



# Physics 604

# Electromagnetic

# Theory I

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# Work



- Will seldom use this expression to find potential: Solve a differential equation!
- Moving a charge  $q$  from location  $\vec{x}_{before}$  to  $\vec{x}_{after}$  requires

$$W = q \left( \Phi(\vec{x}_{after}) - \Phi(\vec{x}_{before}) \right) \quad [\text{eV}]$$

- Falling through a potential difference causes energy gain to a positive charge
- Potential measured in volts [V] or [Nt m/C]

# Equations for Potential



- Poisson Equation

$$\nabla^2 \Phi \equiv \left[ \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right] \Phi = \frac{-\rho}{\epsilon_0}$$

- Laplace Equation ( $\rho = 0$ )

$$\nabla^2 \Phi = 0$$

- Verify Solution with vanishing Boundary Conditions at infinity

$$\begin{aligned}\Phi(\vec{x}) &= \frac{1}{4\pi\epsilon_0} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\rho(\vec{x}')}{|\vec{x} - \vec{x}'|} dx' dy' dz' \\ &\equiv \lim_{a \rightarrow 0} \frac{1}{4\pi\epsilon_0} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\rho(\vec{x}')}{\sqrt{(\vec{x} - \vec{x}')^2 + a^2}} dx' dy' dz'\end{aligned}$$

- Equation for scalar potential

$$\nabla^2 \Phi \equiv \lim_{a \rightarrow 0} \frac{-1}{4\pi\epsilon_0} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\rho(\vec{x}') 3a^2}{\left( (\vec{x} - \vec{x}')^2 + a^2 \right)^{5/2}} dx' dy' dz'$$

- Now

$$\lim_{a \rightarrow 0} \frac{3a^2}{\left( (\vec{x} - \vec{x}')^2 + a^2 \right)^{5/2}} = 0 \quad \vec{x} \neq \vec{x}'$$

- And

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{3a^2}{\left( (\vec{x} - \vec{x}')^2 + a^2 \right)^{5/2}} dx' dy' dz' = 12\pi \int_0^{\infty} \frac{x^2}{\left( x^2 + 1 \right)^{5/2}} dx = 4\pi$$

$$\therefore \lim_{a \rightarrow 0} \frac{3a^2}{\left( (\vec{x} - \vec{x}')^2 + a^2 \right)^{5/2}} = 4\pi \delta(\vec{x} - \vec{x}')$$

# Solution to Poisson Equation



$$\nabla^2 \Phi \equiv \frac{1}{4\pi\epsilon_0} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \rho(\vec{x}') 4\pi \delta(\vec{x} - \vec{x}') dx' dy' dz' = -\frac{\rho}{\epsilon_0}$$

- Representation

$$\nabla^2 \left[ \frac{1}{|\vec{x} - \vec{x}'|} \right] = -4\pi \delta(\vec{x} - \vec{x}')$$

- First Example of a *Green Function*. Made rigorous in the mathematical theory of *distributions*.

# Green's Identities



- Green's First Identity

$$\int_V \vec{\nabla} \cdot \vec{A} d^3x = \int_{S=\partial V} \vec{A} \cdot \vec{n} da$$

$$\vec{\nabla} \cdot (\phi \vec{\nabla} \psi) = \phi \nabla^2 \psi + \vec{\nabla} \phi \cdot \vec{\nabla} \psi$$

$\phi \vec{\nabla} \psi \cdot \vec{n} \equiv \phi \frac{\partial \psi}{\partial n}$  is the normal derivative of the field

$$\therefore \int_V (\phi \nabla^2 \psi + \vec{\nabla} \phi \cdot \vec{\nabla} \psi) d^3x = \int_{S=\partial V} \phi \frac{\partial \psi}{\partial n} da$$

- Green's Second Identity

$$\int_V (\phi \nabla^2 \psi - \psi \nabla^2 \phi) d^3x = \int_{S=\partial V} \left( \phi \frac{\partial \psi}{\partial n} - \psi \frac{\partial \phi}{\partial n} \right) da$$

# Formal Solution

$$\phi = \Phi$$

$$\psi = 1 / R = 1 / |\vec{x} - \vec{x}'|$$

$$\begin{aligned} & \int_V \left( -4\pi\Phi(\vec{x}')\delta(\vec{x} - \vec{x}') + \frac{1}{\epsilon_0 R}\rho(\vec{x}') \right) d^3x' \\ &= \int_{S=\partial V} \left( \Phi \frac{\partial}{\partial n'} \left( \frac{1}{R} \right) - \frac{1}{R} \frac{\partial \Phi}{\partial n'} \right) da' \end{aligned}$$

$$\vec{x} \in V \rightarrow$$

$$\Phi(\vec{x}) = \int_V \frac{1}{4\pi\epsilon_0 R} \rho(\vec{x}') d^3x' + \frac{1}{4\pi} \int_{S=\partial V} \left( \frac{1}{R} \frac{\partial \Phi}{\partial n'} - \Phi \frac{\partial}{\partial n'} \left( \frac{1}{R} \right) \right) da'$$

Surface integral vanishes at  $\infty$  if solutions vanish fast enough.

Surface integral carries the boundary conditions in the solutions!

# Uniqueness of Solutions



- Dirichlet Boundary Conditions

specify  $\phi(\vec{x})$  on a surface, i.e.  $\forall \vec{x} \in S$

- Neumann Boundary Conditions

specify  $\partial\phi(\vec{x})/\partial n$  on a surface, i.e.  $\forall \vec{x} \in S$

- Consider two different solutions  $\Phi_1$  and  $\Phi_2$ , and their difference  $U = \Phi_2 - \Phi_1$

$$\int_V \left( U \nabla^2 U + \vec{\nabla} U \cdot \vec{\nabla} U \right) d^3x = \int_{S=\partial V} U \frac{\partial U}{\partial n} da$$

$$\therefore \int_V \vec{\nabla} U \cdot \vec{\nabla} U d^3x = 0$$

# Solution with Specified BCs



- Use Green function method to solve Poisson Equation with specified BCs.

$$\nabla'^2 \left( \frac{1}{|\vec{x} - \vec{x}'|} \right) = -4\pi\delta(\vec{x} - \vec{x}')$$

Green Function Satisfies

$$\nabla'^2 G(\vec{x}, \vec{x}') = -4\pi\delta(\vec{x} - \vec{x}')$$

but with the boundary conditions of the problem!

$$G(\vec{x}, \vec{x}') = \frac{1}{|\vec{x} - \vec{x}'|} + F(\vec{x}, \vec{x}') \quad \nabla'^2 F(\vec{x}, \vec{x}') = 0$$

# Dirichlet Conditions



$$\begin{aligned}\Phi(\vec{x}) = & \int_V \frac{1}{4\pi\epsilon_0} \rho(\vec{x}') G(\vec{x}, \vec{x}') d^3x' \\ & + \frac{1}{4\pi} \int_{S=\partial V} \left( G(\vec{x}, \vec{x}') \frac{\partial \Phi}{\partial n'} - \Phi \frac{\partial G(\vec{x}, \vec{x}')}{\partial n'} \right) da'\end{aligned}$$

Demand

$$G_D(\vec{x}, \vec{x}') = 0 \quad \vec{x}' \in S$$

$$\Phi(\vec{x}) = \int_V \frac{1}{4\pi\epsilon_0} \rho(\vec{x}') G_D(\vec{x}, \vec{x}') d^3x' - \frac{1}{4\pi} \int_{S=\partial V} \Phi \frac{\partial G_D(\vec{x}, \vec{x}')}{\partial n'} da'$$