

Physics 604

Electromagnetic

Theory I

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Electrostatics



- Equations simplify enormously if the fields and sources do not depend on time
- Roughly the first half of this course deals with Static (time independent) Electric fields.
- Roughly the second half of the course will deal with Static Magnetic fields.
- In next semester (Physics 704/804) we'll study dynamic (time dependent phenomena)

- Electric field of a static charge q_1 at location \vec{x}_1 given by Coulomb's Law

$$\vec{E}(\vec{x}) = \frac{q_1}{4\pi\epsilon_0} \frac{\vec{x} - \vec{x}_1}{|\vec{x} - \vec{x}_1|^3}$$

Coulomb Interaction

- Notation

$$|\vec{v}| \equiv \sqrt{v_x^2 + v_y^2 + v_z^2} \rightarrow$$

$$|\vec{x} - \vec{x}_1|^3 = \left((x - x_1)^2 + (y - y_1)^2 + (z - z_1)^2 \right)^{3/2}$$

- Force on test particle with charge q at location \vec{x}

$$\vec{F}(\vec{x}) = q\vec{E}(\vec{x}) = \frac{qq_1}{4\pi\epsilon_0} \frac{\vec{x} - \vec{x}_1}{\left((x - x_1)^2 + (y - y_1)^2 + (z - z_1)^2 \right)^{3/2}}$$

- Dependence on the distance between charges

$$|\vec{F}| = \frac{qq_1}{4\pi\epsilon_0} \frac{1}{(x - x_1)^2 + (y - y_1)^2 + (z - z_1)^2} \propto \frac{1}{r^2}$$

Superposition Principal



- Note that likes repel and opposite signs attract
- Field from many charges

$$\vec{E}_{tot}(\vec{x}) = \sum_{i=1}^n \frac{q_i}{4\pi\epsilon_0} \frac{\vec{x} - \vec{x}_i}{|\vec{x} - \vec{x}_i|^3}$$

- Example of **SUPERPOSITION**, which arises because LHS of Maxwell Equations are linear differential operators.

$E_1(\vec{x})$ is a solution for M.E. with $\rho_1(\vec{x})$

$E_2(\vec{x})$ is a solution for M.E. with $\rho_2(\vec{x})$



$E_1(\vec{x}) + E_2(\vec{x})$ is a solution for M.E. with $\rho_1(\vec{x}) + \rho_2(\vec{x})$

Charge Density



- For a continuous static distribution of charge with charge density ρ [C/m³]

$$\vec{E}(\vec{x}) = \lim_{\Delta x, \Delta y, \Delta z \rightarrow 0} \sum_{i=-\infty}^{\infty} \sum_{j=-\infty}^{\infty} \sum_{k=-\infty}^{\infty} \frac{1}{4\pi\epsilon_0} \frac{\Delta q(i, j, k)}{\Delta x \Delta y \Delta z} \frac{\vec{x} - \vec{x}_{i,j,k}}{\left| \vec{x} - \vec{x}_{i,j,k} \right|^3} \Delta x \Delta y \Delta z$$
$$\rightarrow \frac{1}{4\pi\epsilon_0} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \rho(\vec{x}') \frac{(\vec{x} - \vec{x}')}{\left| \vec{x} - \vec{x}' \right|^3} dx' dy' dz'$$

- Recover Coulomb's law for a single charge if distribution for a single charge is Dirac delta-function

$$\rho(\vec{x}) = q_1 \delta(x - x_1) \delta(y - y_1) \delta(z - z_1)$$
$$\rightarrow \vec{E}(\vec{x}) = \frac{q_1}{4\pi\epsilon_0} \frac{\vec{x} - \vec{x}_1}{\left| \vec{x} - \vec{x}_1 \right|^3}$$

Properties of Dirac Delta “function”



- Zero-valued

$$\delta(x - a) = 0 \quad x \neq a$$

- Integrates to 1 if the domain of the integral contains a

$$\int_I \delta(x - a) dx = 1 \quad a \subset I,$$
$$= 0 \quad a \not\subset I$$

- Sifting property

$$\int_I f(x) \delta(x - a) dx = f(a) \quad a \subset I,$$
$$= 0 \quad a \not\subset I$$

