\vec{x} - \vec{r}(t')|/c - (t - t')]} \quad (2)$$

$$\vec{E} = -\vec{\nabla} \phi - \frac{\partial \vec{A}}{\partial t} \quad \vec{B} = \vec{\nabla} \times \vec{A}$$

$$\vec{E} = \frac{q}{4\pi\epsilon_0} \left[\frac{\hat{n} - \vec{\beta}}{\gamma^2 (1 - \hat{n} \cdot \vec{\beta})^3 R^2} \right]_{ret} + \frac{q}{4\pi\epsilon_0 c} \left[\frac{\hat{n} \times \{(\hat{n} - \vec{\beta}) \times \dot{\vec{\beta}}\}}{(1 - \hat{n} \cdot \vec{\beta})^3 R} \right]_{ret} \quad (3)$$

$$\vec{B} = \hat{n} \times \vec{E} / c$$

$$\vec{\nabla} \frac{1}{|\vec{x} - \vec{r}(t')|} = -\frac{\hat{n}}{|\vec{x} - \vec{r}(t')|^2}$$

$$\vec{\nabla} |\vec{x} - \vec{r}(t')| = \hat{n}$$

$$\frac{d}{dt'} \frac{1}{|\vec{x} - \vec{r}(t')|} = \frac{\hat{n} \cdot \vec{\beta} c}{|\vec{x} - \vec{r}(t')|^2}$$

$$\frac{d\hat{n}}{dt'} = \frac{-d\vec{r} / dt' + \hat{n}(\hat{n} \cdot d\vec{r} / dt')}{|\vec{x} - \vec{r}(t')|} \quad \dots$$

$$-\nabla \phi(\vec{x}, t) = \frac{q}{8\pi^2 \epsilon_0} \int dt' d\omega \frac{\hat{n} \left(1 - i\omega |\vec{x} - \vec{r}(t')| / c\right)}{|\vec{x} - \vec{r}(t')|^2} e^{i\omega [|\vec{x} - \vec{r}(t')|/c - (t-t')]} \quad \dots$$

$$-\frac{\partial}{\partial t} \vec{A}(\vec{x}, t) = \frac{q \mu_0}{8\pi^2} \int dt' d\omega \frac{\vec{v}(t') i\omega}{|\vec{x} - \vec{r}(t')|} e^{i\omega [|\vec{x} - \vec{r}(t')|/c - (t-t')]} \quad \dots$$

$$\frac{d}{dt'} e^{i\omega [|\vec{x} - \vec{r}(t')|/c - (t-t')]} = i\omega \left(1 - \vec{\beta}(t') \cdot \hat{n}(t')\right) e^{i\omega [|\vec{x} - \vec{r}(t')|/c - (t-t')]} \quad \dots$$

$$\vec{E}(\vec{x}, t) = \frac{q}{8\pi^2 \epsilon_0} \int dt' d\omega e^{i\omega [|\vec{x} - \vec{r}(t')|/c - (t - t')]} \left[\frac{\hat{n}}{|\vec{x} - \vec{r}(t')|^2} + \frac{i\omega (\vec{\beta} - \hat{n})}{c |\vec{x} - \vec{r}(t')|} \right]$$

integrate by parts to get final result

$$\vec{E}(\vec{x}, t)_{vel} = \frac{q}{8\pi^2 \epsilon_0} \int dt' d\omega \frac{e^{i\omega [|\vec{x} - \vec{r}(t')|/c - (t - t')]} }{(1 - \vec{\beta} \cdot \hat{n})^2 |\vec{x} - \vec{r}(t')|^2}$$

$$\times \left[\begin{aligned} &\hat{n} \left(1 - 2\vec{\beta} \cdot \hat{n} + (\vec{\beta} \cdot \hat{n})^2 + \vec{\beta} \cdot \hat{n} - (\vec{\beta} \cdot \hat{n})^2 + \vec{\beta} \cdot \hat{n} - (\vec{\beta} \cdot \hat{n})^2 \right) \\ &- \beta^2 + (\vec{\beta} \cdot \hat{n})^2 \\ &- \vec{\beta} \left(1 - \vec{\beta} \cdot \hat{n} + \vec{\beta} \cdot \hat{n} - (\vec{\beta} \cdot \hat{n})^2 - \beta^2 + (\vec{\beta} \cdot \hat{n})^2 \right) \end{aligned} \right]$$

$$\begin{aligned}
 \vec{E}(\vec{x}, t)_{acc} &= \frac{q}{8\pi^2 \epsilon_0 c} \int dt' d\omega \frac{e^{i\omega[|\vec{x}-\vec{r}(t')|/c-(t-t')]}}{(1-\vec{\beta} \cdot \hat{n})^2 |\vec{x}-\vec{r}(t')|} \\
 &\times \begin{bmatrix} \hat{n}(\dot{\vec{\beta}} \cdot \hat{n}) \\ -\dot{\vec{\beta}}(1-\vec{\beta} \cdot \hat{n}) - \vec{\beta}(\dot{\vec{\beta}} \cdot \hat{n}) \end{bmatrix} \\
 &= \frac{q}{8\pi^2 \epsilon_0 c} \int dt' d\omega \frac{e^{i\omega[|\vec{x}-\vec{r}(t')|/c-(t-t')]}}{(1-\vec{\beta} \cdot \hat{n})^2 |\vec{x}-\vec{r}(t')|} \left[\hat{n} \times \left\{ (\hat{n} - \vec{\beta}) \times \dot{\vec{\beta}} \right\} \right]
 \end{aligned}$$

Larmor's Formula



For small velocities can neglect retardation

$$\vec{E}(\vec{x}, t)_{acc} = \frac{q}{4\pi\epsilon_0 c} \left[\hat{n} \times \left\{ \hat{n} \times \dot{\vec{\beta}} \right\} \right] / R$$

$$\frac{dP}{d\Omega} = \frac{q^2}{16\pi^2 \epsilon_0 \mu_0 c^3} \left| \left[\hat{n} \times \left\{ \hat{n} \times \dot{\vec{\beta}} \right\} \right] \right|^2$$

$$= \frac{q^2}{16\pi^2 \epsilon_0 c^3} \left| \dot{\vec{v}} \right|^2 \sin^2 \theta$$

$$P = \frac{q^2}{6\pi\epsilon_0 c^3} \left| \dot{\vec{v}} \right|^2$$

Lienard's Relativistic Generalization



For photon emission, the emitted energy increment dE and duration increment dt are time components of (different) 4-vectors. Therefore their ratio, the power emitted

$$P = \frac{dE}{dt}$$

is a Lorentz invariant! Must be of form

$$P = -\frac{q^2}{6\pi\varepsilon_0 m^2 c^3} \left(\frac{dp^\mu}{d\tau} \frac{dp_\mu}{d\tau} \right)$$

$$E = \gamma mc^2 \quad \vec{p} = \gamma m\vec{v} \quad d\tau = dt / \gamma$$

$$P = -\frac{q^2}{6\pi\varepsilon_0 m^2 c^3} \left(\frac{dp^\mu}{d\tau} \frac{dp_\mu}{d\tau} \right) = \frac{q^2}{6\pi\varepsilon_0 c} \gamma^6 \left(\left(\dot{\vec{\beta}} \right)^2 - \left(\vec{\beta} \times \dot{\vec{\beta}} \right)^2 \right)$$

Very strong energy dependence!

For Circular Motion

$$\vec{x}(t) = \rho [\cos \omega t \hat{x} + \sin \omega t \hat{y}]$$

$$\vec{v}(t) = \rho \omega [-\sin \omega t \hat{x} + \cos \omega t \hat{y}] \quad \dot{\vec{v}}(t) = -\omega^2 \vec{x}$$

$$\dot{\vec{\beta}}^2 = \left(\frac{\rho \omega^2}{c} \right)^2 \quad (\vec{\beta} \times \dot{\vec{\beta}})^2 = \left(\frac{\rho \omega}{c} \frac{\rho \omega^2}{c} \right)^2$$

$$P = \frac{q^2}{6\pi\epsilon_0 c} \gamma^4 \left(\frac{\rho \omega^2}{c} \right)^2 = \frac{q^2}{6\pi\epsilon_0 c} \gamma^4 \beta^4 \frac{c^2}{\rho^2} = \frac{q^2 c}{6\pi\epsilon_0 \rho^2} \gamma^4 \beta^4$$

Energy loss per turn

$$\delta E = \frac{2\pi\rho}{\beta c} \frac{q^2 c}{6\pi\epsilon_0 \rho^2} \gamma^4 \beta^4 = \frac{q^2}{3\epsilon_0 \rho} \gamma^4 \beta^3$$

$$\delta E (\text{MeV}) = 8.85 \times 0.01 \frac{[E(\text{GeV})]^4}{\rho(\text{m})}$$