- All domains of integration contain a going forward; same pattern applies and get zero if a not in the domain

- Derivative property

$$\int f(x)\delta'(x-a)dx = -f'(a)$$

- Delta of a function (x_i are the zeros of $f(x)$)

$$\delta(f(x)) = \sum_i \frac{1}{\left| \frac{df}{dx}(x_i) \right|} \delta(x - x_i)$$

- 3-D delta function

$$\delta(\vec{x} - \vec{X}) \equiv \delta(x - X)\delta(y - Y)\delta(z - Z)$$

$$\int_{\Delta V} \delta(\vec{x} - \vec{X}) d^3x = 1 \quad \vec{X} \subset \Delta V, \quad 0 \text{ otherwise}$$

Gauss's Law

- Easy proof:

$$\vec{\nabla} \cdot \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} = \frac{1}{|\vec{x} - \vec{x}'|^3} - \frac{3(x - x')^2}{|\vec{x} - \vec{x}'|^5}$$

$$+ \frac{1}{|\vec{x} - \vec{x}'|^3} - \frac{3(y - y')^2}{|\vec{x} - \vec{x}'|^5} + \frac{1}{|\vec{x} - \vec{x}'|^3} - \frac{3(z - z')^2}{|\vec{x} - \vec{x}'|^5} = 0 \quad \vec{x} \neq \vec{x}'$$

- Evaluate $\int (\vec{x} - \vec{x}') / |\vec{x} - \vec{x}'|^3 \cdot \vec{n} dA$ on a sphere surrounding \vec{x}'

$$\int_S \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} \cdot \vec{n} dA = \frac{1}{R^2} \int R^2 d\cos\theta d\phi = 4\pi$$

- Independent of R ! This is it! Do you see it!

Formal Statement

- Gauss's Law: for an *arbitrary* surface S (doesn't need to be a sphere!)

$$\int_S \vec{E} \cdot \vec{n} dA = \frac{q_{inside}}{\epsilon_0}$$

- Pf: If a charge is outside, calculation 1 shows it doesn't contribute to the integral. If a charge is inside, calculation 2 shows it contributes to the integral as given.
- For any V and $S = \partial V$

$$\int_V \nabla \cdot \vec{E} dx dy dz = \int_S \vec{E} \cdot \vec{n} dA = \frac{q_{inside}}{\epsilon_0} = \frac{1}{\epsilon_0} \int_V \rho(\vec{x}) dx dy dz$$

$$\therefore \nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0} \quad \text{1st Maxwell Eqn. in vacuum}$$

Electrostatic Approximation



- To fully solve for $\vec{E}(\vec{x})$ need equation for $\vec{\nabla} \times \vec{E}$

$$\vec{\nabla} \times \vec{E}(\vec{x}) = 0$$

- Consistent second Maxwell Eqn. when no induction
- Consistent with general expressions for field because

$$\vec{\nabla} \times \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} = 0$$

- In Mechanics, such force fields are called **conservative** because

$$\int_L \vec{E} \cdot d\vec{l} = \int_S \vec{\nabla} \times \vec{E} \cdot \vec{n} dA = 0$$

Scalar Potential

- There is a potential energy function, the **Scalar Potential** with

$$\vec{E} = -\vec{\nabla}\Phi$$

- Is a potential function for particle movement in the electric field
- Large simplification: computing one scalar function easier than three vector component functions
- General expression for scalar potential

$$\frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} = -\vec{\nabla} \left(\frac{1}{|\vec{x} - \vec{x}'|} \right) \rightarrow$$

$$\Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\rho(\vec{x}')}{|\vec{x} - \vec{x}'|} dx' dy' dz'$$

Work



- Will seldom use this expression to find potential: Solve a differential equation!
- Moving a charge q from location \vec{x}_{before} to \vec{x}_{after} requires

$$W = q \left(\Phi(\vec{x}_{after}) - \Phi(\vec{x}_{before}) \right) \quad [\text{eV}]$$

- Falling through a potential difference causes energy gain to a positive charge
- Potential measured in volts [V] or [Nt m/C]